

Anomaly polynomial of general 6d SCFTs

Kantaro Ohmori

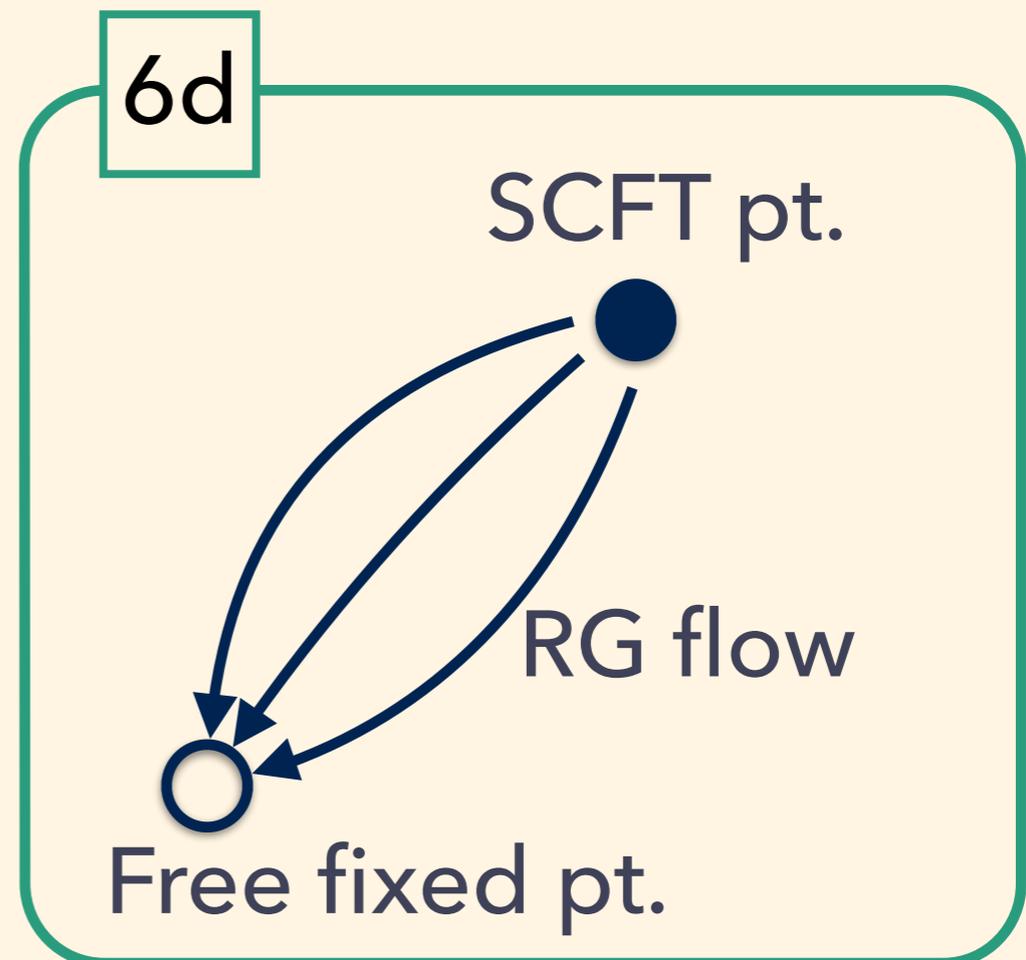
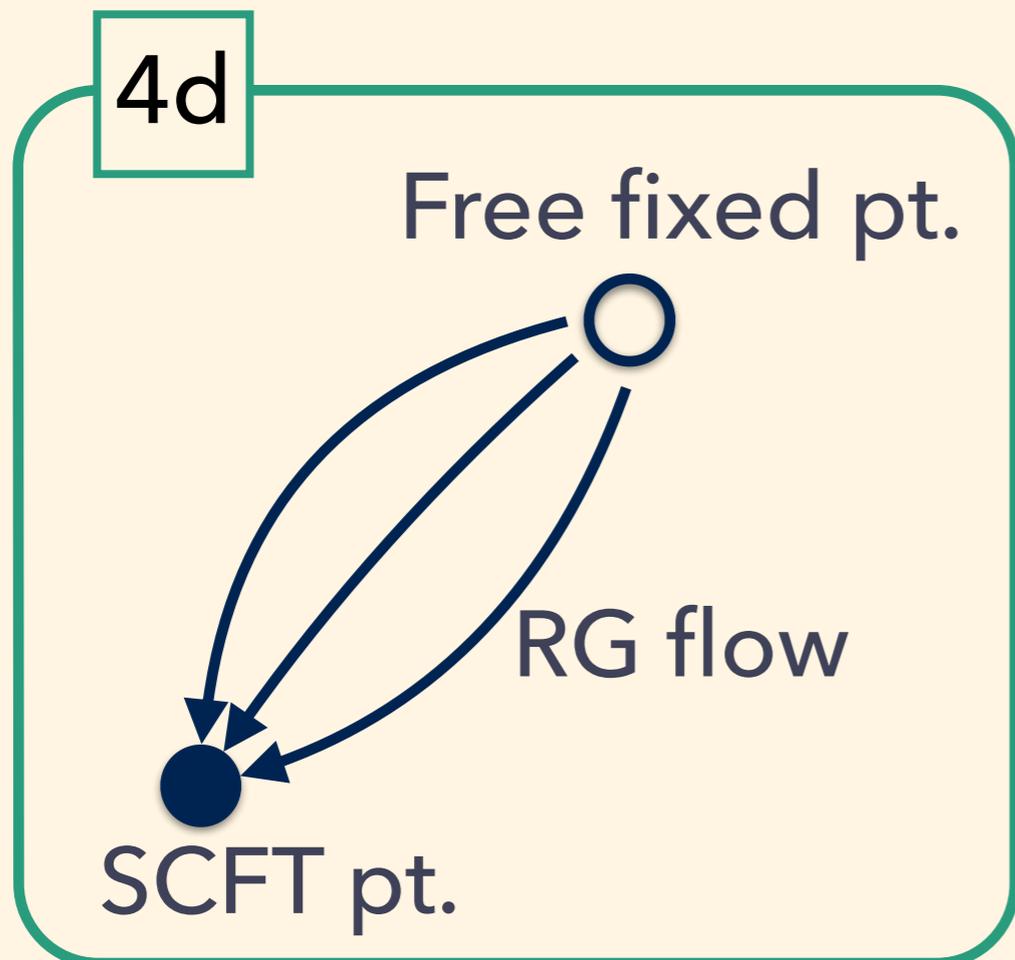
w/ Hiroyuki Shimizu, Yuji Tachikawa, Kazuya Yonekura

arXiv:1408.5572, 1????.?????

Introduction

Interacting 6d SCFTs

- $\mathcal{N}=(2,0)$ or $\mathcal{N}=(1,0)$ (chiral theories)
- UV fixed point



Why 6d SCFTs?

- Strongly coupled system
: no known Lagrangian
- Compactification \Rightarrow low dimensional systems
- Might control low dimensional dualities
e.g. $\mathcal{N}=(2,0)$ theories of ADE type

What's done?

- String (M,F-) theoretical constructions

[Witten '96],[Ganor,Hanany '96],[Seiberg '97],[Intrilligator Blum '97]

[Brunner,Karch '97],[Hanany,Zaffaroni '97],[Aspinwall,Morrison '97] etc.

[Heckman,Morrison,Vafa '13][Gaiotto Tomasiello '14],

[Del-Zotto,Heckman,Tomasiello,Vafa '14]

- Enormous works for $\mathcal{N}=(2,0)$ theories

[Gaiotto '09]...

- Anomaly polynomials

for above $\mathcal{N}=(1,0)$ theories

- Inflow:[Freed, Harvey, Minasian, Moore, '98],[KO,Shimizu,Tachikawa '14]

- Tensor branch:[Intrilligator '14],[KO,Shimizu,Tachikawa,Yonekura '14]

Plan

- Brane constructions of 6d SCFTs
- Anomaly polynomial
- Tensor branch anomaly matching
- Torus compactifications (if time permits)

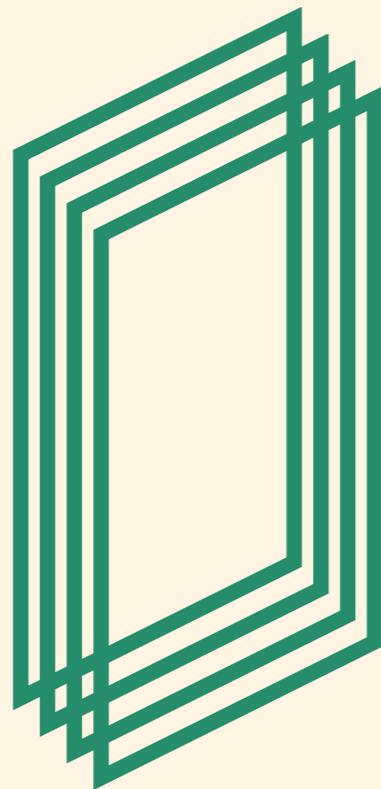
Brane constructions of 6d SCFTs

6d $\mathcal{N}=(1,0)$ Supermultiplets

- 8 supercharges, $SU(2)$ R-symmetry
- Tensor multiplet: $(B_{\mu\nu}^+, \psi^+, a)$
 $a \in \mathbb{R}$: “tensor branch” vev (preserves R-sym)
- Vector multiplet: (A_μ, λ^-)
No scalar
- Hyper multiplet: (ϕ_i, ψ^+) $i = 1, 2, 3, 4$
 $\phi \in \mathbb{R}^4$: “Higgs branch” vev (breaks R-sym)

$\mathcal{N}=(2,0)$ theory of A -type

Q coincident M5-branes

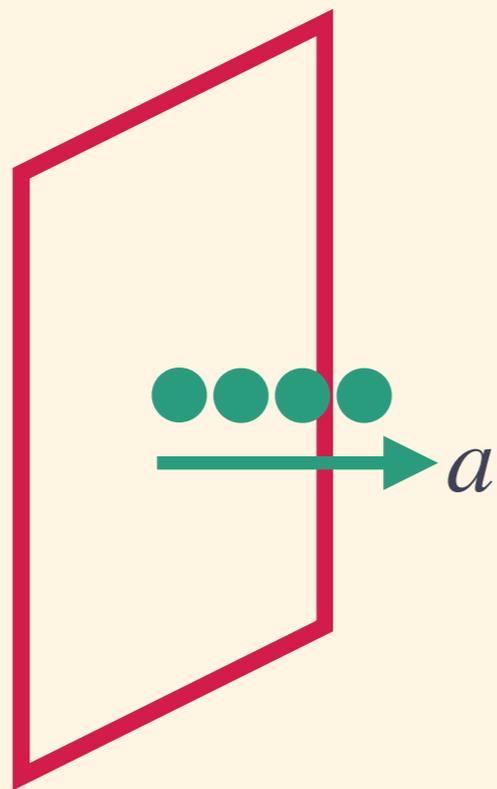


$\Rightarrow A_{Q-1}$ -type $(2,0)$ theory
+ center of mass mode

E-string theory

Q M5-branes on

“End-of-the-world” brane (10d)



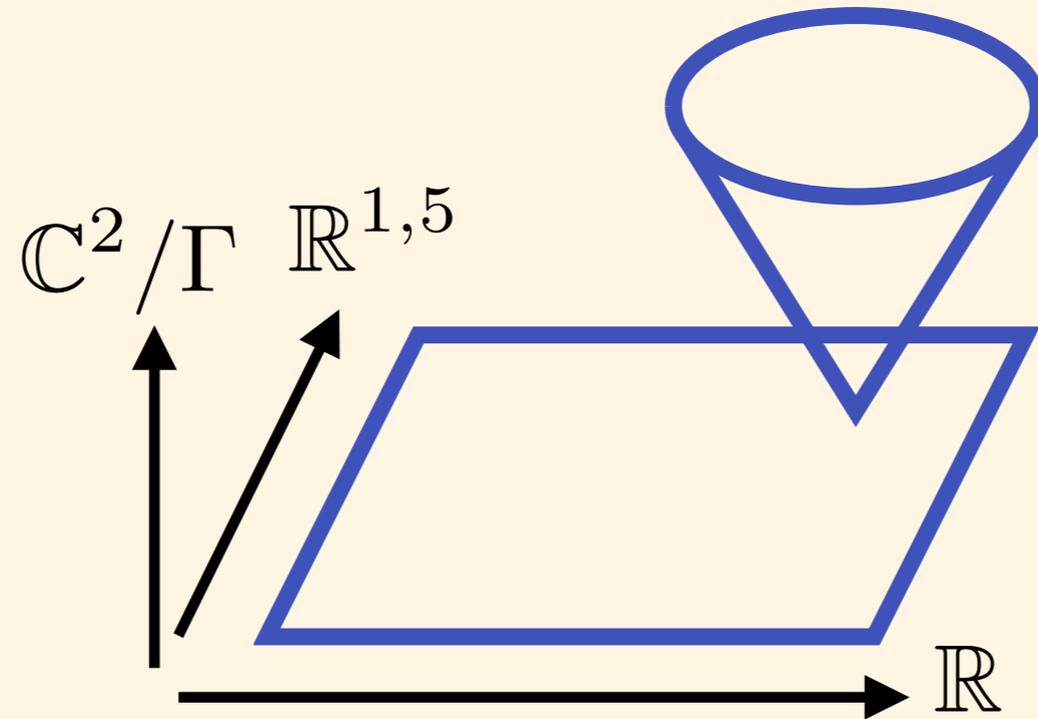
\Rightarrow E-string theory of rank Q
with E_8 flavor symmetry
+ free hyper
(center of mass mode)

M5-branes on $\mathbb{C}^2/\Gamma_{A,D,E}$

$\Gamma_G \subset SU(2)$: Finite Subgroup

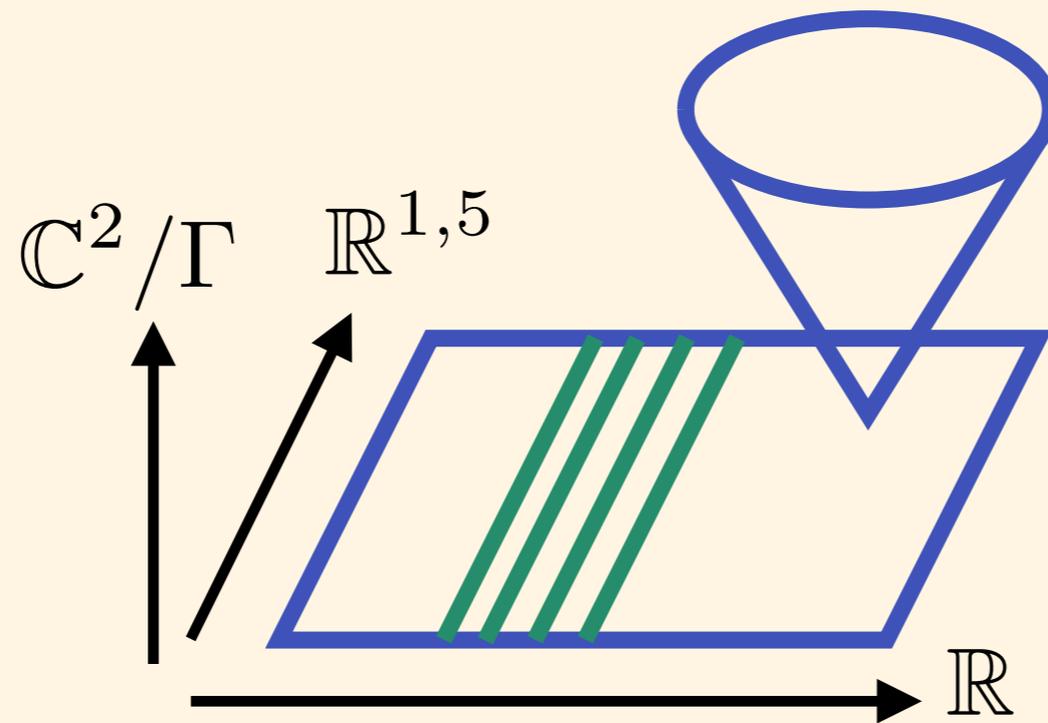
\Rightarrow 7d Vector mult. with gauge group G

On singular locus of $\mathbb{C}^2/\Gamma \times \mathbb{R}^{1,6}$



M5-branes on $\mathbb{C}^2/\Gamma_{A,D,E}$

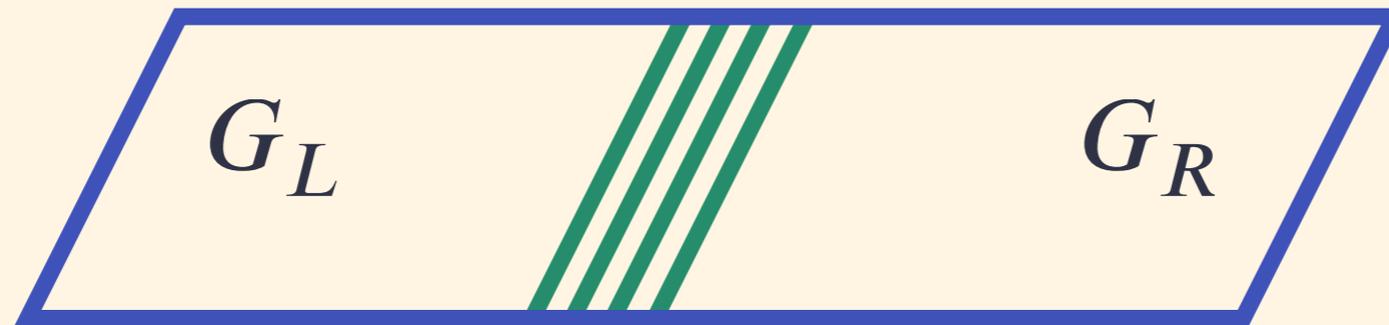
Q M5-branes on
Singular locus of $\mathbb{C}^2/\Gamma \times \mathbb{R}^{1,6}$



$G \times G$ flavor symmetry

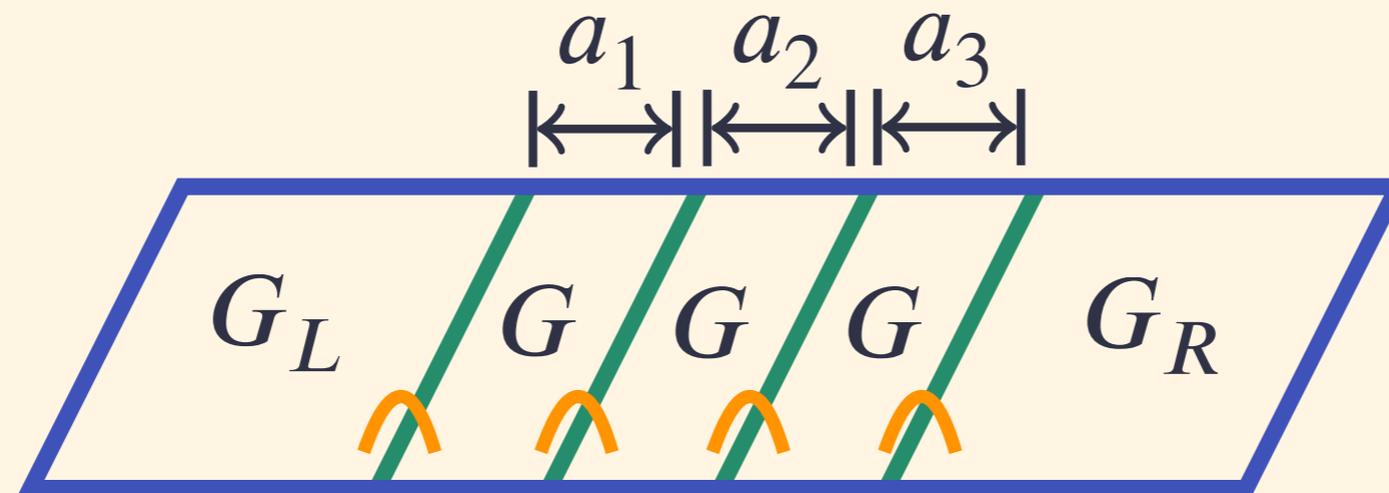
Tensor branch theory

A_k -type singularity case



Tensor branch theory

A_k -type singularity case



$$G = SU(k + 1) + \text{bifundamentals}$$

$Q-1$ dynamical vector mult.s of $SU(k + 1)$

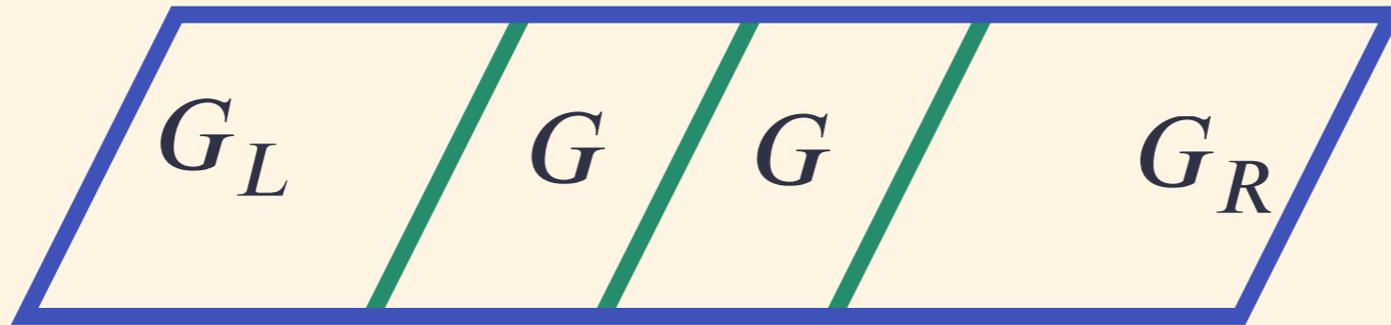
+ bifundamental hypers of neighboring $SU(k+1)$'s

+ $Q-1$ Tensor multiplets with tensor vev.s a_i (dynamical)

+ massive string (M2 branes)

Tensor branch theory

D_k -type singularity case

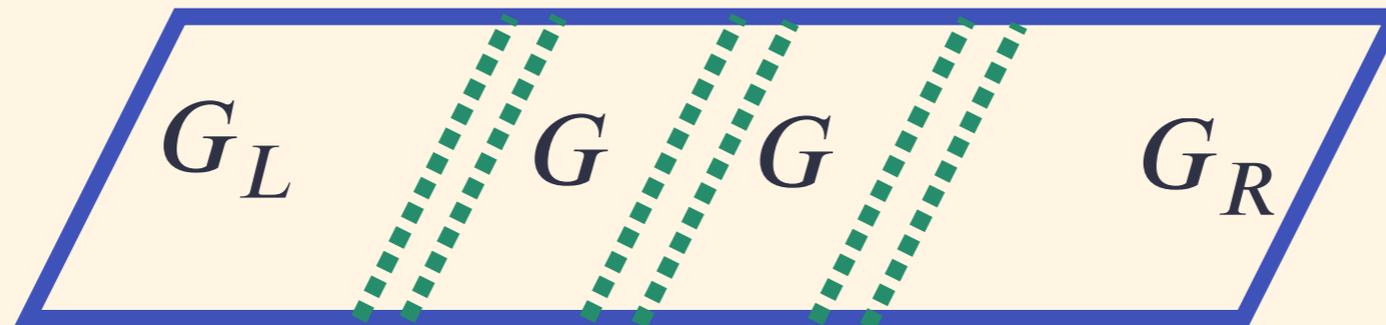


$$G = SO(2k)$$

Tensor branch theory

D_k -type singularity case

“Fractional M5”



$$G = SO(2k)$$

Tensor branch theory

D_k -type singularity case

“Fractional M5”



$$G = SO(2k) \quad +\text{half bifundamentals}$$
$$G' = USp(2k - 8)$$

$SO(2k)$ - $USp(2k-8)$ alternating quiver theory

Gauge couplings are governed by tensor vev.s (dynamical)

This system is also realized by D6-O6-NS5 system in IIA

Tensor branch theory

D_k -type singularity case

$$k = 4$$

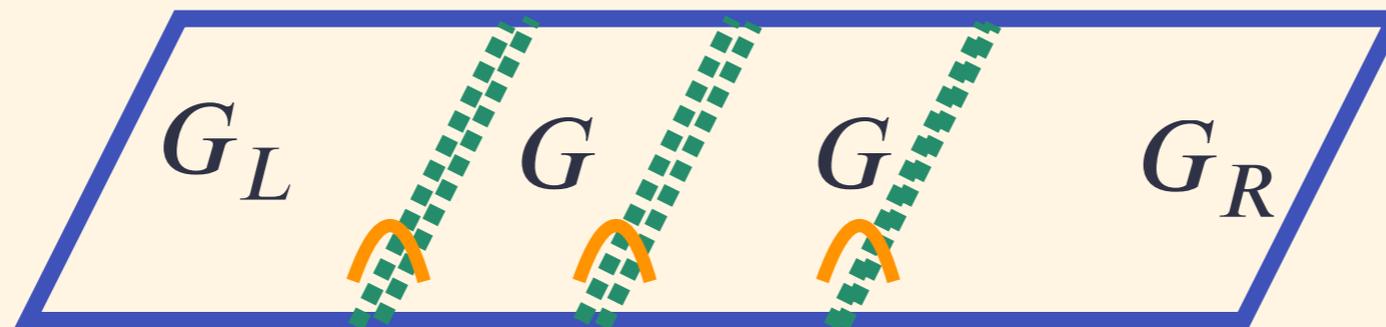


$$G = SO(8)$$
$$G' = USp(0)$$

Tensor branch theory

D_k -type singularity case

$$k = 4$$



$G = SO(8)$ +rank 1 E-string theories

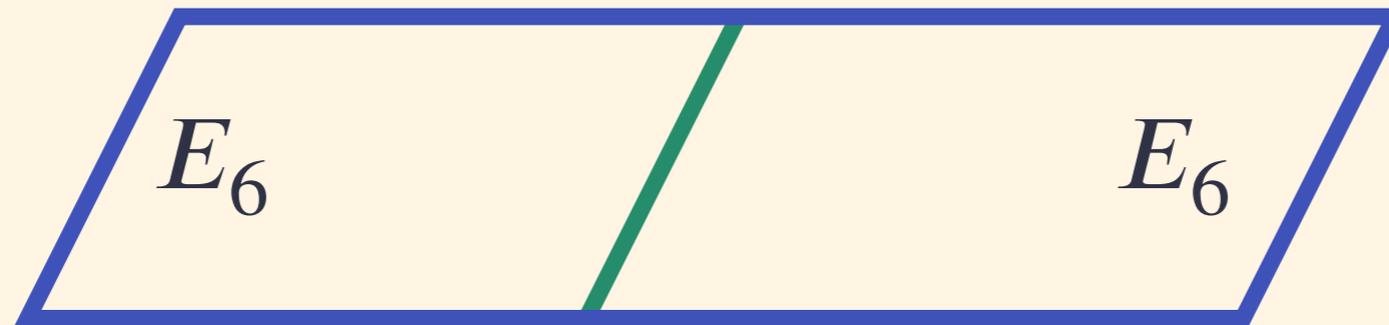
“(SO(8), SO(8)) conformal matter” = E-string

$$SO(8) \times SO(8) \subset E_8$$

[Del-Zotto, Heckman, Tomasiello, Vafa '14]

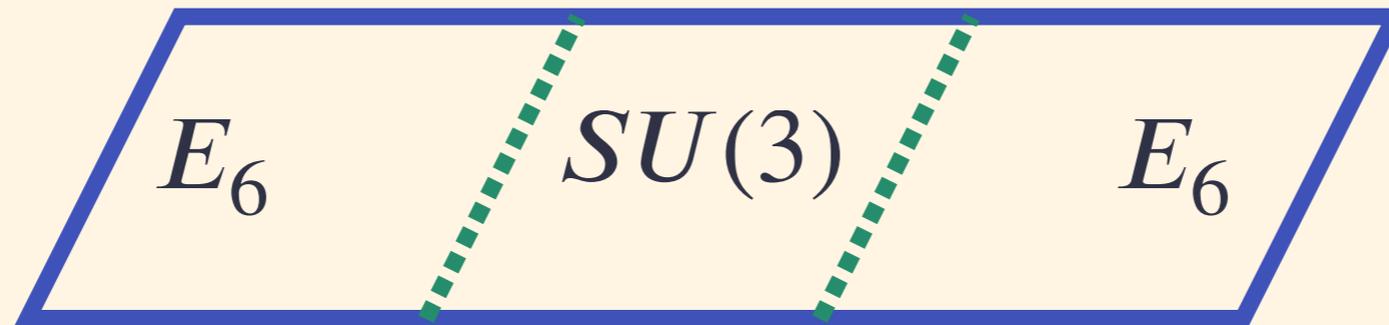
Tensor branch theory

E_6 -type singularity case



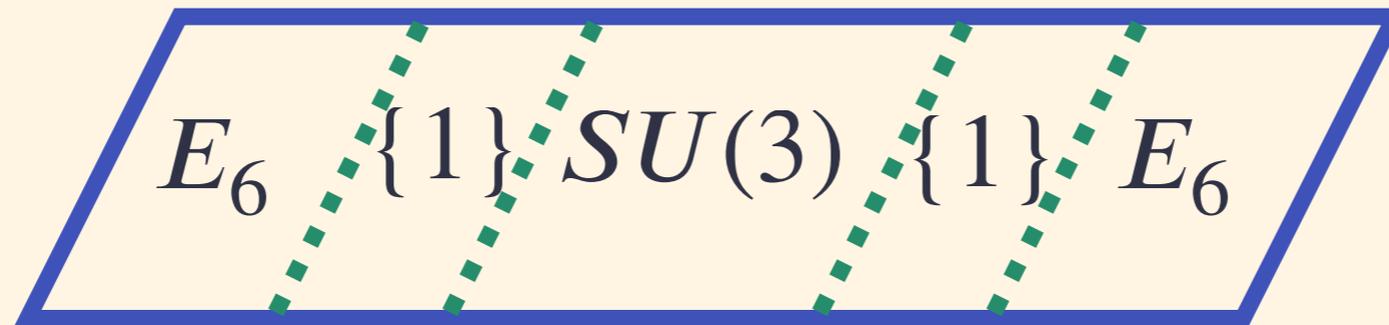
Tensor branch theory

E_6 -type singularity case



Tensor branch theory

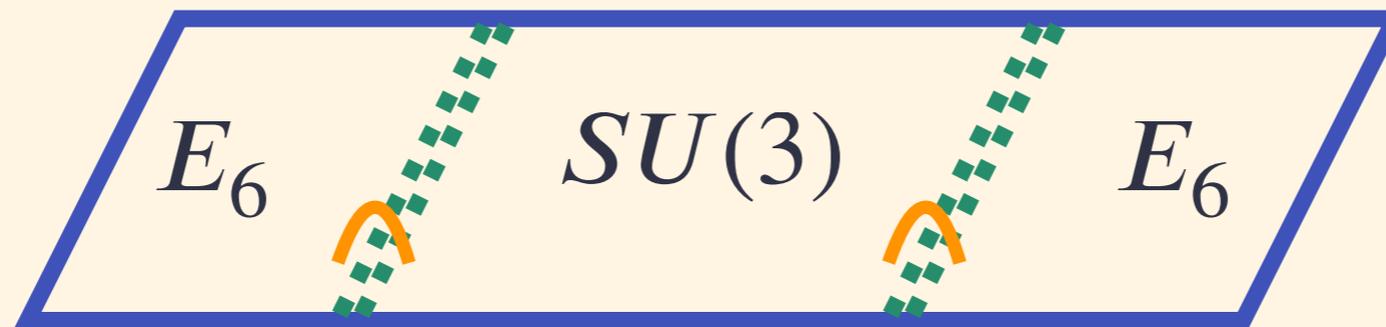
E_6 -type singularity case



Tensor branch theory

E_6 -type singularity case

[Del-Zotto, Heckman, Tomasiello, Vafa '14]



rank 1 E-string theories

“(SU(3), E_6) conformal matter” = E-string

$$SU(3) \times E_6 \subset E_8$$

Other 6d SCFTs

- End-of-the-world brane + ALE singularity
+ (multiple) M5 branes
[Del-Zotto, Heckman, Tomasiello, Vafa '14]
- F-theory construction
[Heckman, Morrison, Vafa '13]
- Lagrangian constructions
[Smilga '07], [Samtleben, E. Sezgin, and R. Wimmer '11, '12]
[Ho, Matsuo '14] for (2,0) theories

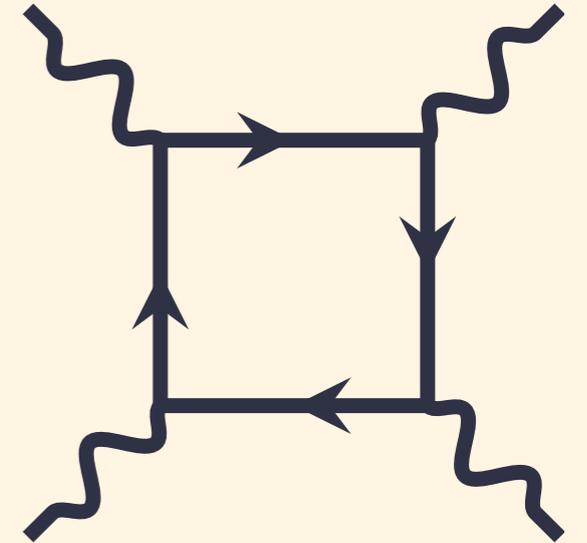
Anomaly polynomials

Anomaly in 6d SCFTs

- Anomaly polynomial 8-form I_8

$$\begin{aligned} I_8 \supset & (\text{Tr } F_R^2)^2, & & SU(2)_R^4 \\ & (\text{Tr } R^2)^2, (\text{Tr } R^4), & & \text{grav.}^4 \\ & (\text{Tr } F_G^2)^2, & & G^4 \\ & \text{Tr } F_R^2 \text{Tr } R^2, \text{Tr } F_R^2 \text{Tr } F_G^2, \dots & & \text{mixed} \end{aligned}$$

- Anomaly polynomial should exist even for non-perturbative SCFTs.



6d Green-Schwarz term

- $\int B \wedge X_4, X_4 \supset \text{Tr} F_R^2, \text{Tr} F_G^2, \text{Tr} R^2$

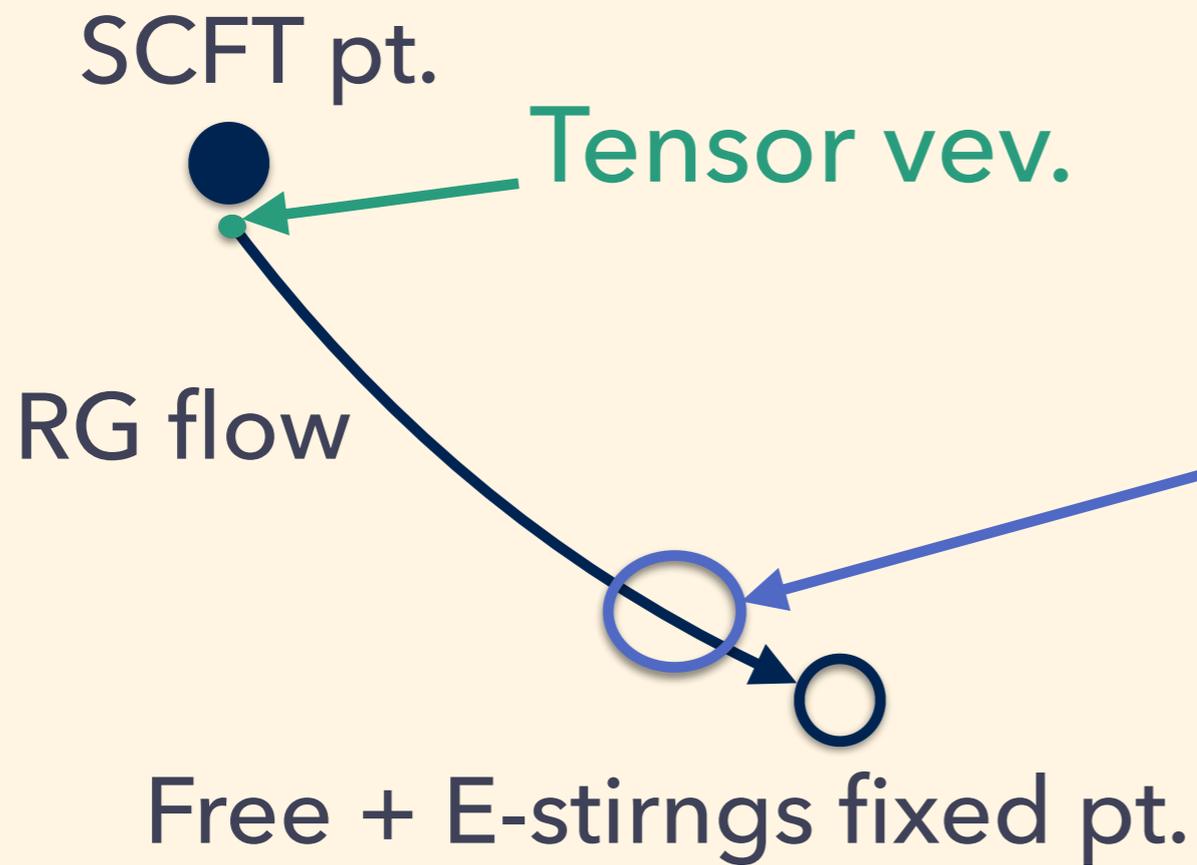
induces modified Bianchi identity $dH \propto X_4$
and additional anomaly $\delta I_8 = \frac{1}{2} X_4^2$ [Monnier '13]

- This contribution is important for calculating anomaly polynomial of 6d SCFTs.

Tensor branch anomaly matching

Basic idea

[Intrilligator '14],[KO,Shimizu,Tachikawa,Yonekura '14]



vector multiplets,
"bifundamentals",
(= hyper or E-string)
tensor multiplets
+ corrections
from massive string

GS terms $\int B \wedge X$

$$I^{UV} = I^{\text{Naive}} + I^{\text{GS}} \quad \leftarrow \quad I^{\text{GS}} = \frac{1}{2} X^2$$

Tensor branch anomaly matching

- Vector, tensor, hyper mult. and E-strings generate $I^{\text{Naive}} \supset (\text{Tr } F_{\text{gauge}}^2)^2, \text{Tr } F_{\text{gauge}}^2 \text{Tr } R^2, \text{Tr } F_{\text{gauge}}^2 \text{Tr } F_R^2, \text{Tr } F_{\text{gauge}}^2 \text{Tr } F_{\text{flavor}}^2$
- Anomaly for E-strings are calculated from anomaly inflow: [\[KO, Shimizu, Tachikawa '14\]](#)
- Because the UV theory are superconformal, $I^{UV} = I^{\text{Naive}} + I^{\text{GS}}$ should not contain F_{gauge}

Tensor branch anomaly matching

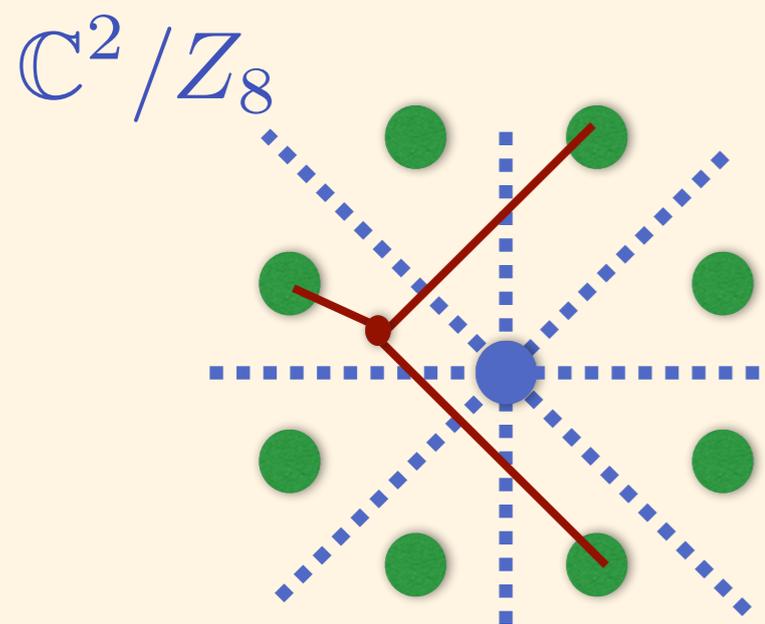
- Example: $I^{\text{Naive}} \supset -\alpha(\text{Tr} F_{\text{gauge}}^2)^2 - \beta \text{Tr} F_{\text{gauge}}^2 \text{Tr} R^2$
- To cancel this by $I^{\text{GS}} = \frac{1}{2} X^2$
 X should contain
$$X \supset \sqrt{2\alpha} \text{Tr} F_{\text{gauge}} + \frac{\beta}{\sqrt{2\alpha}} \text{Tr} R^2$$
- We can fix all of the GS terms $\int B_i \wedge X^i$
in this manner. (If # of $X^i = \#$ of gauge fields)
- $I^{UV} = I^{\text{Naive}} + I^{\text{GS}}$: Do square completion!

Tensor branch anomaly matching

- Result for Q M5's on Singularities:

$$I^{\text{UV}} \supset \frac{Q^3 |\Gamma|^2}{96} (\text{Tr} F_R^2)^2$$

- The leading behavior can be understood from Higgs branch and **M2 junction**:



Q^3 behavior \leftarrow **M2 junction**.
 $Q^3 |\Gamma|^3$ ways of suspending junction
 divided by Γ

M5 branes and mirrors

Torus compactifications

(on going with the same collaborators)

Torus compactifications

- What is the 4d theory of M5 branes on torus-compactified singularities?

	$\mathbb{R}^{1,3}$	T^2	\mathbb{R}	\mathbb{C}^2/Γ
M5	•	•		

- Is that SCFT?
- Can we find candidates in Class S theories?
- How to check?

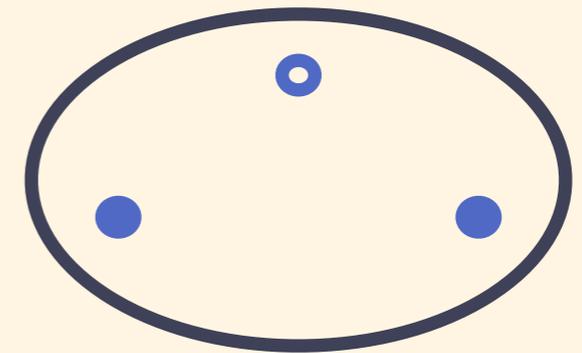
Known examples and Class S

- Single M5 on $\Gamma = A_k$ singularity
⇒ bifundamental both in 6d and 4d.
- Single M5 on $\Gamma = D_4$ singularity
⇒ Rank 1 E-string theory in 6d
⇒ Minahan-Nemeshansky E_8 theory in 4d

Known examples and Class S

- Single M5 on $\Gamma = A_k$ singularity
 \Rightarrow bifundamental both in 4d.

= Class S of type A_k def'ed by

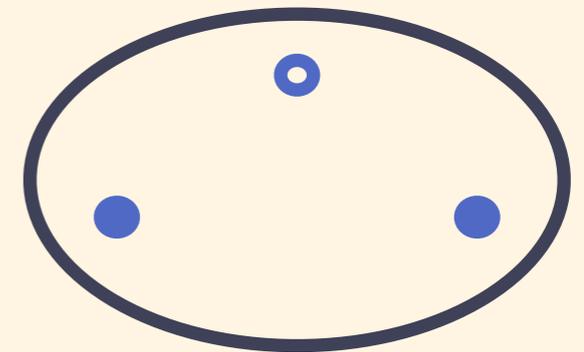


- : “full” puncture (full symmetry)
- : “simple” puncture ($U(1)$ or no sym.)

Known examples and Class S

- Single M5 on $\Gamma = D_4$ singularity
⇒ Rank 1 E-string theory in 6d
⇒ Minahan-Nemeshansky E_8 theory in 4d

= Class S of type D_4 def'ed by

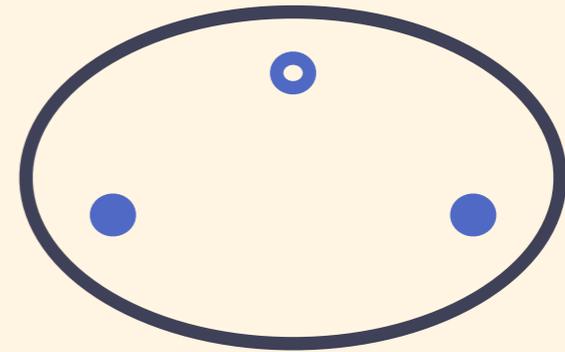


- : “full” puncture (full symmetry)
- : “simple” puncture ($U(1)$ or no sym.)

Guess for $Q=1$

Single M5 on \mathfrak{g} -sing. $\times T^2$ without center of mass
(`(G,G) conformal matter`)

?
= Class S of type \mathfrak{g} def'ed by



- : “full” puncture (full symmetry)
- : “simple” puncture ($U(1)$ or no sym.)

“4d conformal matter”

Checks

- Dim.s of Higgs/Coulomb branches matches
- The geometry of Higgs branch = $\mathbb{C}^2/\Gamma_{\mathfrak{g}}$
- With plausible assumption, we can calculate the conformal anomalies and flavor levels for one M5 on \mathfrak{g} -sing. $\times T^2$
 \Rightarrow matches with the class S conjecture.
- M/IIB duality chain? Not so clear.

Conclusion

- The anomaly polynomials of 6d SCFTs can be calculated by adding tensor branch contributions and GS contributions.
- For “6d (G,G) conformal matter”, the candidate for the torus compactified theories can be found in the Class S theories.

Further directions

- Comactifications of 6d (1,0) SCFTs
 - 4d anomaly polynomials for T^2 compactified theories?
 - What and how the 4d theories are?
- 6d a-theorems?

Appendix

Tensor branch

anomaly matching wo/ vector

- $N=(2,0)$ and E-string theories :
there is no vector in tensor branch theory
- We know 5d theories obtained by circle compactifications for those theories.
- In 5d coulomb branch corresponds to 6d tensor branch, we have $U(1)$ vector A_i come from 6d self-dual tensor B_i and massive mode because of coulomb vev.

Tensor branch

anomaly matching wo/ vector

- In 5d coulomb branch corresponds to 6d tensor branch, we have U(1) vector A_i come from 6d self-dual tensor B_i and massive mode because of coulomb vev.
- Those massive modes generates 5d CS terms $\int A_i \wedge X^i$, which is calculable.
- This should come from 6d GS term $\int B_i \wedge X^i$