

余剰次元と Gauge-Higgs 統一

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「LHCでの余剰次元研究」研究会

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1. Standard Model では
2. ゲージ・ヒッグス統一とは
3. ヒッグス場がゲージ場の一部になる
4. ヒッグス場はAB位相のゆらぎ
5. Effective Lagrangianでみる
6. Higgs couplings の予言 - SMからの大きなずれ
7. AB位相の量子効果でEW対称性が破れる
top quark がひきおこす

8. なんと $\theta_H = \frac{\pi}{2}$ が選ばれる
ZZH, Yukawa = 0 となる

9. ヒッグス粒子が安定になる

10. 暗黒物質 = ヒッグス

11. WMAP からヒッグス質量が決まる

12. LHC, ILC でどう見るか

13. Weak universalityの破れ

14. まとめ

Standard Model では

$$|D_\mu \Phi|^2 \quad y \bar{q}_L \Phi u_R \quad V[\Phi]$$

$$\Phi \sim \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix} \quad m_W = \frac{1}{2} g v$$

Gauge sector : established.

Higgs sector : yet to be unveiled.

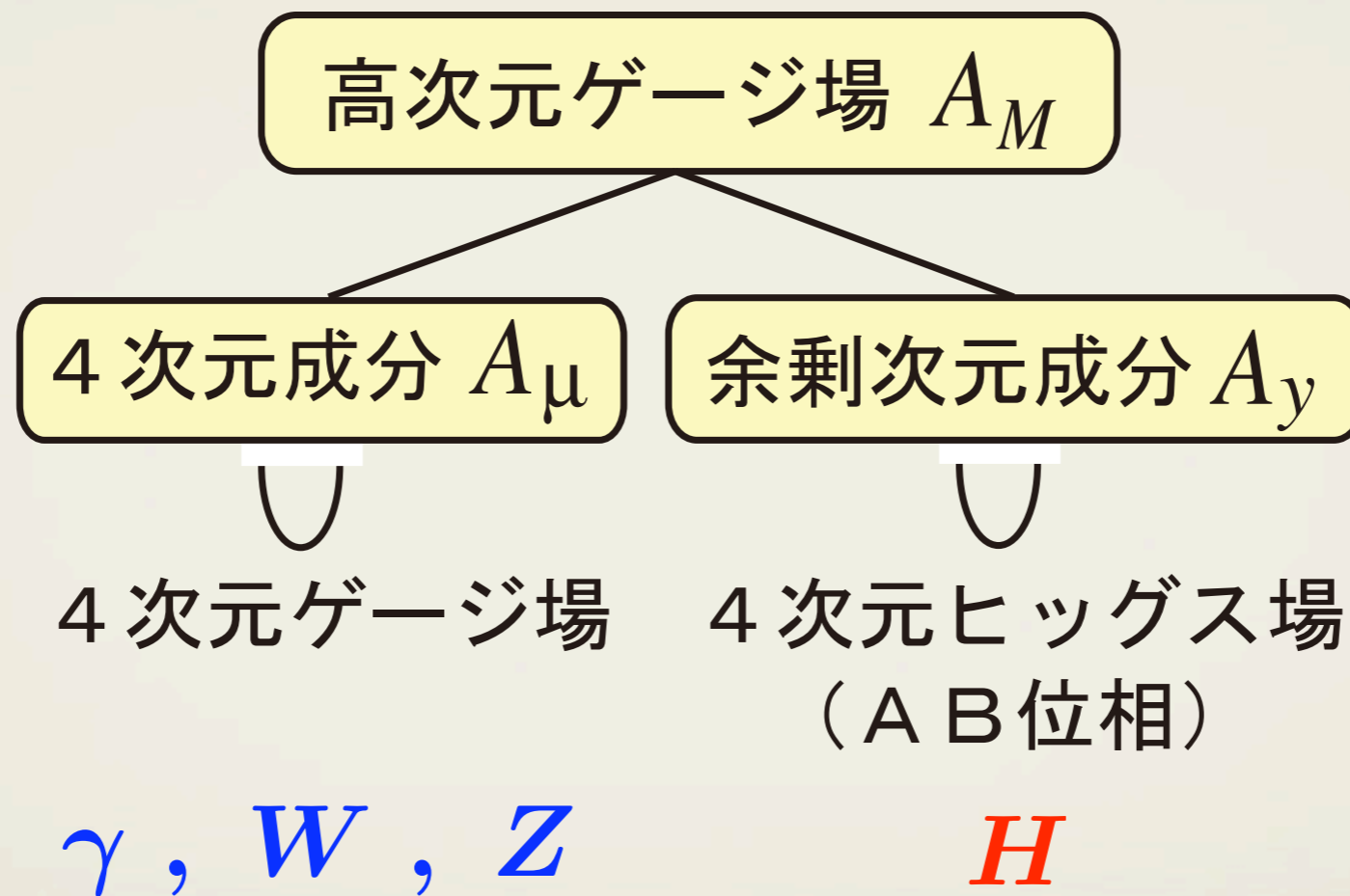
No principle

Many free parameters

Unnatural against radiative corrections

Hierarchy in fermion masses

ゲージ・ヒッグス統一とは



ヒッグス場はゲージ場の一部になる

Higgs mass

m_H dynamically generated
finite

Higgs couplings

H^3, H^4, \dots
 WWH, WWH^2, \dots
 ZZH, ZZH^2, \dots
 $\bar{\Psi}\Psi H, \dots$

all determined

Kaluza-Klein モード

$$M^4 \times S^1$$

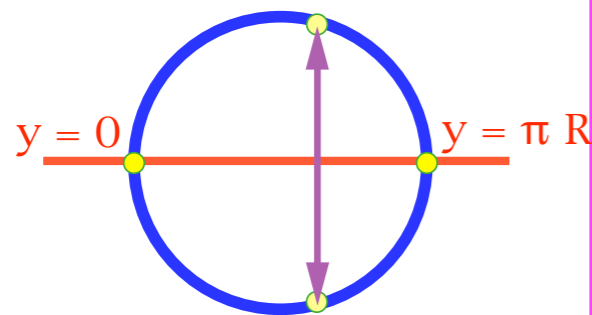
$$A_M(x, y) = \sum_n A_M^{(n)}(x) e^{iny/R}$$

KKモード $m_n = \frac{|n|}{R}$ KKスケール $m_{KK} = \frac{1}{R}$

$$M^4 \times (S^1/Z_2)$$

$$A_\mu(x, y) = \sum_{n=0}^{\infty} A_\mu^{(n)}(x) \cos \frac{ny}{R}$$

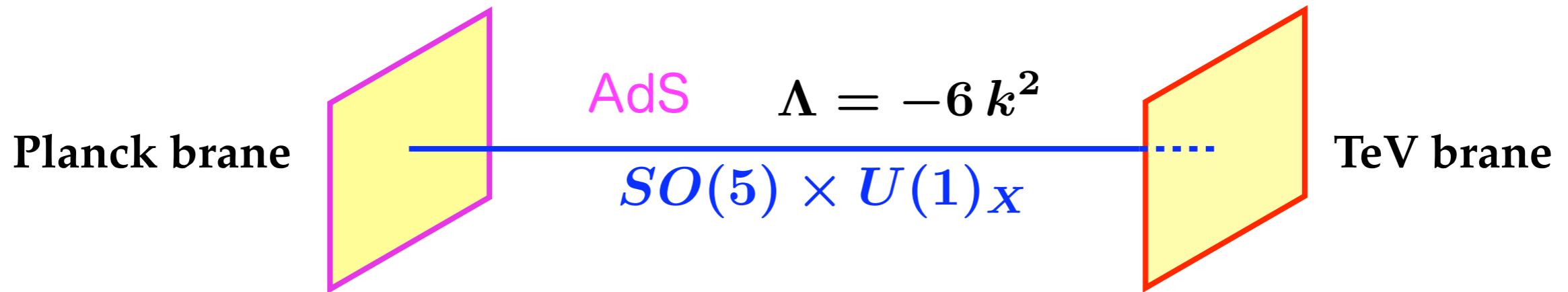
$$A_y(x, y) = \sum_{n=1}^{\infty} A_y^{(n)}(x) \sin \frac{ny}{R}$$



orbifold

A_μ のみ零モード

Randall-Sundrum フープ時空



$$\begin{pmatrix} A_\mu \\ A_y \end{pmatrix} (x, -y) = P_0 \begin{pmatrix} A_\mu \\ -A_y \end{pmatrix} (x, y) P_0^\dagger$$

$$\begin{pmatrix} A_\mu \\ A_y \end{pmatrix} (x, \pi R - y) = P_1 \begin{pmatrix} A_\mu \\ -A_y \end{pmatrix} (x, \pi R + y) P_1^\dagger$$

Orbifold BC: P_0, P_1

Origin of the Higgs doublet

Agashe, Contino, Pomarol 2005
 Hosotani, Sakamura 2006
 Medina, Shah, Wagner 2007

$$SO(5) \times U(1)_X$$

$$P_0 = P_1 = \begin{pmatrix} -1 & & & & \\ & -1 & & & \\ & & -1 & & \\ & & & -1 & \\ & & & & +1 \end{pmatrix}$$



$$W \quad Z \quad \gamma$$

$$A_\mu \sim \begin{pmatrix} \square \end{pmatrix}$$

$$SO(5) \rightarrow SO(4) \simeq SU(2)_L \times SU(2)_R$$



Higgs

$$A_y \sim \begin{pmatrix} \begin{matrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \end{matrix} \\ \square \end{pmatrix} \quad \Phi = \begin{bmatrix} \phi_1 + i\phi_2 \\ \phi_4 - i\phi_3 \end{bmatrix}$$

Model on RS

YH, Oda, Ohnuma, Sakamura 2008
(YH, Noda, Uekusa 2009)

Planck brane

$$A_M \times B_M$$

$$SO(5) \times U(1)_X$$

TeV brane

$$\begin{pmatrix} \hat{T}_R \\ \hat{B}_R \end{pmatrix}$$

$$\begin{pmatrix} \hat{U}_R \\ \hat{D}_R \end{pmatrix}$$

$$\begin{pmatrix} \hat{X}_R \\ \hat{Y}_R \end{pmatrix}$$

$$\begin{pmatrix} T \\ B \\ t \\ b \\ t' \end{pmatrix}_{\frac{2}{3}}$$

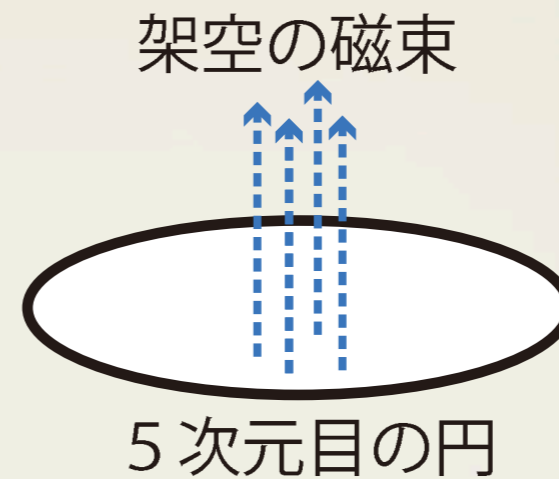
$$\begin{pmatrix} U \\ D \\ X \\ Y \\ b' \end{pmatrix}_{-\frac{1}{3}}$$

Anomaly-free

At low energies

$$\gamma, W, Z, H, \begin{pmatrix} t_L \\ b_L \end{pmatrix}, t'_R, b'_R, \dots$$

ヒッグス場はAB位相のゆらぎ



$$A_y^{45}(x, y) = h_0(y) \left\{ \theta_H + \frac{H(x)}{f_H} \right\}$$

$$\theta_H \sim \theta_H + 2\pi$$

Effective interactions

$$\text{AB phase} \quad \hat{\theta}_H = \theta_H + \frac{H}{f_H} \quad f_H = \frac{2}{\sqrt{kL}} \frac{m_{KK}}{\pi g}$$

$$\mathcal{L}_{\text{eff}} \sim -V_{\text{eff}}(\hat{\theta}_H) \quad \text{YH 1983}$$

$$-m_W (\hat{\theta}_H)^2 W_\mu^\dagger W^\mu - \frac{1}{2} m_Z (\hat{\theta}_H)^2 Z_\mu Z^\mu$$

YH-Sakamura 2006, 2007

$$-m_f (\hat{\theta}_H) \bar{\psi}_f \psi_f \quad \text{YH-Kobayashi 2008}$$

$$\mathcal{L}_{\text{eff}}(\hat{\theta}_H + 2\pi) = \mathcal{L}_{\text{eff}}(\hat{\theta}_H) \quad \textit{finite}$$

SO(5)xU(1) model

SM

$$m_W(\hat{\theta}_H) \sim \frac{1}{2} g f_H \sin \hat{\theta}_H$$

$$\frac{1}{2} g(v + H)$$

$$m_Z(\hat{\theta}_H) \sim \frac{1}{2 \cos \theta_W} g f_H \sin \hat{\theta}_H$$

$$\frac{1}{2 \cos \theta_W} g(v + H)$$

$$m_f(\hat{\theta}_H) \sim \lambda_f \sin \hat{\theta}_H$$

$$y_f(v + H)$$

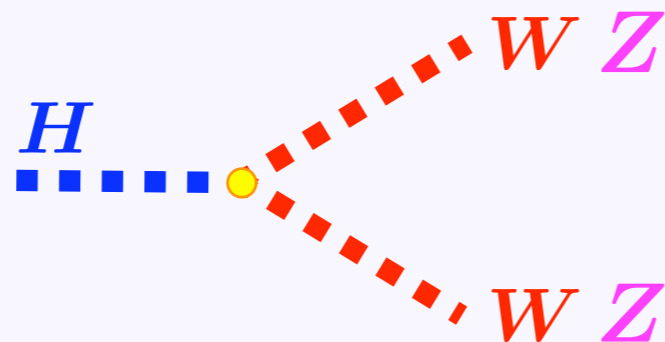
$$\hat{\theta}_H = \theta_H + \frac{H}{f_H}$$

周期性、非線形性

Prediction 1

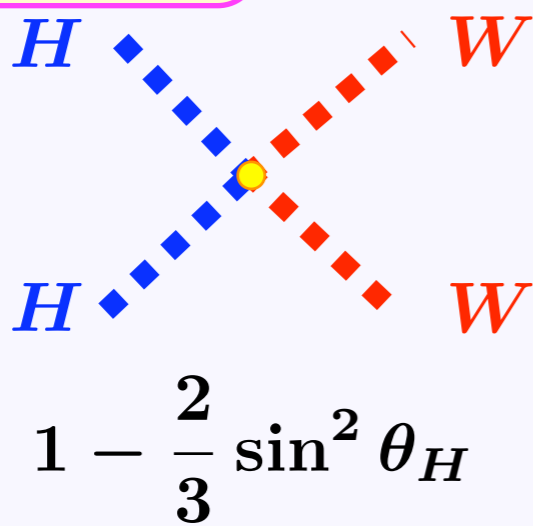
$$-m_W (\hat{\theta}_H)^2 W_\mu^\dagger W^\mu - \frac{1}{2} m_Z (\hat{\theta}_H)^2 Z_\mu Z^\mu$$

WWH
ZZH



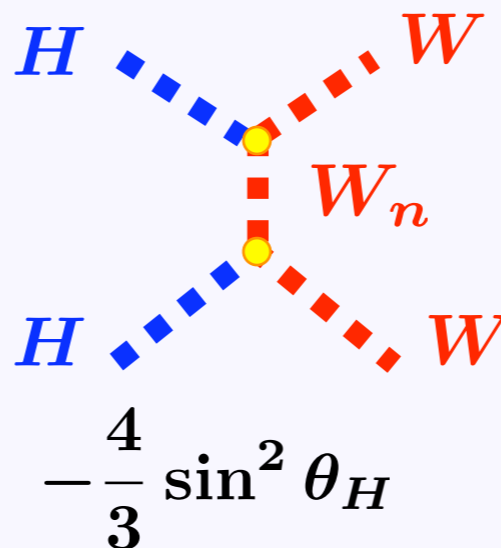
$$SM \times \cos \theta_H$$

WWHH
ZZHH



$$1 - \frac{2}{3} \sin^2 \theta_H$$

$$+ \sum_n$$



$$-\frac{4}{3} \sin^2 \theta_H$$

$$SM \times \cos 2\theta_H$$

4D gauge and Higgs couplings

$$F_{MN}^2 \sim (\partial_\mu A_\nu - \partial_\nu A_\mu + g[A_\mu, A_\nu])^2$$

WWZ
WWZZ
WWWW

Almost universal in RS

(Large deviation in flat space)

$$+(\partial_\mu A_y - \partial_y A_\mu + g[A_\mu, A_y])^2$$

WWH
ZZH
WWHH
ZZHH

significant θ_H -dependence

Sakamura-Hosotani 2006

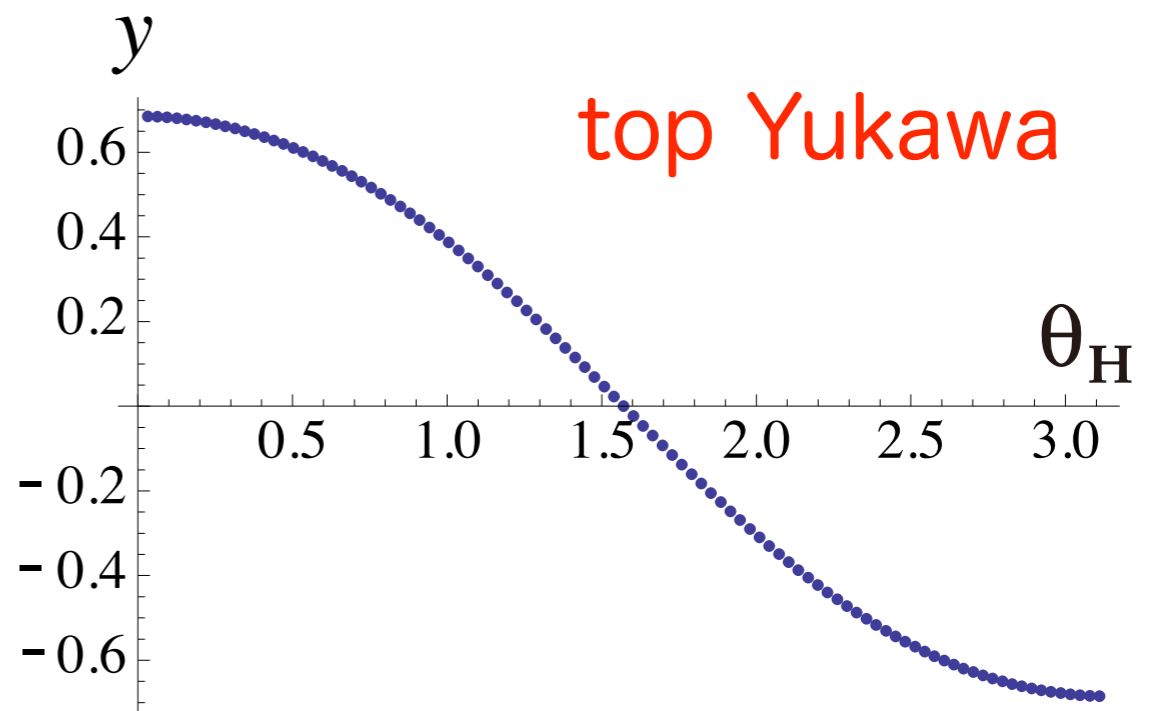
Prediction 2

$$m_f \sim \lambda_f \sin \theta_H$$

$$y_f \sim \frac{\lambda_f}{f_H} \cos \theta_H$$

Yukawa couplings

$$-m_f(\hat{\theta}_H) \bar{\psi}_f \psi_f$$

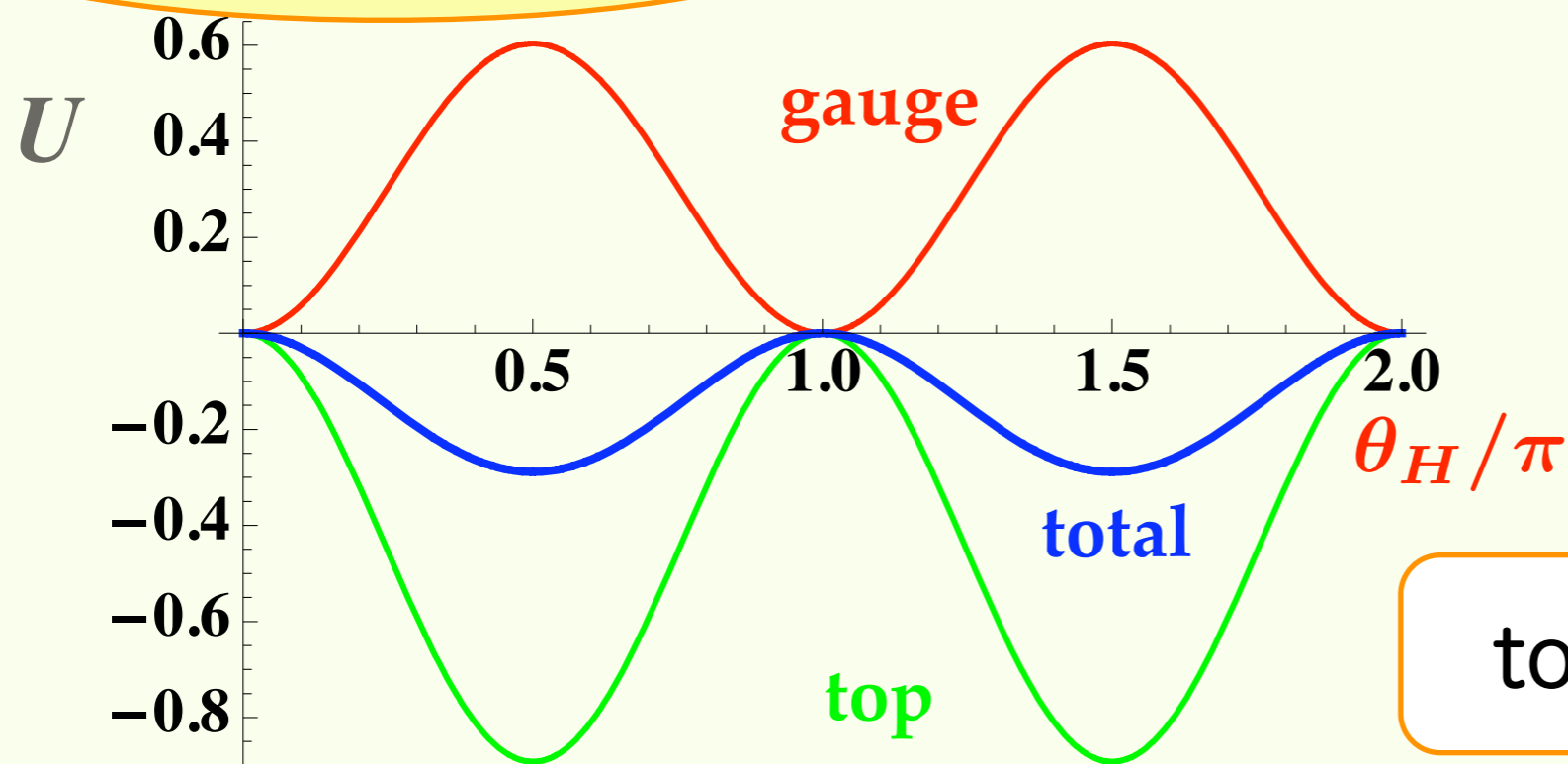


Prediction 3

AB位相の量子効果でEW対称性が破れる

Hosotani mechanism 1983

Effective potential



top quark がひきおこす

このモデルでは

$$m_H \sim 50 \text{ GeV} !$$

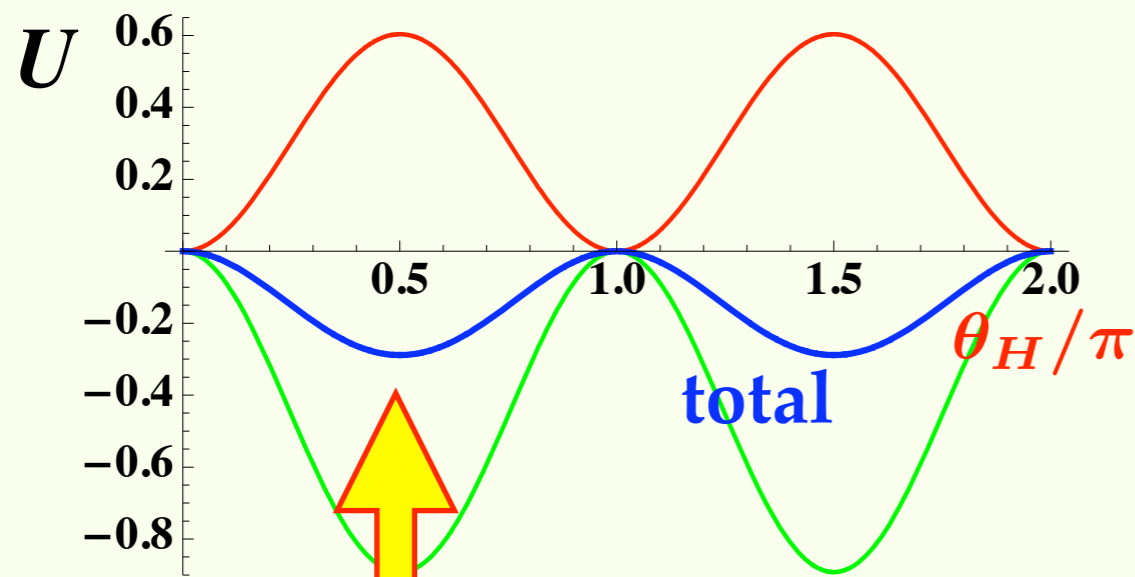
Prediction 4

$z_L = e^{kL}$	k (GeV)	m_{KK} (TeV)	c	m_H (GeV)
10^{17}	5.0×10^{19}	1.58	0.435	54.4
10^{15}	4.7×10^{17}	1.48	0.426	50.8
10^{13}	4.4×10^{15}	1.38	0.413	47.0
10^{10}	3.9×10^{12}	1.21	0.384	40.6
1.30×10^4	3.2×10^6	0.78	0.	24.5

$$m_{KK} \sim 1.5 \text{ TeV}$$

Light Higgs

(Its value depends on details of the model.)



Minima at

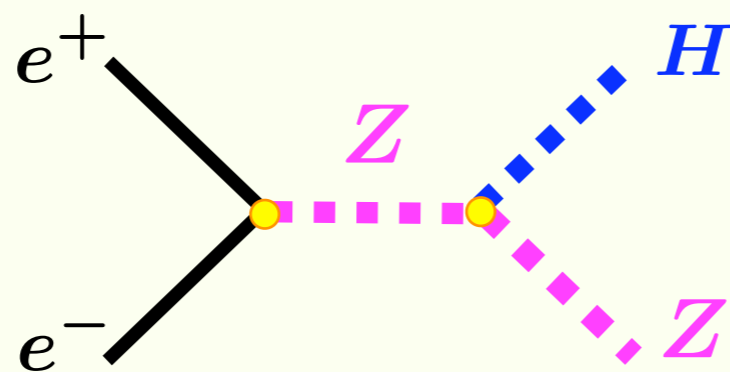
$$\theta_H = \frac{\pi}{2}$$



$$\theta_H = \frac{\pi}{2} \text{ が選ばれる}$$



WWH, ZZH, Yukawa = 0 となる!



$$= 0$$

LEP2 bound is evaded.

安定なヒッグスボゾン



Under certain conditions, Higgs bosons become absolutely stable to all orders.

$$\text{“ } \theta_H = \frac{\pi}{2} \text{ ”}$$

SO(5)×U(1) ゲージ・ヒッグス統合理論の対称性

YH, P. Ko, M. Tanaka, arXiv:0908.0212 (PLB)

Mirror reflection symmetry

$$\begin{aligned}(x^\mu, y) &\rightarrow (x^\mu, -y) \\ (A_\mu, A_y) &\rightarrow (A_\mu, -A_y), \quad \Psi \rightarrow \pm \gamma^5 \Psi\end{aligned}$$

invariant under $\hat{\theta}_H = \theta_H + \frac{H}{f_H} \rightarrow -\hat{\theta}_H$

Enhanced gauge symmetry

条件: all bulk fermions は SO(5)のvector rep.

invariant under $\hat{\theta}_H \rightarrow \hat{\theta}_H + \pi$

top quark



$$\theta_H = \frac{\pi}{2}$$

$$\frac{\pi}{2} + \frac{H}{f_H} \sim -\frac{\pi}{2} - \frac{H}{f_H} \sim \frac{\pi}{2} - \frac{H}{f_H}$$

理論は $H \rightarrow -H$ で不変

H-parity

$H : -$

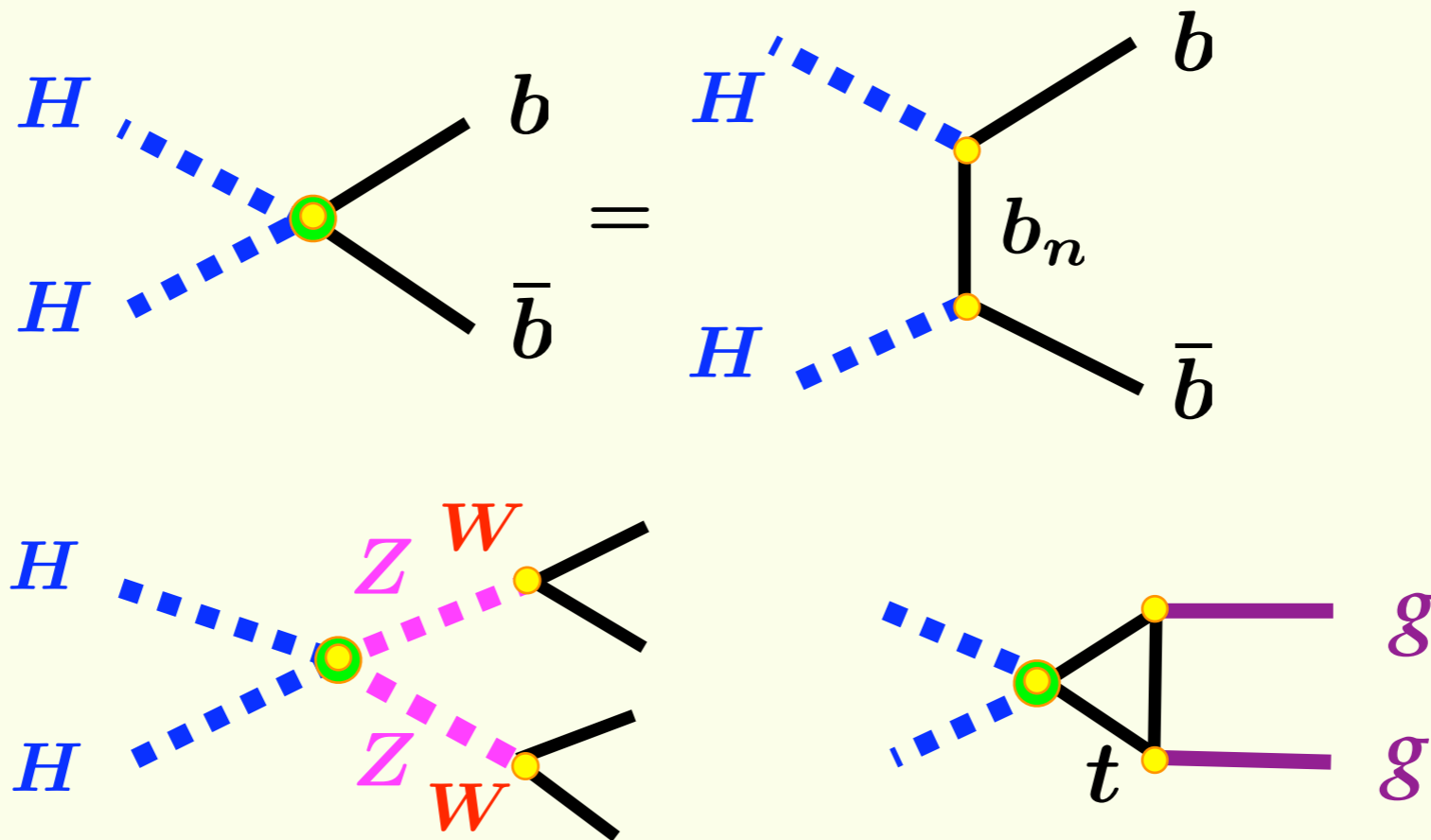
all other SM particles : +

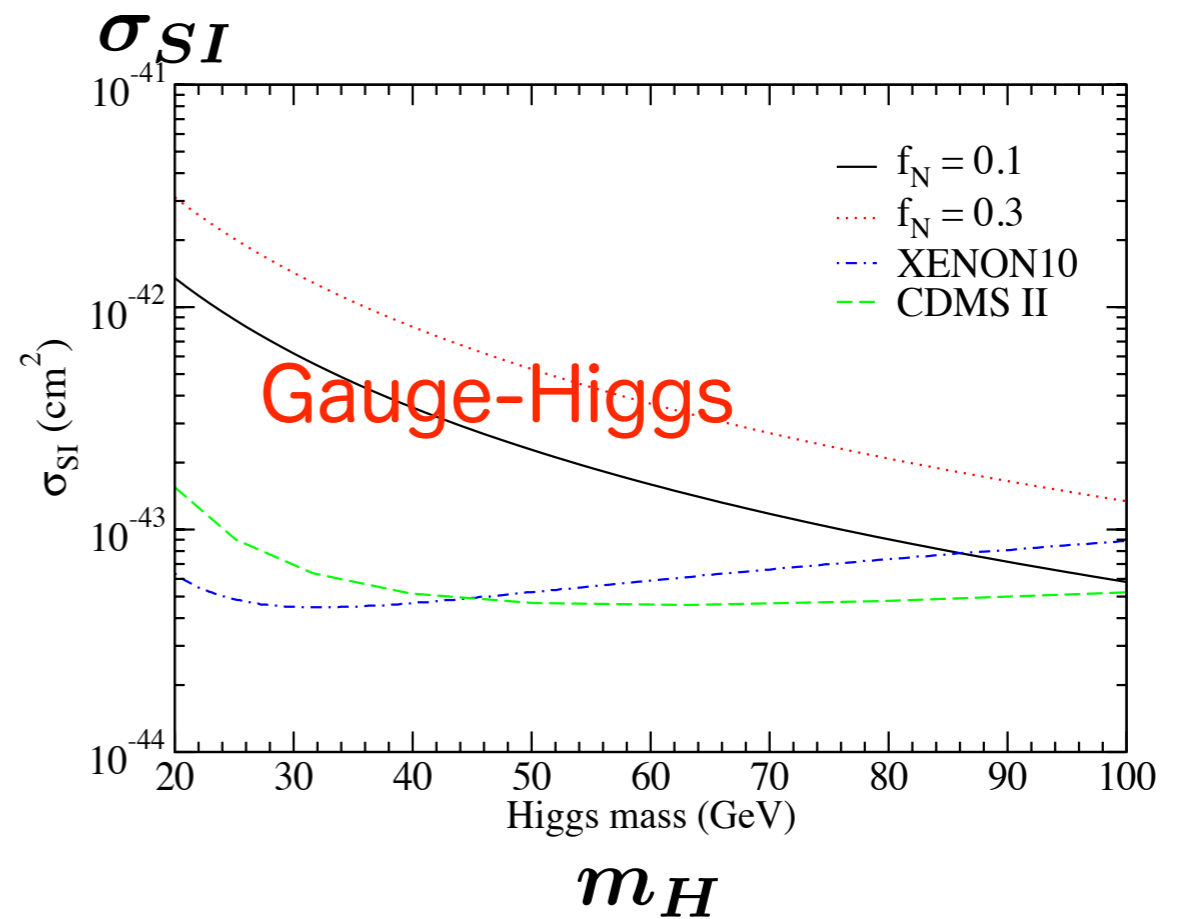
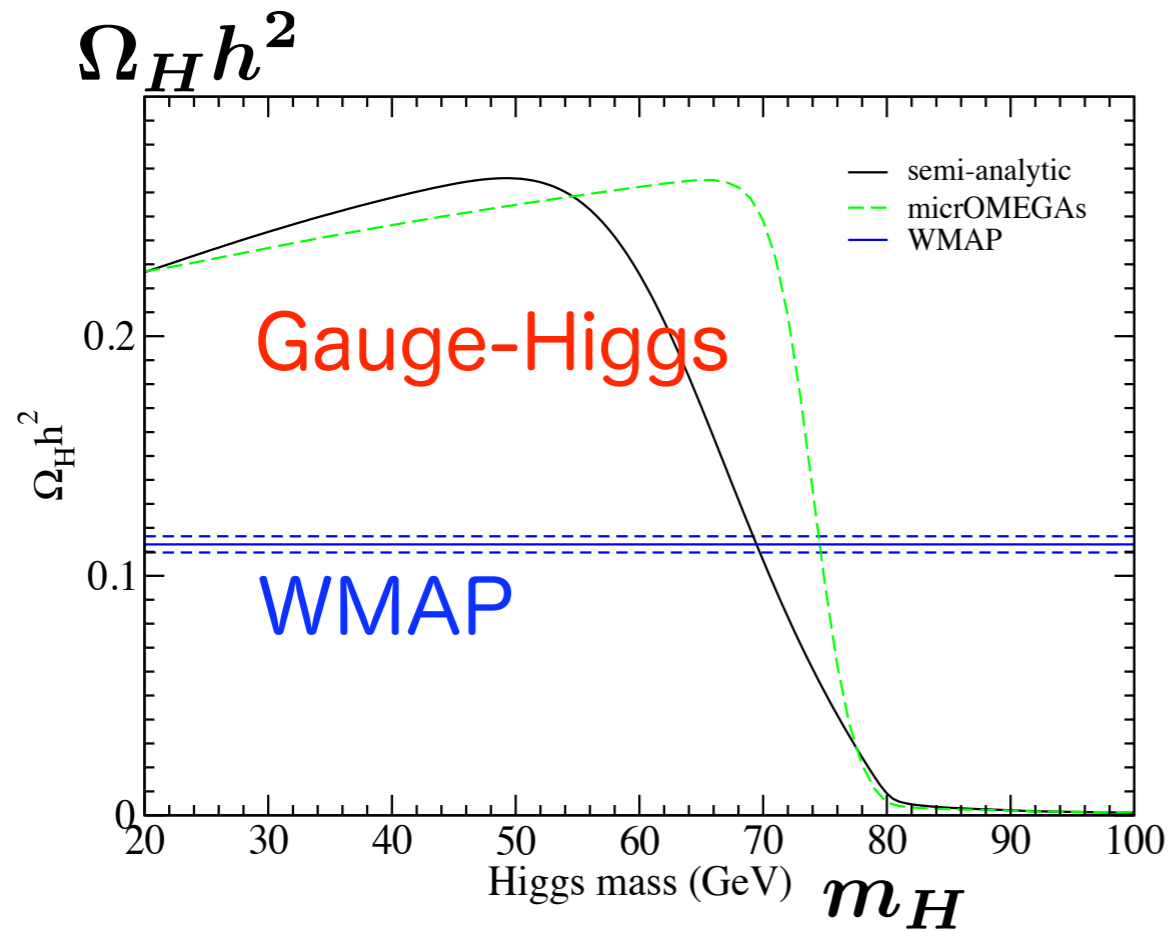
Higgs bosons become absolutely stable.

$WWH, ZZH, \text{Yukawa} = 0$

暗黒物質 = ヒッグス

Annihilation



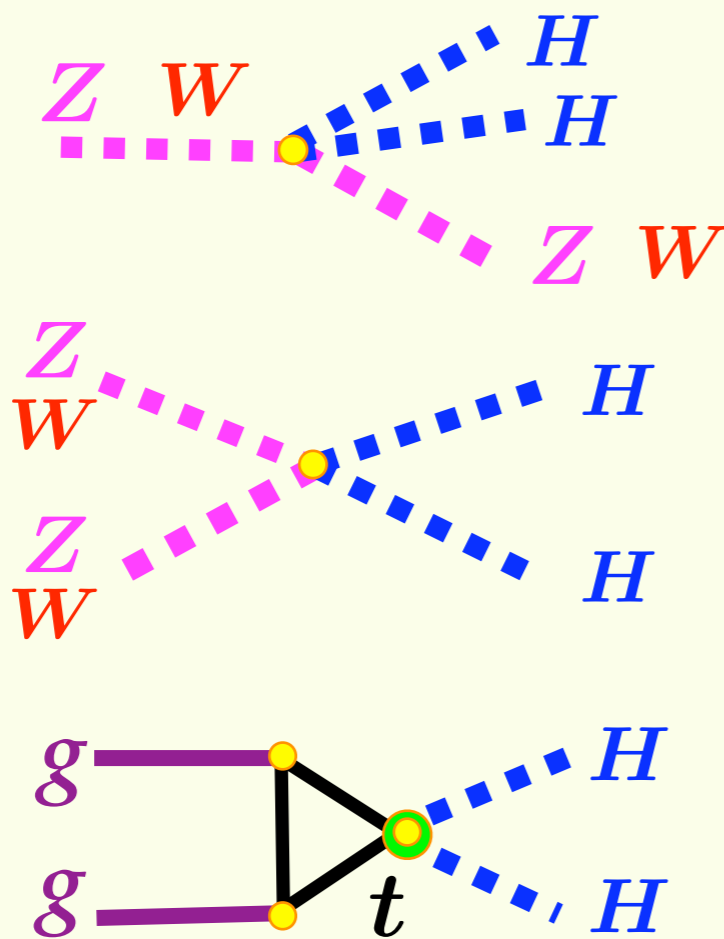


WMAP data $\rightarrow m_H \sim 70 \text{ GeV}$

With many uncertainties it is premature to exclude the Higgs DM scenario from the CDMS II and XENON10 bounds.

LHC, ILC でヒッグスをどうみるか

Production:



だが、ヒッグスは安定

ヒッグス粒子

=

missing energy,
missing momentum

KK モード ?

g_n

W_n, Z_n, γ_n

Weak universality の破れ

WWZ coupling $\sim 0.2\%$

μ - e , τ - e , t - e universality violation

2×10^{-8} 4×10^{-2}
 5×10^{-6}

Summary

ゲージ・ヒッグス統一では

Higgs couplings がSMより大きくずれる

安定なヒッグス粒子

暗黒物質=ヒッグス $m_H \sim 70 \text{ GeV}$

新しい可能性

ヒッグス=missing energy, momentum