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The Valley Hall Effect in MoS₂ Transistors Kagimura

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Valleytronics



Prospective

Outline

- I. Anomalous velocity
- 2. The valley Hall effect
- 3. Valley-selective circular dichroism of monolayer MoS₂
- 4. The valley Hall effect in MoS₂ transistors

I. Anomalous velocity

$$\Omega_{\mu\nu}^{n} = i \left[\left\langle \frac{\partial n(\mathbf{R})}{\partial R^{\mu}} \middle| \frac{\partial n(\mathbf{R})}{\partial R^{\nu}} \right\rangle - (\nu \leftrightarrow \mu) \right]$$

 $|n(\mathbf{R})\rangle$: instantaneous eigenstates of Hamiltonian at each value of \mathbf{R} i.e. $H(\mathbf{R}) |n(\mathbf{R})\rangle = \varepsilon_n(\mathbf{R}) |n(\mathbf{R})\rangle$

Hamiltonian and eigenstates

Independent electron approximation

$$H = \frac{\hat{\mathbf{p}}^2}{2m} + V(\mathbf{r})$$
Bloch's theorem: $\psi_{n,\mathbf{q}}(\mathbf{r} + \mathbf{a}) = e^{i\mathbf{q}\cdot\mathbf{a}}\psi_{n\mathbf{q}}(\mathbf{r})$
crystal momentum
 \mathbf{q} -dependent Hamiltonian: $H(\mathbf{q}) = e^{-i\mathbf{q}\cdot\mathbf{r}}He^{i\mathbf{q}\cdot\mathbf{r}}$
 $= \frac{(\hat{\mathbf{p}} + \mathbf{q})^2}{2m} + V(\mathbf{r})$
 $u_{n,\mathbf{q}}(\mathbf{r}) = e^{-i\mathbf{q}\cdot\mathbf{r}}\psi_{n\mathbf{q}}(\mathbf{r})$

Wave function

 $\mathbf{q}(t)$: changes slowly in time

time-dependent Scrhodinger equation: $i \frac{\partial}{\partial t} |\psi(t)\rangle = H(t) |\psi(t)\rangle$

instantaneous eigenstates

$$|\psi(t)\rangle = \sum_{n} \exp\left(-i \int_{t_0}^t dt' E_n(t')\right) a_n(t) |n(t)\rangle$$

$$\dot{a}_n(r) = -\sum_l a_l(t) \langle n(t) | \frac{\partial}{\partial t} | l(t) \rangle \exp\left(-i \int_{t_0}^t dt' [E_l(t') - E_n(t')]\right)$$

First-order wave function

We have an initial condition: $a_n(0) = 1, \ a_{n'} = 0$ For $n' \neq n$

$$\frac{\partial}{\partial t}a_{n'} = -\left\langle n'\right|\frac{\partial}{\partial t}\left|n\right\rangle\exp\left(-i\int_{t_0}^t dt'[E_n(t') - E_{n'}(t')]\right)$$

Integrate by parts

 ψ

$$a_{n'} = -\frac{\langle n'|\partial/\partial t |n\rangle}{E_n - E_{n'}} i \exp\left(-i \int_{t_0}^t dt' [E_n(t') - E_{n'}(t')]\right)$$
$$t) \rangle = \exp\left(-i \int_{t_0}^t dt' E_n(t')\right) \left\{ |n(t)\rangle - i \sum_{n \neq n'} |n(t)\rangle \frac{\langle n|\partial/\partial t |n\rangle}{E_n - E_{n'}} \right\}$$

The average velocity

Apart from the overall phase factor,

$$|u_n\rangle - i\sum_{n'\neq n} \frac{|n_{n'}\rangle \langle u_{n'}| \partial u_n/\partial t\rangle}{\varepsilon_n - \varepsilon_{n'}}$$

velocity operator in the q representation : $v(q,t) = \partial H(q,t)/\partial q$

$$v_{n}(q) = \frac{\partial \epsilon_{n}(q)}{\partial q} - i \sum_{n' \neq n} \left\{ \frac{\langle n_{n} | \partial H / \partial q | u_{n'} \rangle \langle u_{n'} \partial u_{n} / \partial t \rangle}{\varepsilon_{n} - \varepsilon_{n'}} - \text{c.c.} \right\}$$

$$v_n(q) = \frac{\partial \varepsilon_n(q)}{\partial q} - i \left[\left\langle \frac{\partial u_n}{\partial q} \middle| \frac{\partial u_n}{\partial t} \right\rangle - \left\langle \frac{\partial u_n}{\partial t} \middle| \frac{\partial u_n}{\partial q} \right\rangle \right]$$

Berry curvature Ω_{qt}^n

Der

With electric field

uniform vector potential A(t)

$$H(t) = \frac{[\hat{\mathbf{p}} + e\mathbf{A}(t)]^2}{2m} + V(\mathbf{r})$$

q-space $H(\mathbf{q}, t) = H(\mathbf{q} + e\mathbf{A}(t))$

gauge invariant crystal momentum: $\mathbf{k} = \mathbf{q} + e\mathbf{A}(t)$ eigenstates: $u_n(\mathbf{k})$

 ${f A}(t)$ preserve the translational symmetry, ${f \dot q}=0$ $\dot {f k}=-e{f E}$



Berry curvature: $\Omega_n(\mathbf{k}) = i \langle \nabla_{\mathbf{k}} u_n(\mathbf{k}) | \times | \nabla_{\mathbf{k}} u_n(\mathbf{k}) \rangle$

anomalous velocity: $-e\mathbf{E} \times \mathbf{\Omega}_n(\mathbf{k})$

transverse to the electric field give rise a Hall current

2. The valley Hall effect

Symmetry consideration

$$\mathbf{v}_n(\mathbf{k}) = \frac{\partial \varepsilon_n(\mathbf{k})}{\partial \mathbf{k}} - e\mathbf{E} \times \mathbf{\Omega}_n(\mathbf{k})$$

 $\mathbf{v} \rightarrow -\mathbf{v}, \ \mathbf{k} \rightarrow -\mathbf{k}, \ \mathbf{E} \rightarrow \mathbf{E}$
time-reversal sym. $\mathbf{\Omega}_n(-\mathbf{k}) = -\mathbf{\Omega}_n(\mathbf{k})$
 $\mathbf{v} \rightarrow -\mathbf{v}, \ \mathbf{k} \rightarrow -\mathbf{k}, \ \mathbf{E} \rightarrow -\mathbf{E}$
spatial inversion sym. $\mathbf{\Omega}_n(-\mathbf{k}) = \mathbf{\Omega}_n(\mathbf{k})$
time-reversal and spatial inversion symmetry
 $\mathbf{\Omega}_n(\mathbf{k}) = 0$

2-valley structure



Valley Hall effect

valley index: $\tau_z = \pm 1$

Valley Hall Effect: $\mathbf{j}^v = \sigma_H^v \hat{z} \times \mathbf{E}$ valley Hall conductivity

valley current:
$$\mathbf{j}^v = \langle au_z \mathbf{v}
angle$$



Berry curvature of graphene



Valley Hall current

Anomalous velocity: $\mathbf{v}_n(\mathbf{k}) = \frac{\partial \varepsilon_n(\mathbf{k})}{\partial \mathbf{k}} - e\mathbf{E} \times \mathbf{\Omega}_n(\mathbf{k})$

 $\mathbf{j}^{v} = \langle \tau_{z} \mathbf{v} \rangle$ $= -2eE \frac{3a^{2}t^{2}\Delta}{2(\Delta^{2} + 3q^{2}a^{2}t^{2})^{3/2}}$

Valley Hall current!

Pathway to valleytronics



It provides a new and standard pathway to potential applications of valleytronics in a broad class of semiconductors.

3. Valley-selective circular dichroism of monolayer MoS₂

[Ting Cao, Ji Feng, Junren Shi, Qian Niu, Enge Wang,(2012)]

Graphene?

graphene: inversion symmetry



Unprecedented control of the lattice structure on the scale of a single atom is required.

Valleytronics in graphene is very difficult



b



inversion symmetry

MoS₂ for valleytronics



Degree of circular polarization

difference between the absorption of left- and righthanded lights, between the top of the valence bands and the bottom of conduction bands

$$\eta(\mathbf{k}, \omega_{cv}) = \frac{|\mathcal{P}_{\pm}^{cv}(\mathbf{k})|^2 - |\mathcal{P}_{-}^{cv}(\mathbf{k})|^2}{|\mathcal{P}_{\pm}^{cv}(\mathbf{k})|^2 + |\mathcal{P}_{-}^{cv}(\mathbf{k})|^2}$$

the sincular polarization:
$$\mathcal{P}_{\pm}^{cv} = \frac{1}{\sqrt{2}} [P_x^{cv}(\mathbf{k}) \pm i P_y^{cv}(\mathbf{k})]$$

inter-band matrix elements: $\mathbf{P}^{cv}(\mathbf{k}) = \langle \psi_{c\mathbf{k}} | \hat{\mathbf{p}} | \psi_{v\mathbf{k}} \rangle$

 $\hbar\omega_{cv}(\mathbf{k}) = \varepsilon_{c}(\mathbf{k}) - \varepsilon_{v}(\mathbf{k})$ $\mathcal{P}_{\pm}^{cv}(\mathbf{k}) = 1/\sqrt{2} [P_x^{cv}(\mathbf{k}) \pm i P_y^{cv}(\mathbf{k})]$ $\mathbf{P}^{cv}(\mathbf{k}) = \langle \psi_{c\mathbf{k}} \mid \mathbf{p} \mid \psi_{v\mathbf{k}} \rangle$ η at K^{\pm} in MoS₂





Berry curvature



Hall effect



4. The valley Hall effect in MoS₂ transistors

Monolayer MoS₂ Hall bar device







Doping dependence of the anomalous Hall conductivity



Open question

- What does cause the observed quick drop of Hall conductance at laser frequency? Inter-corn scattering?
- What are the relevant sources of scattering present in MoS₂, and what are their relative strengths?
- What are the next steps towards the long standing goal of valleytronics?

Summary

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