

Information metric for the matrix geometry

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Introduction

- Noncommutative geometry = Quantum structure of space-time ??

[Madore 1991, Connes 1994, Hanada-Kawai-Kimura 2006, Steinacker 2010, etc.]

- Noncommutative coordinates \rightarrow loss of “locality”

$$[x^\mu, x^\nu] \neq 0 \quad \longrightarrow \quad \Delta x^\mu \Delta x^\nu \gtrsim \Lambda$$

- UV divergences \rightarrow Cutoff by Λ

Matrix geometry

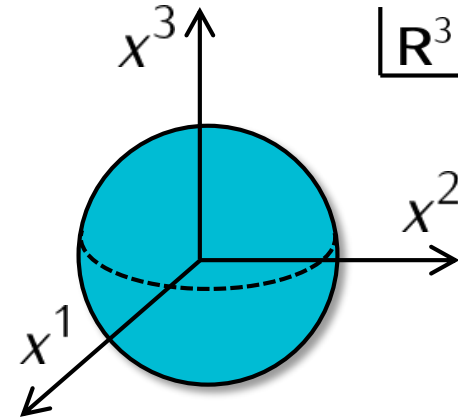
- A typical example of NC geometry
 - Coordinates \rightarrow Hermitian matrices of finite size
 - Many examples: fuzzy sphere, fuzzy torus, fuzzy CP^N , etc.

Example: fuzzy two sphere

[Madore 1992]

– (Unit) two sphere $S^2 \subset \mathbb{R}^3$

$$\left\{ \begin{array}{l} (x^1, x^2, x^3) \in \mathbb{R}^3 : \text{embedding functions} \\ x^i x_i = 1 \\ [x^i, x^j] = 0 \end{array} \right.$$



– Fuzzy two sphere

$$\left\{ \begin{array}{l} X^i = \frac{2}{\sqrt{N^2-1}} L^i \quad L^i: N \text{ dim irrep of SU(2) generators} \\ X^i X_i = 1 \\ [X^i, X^j] = \frac{2}{\sqrt{N^2-1}} i \epsilon^{ijk} X^k \longrightarrow 0 \quad (N \rightarrow \infty) \end{array} \right.$$

Commutative limit

Functions on two sphere

Two sphere

- Functions

$$f(\Omega) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} c_{\ell m} Y_{\ell m}(\Omega)$$

- Spherical harmonics

$$Y_{\ell m}(\Omega) = \sum_{k=0}^{\ell} t_{i_1 i_2 \dots i_k}^{\ell m} X^{i_1} \dots X^{i_k}$$

- Algebra

$$Y_{\ell_1 m_1}(\Omega) Y_{\ell_2 m_2}(\Omega) = \sum_{\ell_3 m_3} C_{\ell_1 m_1 \ell_2 m_2}^{\ell_3 m_3} Y_{\ell_3 m_3}(\Omega)$$

$C_{\ell_1 m_1 \ell_2 m_2}^{\ell_3 m_3}$: Clebsch-Gordan coefficients

Fuzzy two sphere

- Matrices UV cutoff

$$\hat{f} = \sum_{\ell=0}^{N-1} \sum_{m=-\ell}^{\ell} c_{\ell m} \hat{Y}_{\ell m}$$

- Fuzzy spherical harmonics

$$\hat{Y}_{\ell m} = \sum_{k=0}^{\ell} t_{i_1 i_2 \dots i_k}^{\ell m} X^{i_1} \dots X^{i_k}$$

- Algebra

$$\hat{Y}_{\ell_1 m_1} \hat{Y}_{\ell_2 m_2} = \sum_{\ell_3 m_3} \hat{C}_{\ell_1 m_1 \ell_2 m_2}^{\ell_3 m_3} \hat{Y}_{\ell_3 m_3}$$

$$\hat{C}_{\ell_1 m_1 \ell_2 m_2}^{\ell_3 m_3} \xrightarrow{N \rightarrow \infty} C_{\ell_1 m_1 \ell_2 m_2}^{\ell_3 m_3}$$

Differential & integral

Two sphere

- Poisson bracket

$$\{x^i, x^j\} = 2\epsilon^{ijk} x_k$$

- Integral

$$\int \frac{d\Omega}{2\pi} Y_{\ell m}(\Omega)$$

Fuzzy two sphere

- Commutator

$$-iN[X^i, X^j] \sim 2\epsilon^{ijk} X_k$$

- Commutator

$$\frac{1}{N} \text{Tr} \hat{Y}_{\ell m} \xrightarrow{N \rightarrow \infty} \int \frac{d\Omega}{2\pi} Y_{\ell m}(\Omega)$$

– Poisson bracket gives angular momentum operator

$$\{x^i, \cdot\} = 2\epsilon^{ijk} x_k \partial_j$$

Dictionary of fuzzy two sphere

- Fuzzy sphere is given by a mapping

$$x^i \mapsto X^i = \frac{2}{\sqrt{N^2 - 1}} L^i$$

- Dictionary of functions & differential/integral operators

$$Y_{\ell m}(\Omega) \longleftrightarrow \hat{Y}_{\ell m}$$

$$\{x^i, \cdot\} \longleftrightarrow -iN[X^i, \cdot]$$

$$\int \frac{d\Omega}{2\pi} \longleftrightarrow \frac{1}{N} \text{Tr}$$

- This mapping is an example of **matrix regularization**

Matrix regularization (MR)

– MR is a sequence of linear maps $\{T_N\}$ ($N = 1, 2, \dots, \infty$) which satisfy

[Goldstone-Hoppe, 1982; Arnlind-Hoppe-Huisken, 2012]

$$\left\{ \begin{array}{l} T_N : C^\infty(\mathcal{M}) \rightarrow N \times N \text{ Hermitian matrices} \\ \lim_{N \rightarrow \infty} \|T_N(fg) - T_N(f)T_N(g)\| = 0 \\ \lim_{N \rightarrow \infty} \|iN[T_N(f), T_N(g)] - T_N(\{f, g\})\| = 0 \\ \lim_{N \rightarrow \infty} \frac{2\pi}{N} \text{Tr } T_N(f) = \int_{\mathcal{M}} f \omega \end{array} \right.$$

\mathcal{M} : symplectic manifold ω : symplectic form on \mathcal{M}

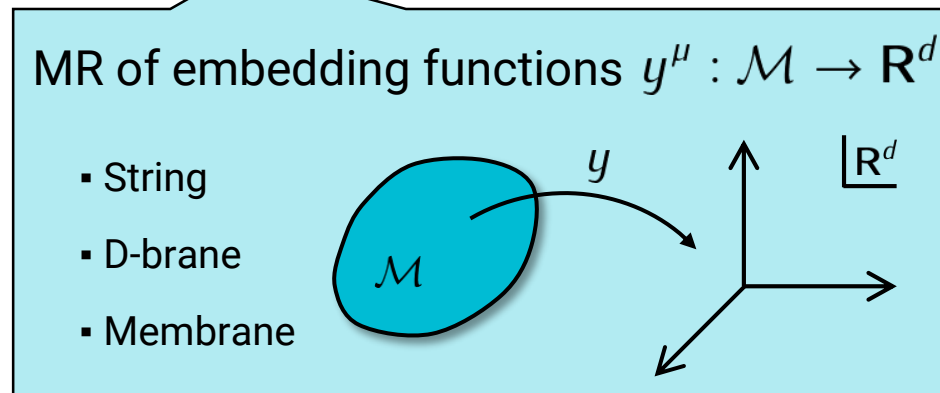
$\| \cdot \|$: norm $\{ , \}$: Poisson bracket on \mathcal{M}

Relation to string/M theories

– Matrix models ← MR of string/M theories

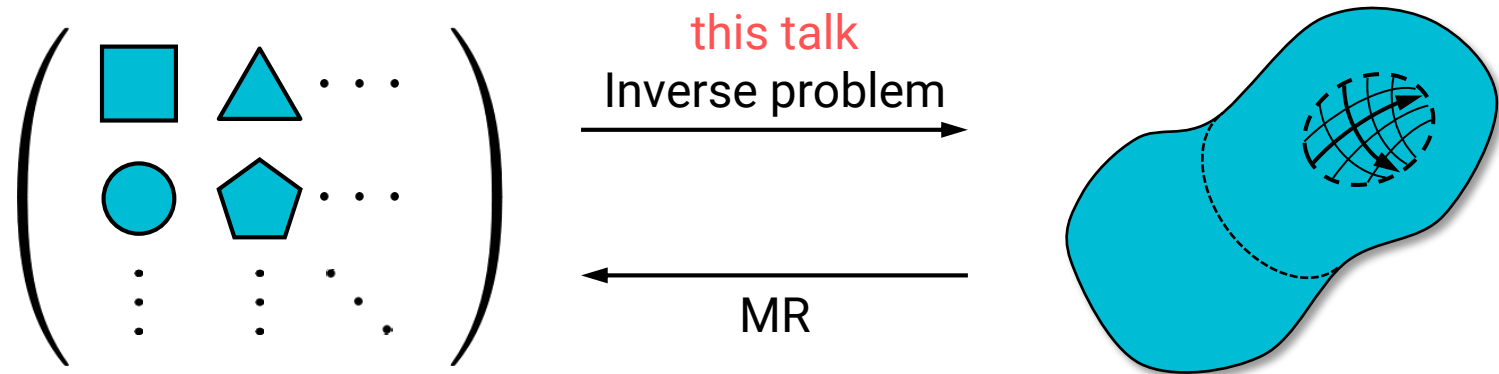
[BFSS 1997, IKKT 1997, DVV 1997]

$$Z = \int \prod_{\mu=1}^d DX^\mu e^{-S} \quad X^\mu : N \times N \text{ Hermitian matrices}$$



– Non-perturbative formulation & second quantization

Matrices & Smooth geometry



– Smooth geometry \rightarrow Matrices: MR

[de Wit-Hoppe-Nicolai 1988, Bordemann-Meinrenken-Schlichenmaier 1994]

– Matrices \rightarrow Smooth geometry: inverse problem of MR \leftarrow this talk

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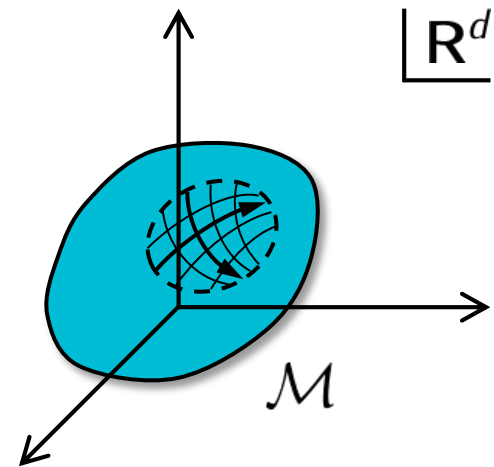
2. Inverse problem of MR

Inverse problem of MR

– For given matrices, how can we associate a smooth geometry \mathcal{M} ??

$N \times N$ Hermitian matrices
 X^μ ($\mu = 1, 2, \dots, d$)

How ??
 $N \rightarrow \infty$



– Many methods

- using Dirac operator, [Asakawa-Sugimoto-Terasima 2002, Berenstein-Dzienkowski 2012]
- using coherent states, [Ishiki 2015, Schneiderbauer-Steinacker 2016] ← **this talk**
- etc., [Hotta-Nishimura-Tsuchiya 1999, Azeyanagi-Hanada-Hirata-Shimada 2009]

Commutative space for X^μ

[Ishiki 2015]

– d $N \times N$ Hermitian matrices

$$\{(X^1, X^2, \dots, X^d) \mid N = 1, 2, \dots\}$$

– “Hamiltonian” using given matrices

$$H(y) = \frac{1}{2} \sum_{\mu=1}^d (X^\mu - y^\mu)^2 \quad y^\mu \in \mathbf{R}$$

– Eigenstates & eigenvalues

$$H(y)|n, y\rangle = E_n(y)|n, y\rangle \quad (n = 0, 1, \dots, N - 1)$$

$$E_0(y) \leq E_1(y) \leq \dots$$

Commutative space for X^μ

[Ishiki 2015]

– Zeros of $E_0(y)$ = points on commutative space

$$\begin{aligned} E_0(y) &= \langle H(y) \rangle = \frac{1}{2} \langle (X^\mu - y^\mu)^2 \rangle = \frac{1}{2} \langle (X^\mu)^2 \rangle - \langle X^\mu \rangle y_\mu + \frac{1}{2} (y^\mu)^2 \\ &= \frac{1}{2} \langle (X^\mu)^2 \rangle - \frac{1}{2} \langle X^\mu \rangle^2 + \frac{1}{2} \langle X^\mu \rangle^2 - \langle X^\mu \rangle y_\mu + \frac{1}{2} (y^\mu)^2 \\ &= (\Delta X^\mu)^2 + (\langle X^\mu \rangle - y^\mu)^2 \end{aligned}$$

$$E_0(y) \rightarrow 0 \quad \Leftrightarrow \quad \begin{cases} \langle 0, y | X^\mu | 0, y \rangle \rightarrow y^\mu \\ \Delta X^\mu \rightarrow 0 \end{cases}$$

The wave packet shrinks to a point at $y \in \mathbb{R}^d$

Commutative space for X^μ

[Ishiki 2015]

– Commutative space for given matrices = zeros of $E_0(y)$

$$\mathcal{M} = \left\{ y \in \mathbf{R}^d \mid \lim_{N \rightarrow \infty} E_0(y) = 0 \right\}$$

There are no zero points for finite N except trivial cases

– This works for fuzzy sphere, fuzzy torus, etc.

Construction method

[Ishiki 2015]

Given matrices

$$N \times N \text{ Hermitian matrices} \\ X^\mu \quad (\mu = 1, 2, \dots, d)$$

Commutative space

$$\text{Set of zeros of } E_0(y) \\ \mathcal{M} = \left\{ y \in \mathbb{R}^d \mid \lim_{N \rightarrow \infty} E_0(y) = 0 \right\}$$

Hamiltonian

$$H(y) = \frac{1}{2} (X^\mu - y^\mu)^2 \\ y^\mu \in \mathbb{R}$$

Eigenstates & energies

$$H(y)|n, y\rangle = E_n(y)|n, y\rangle \\ E_0(y) \leq E_1(y) \leq \dots$$

Example: fuzzy two sphere

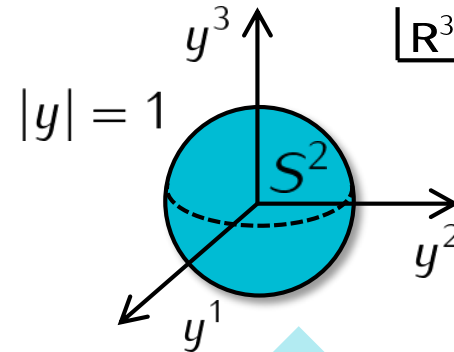
[Ishiki 2015]

Given matrices

$$X^i = \frac{2}{\sqrt{N^2-1}} L^i \quad (i = 1, 2, 3)$$

L^i : N dim irrep of $SU(2)$ generators

Commutative space



Hamiltonian

$$H(y) = \frac{1 + |y|^2}{2} - \frac{2y_i L^i}{\sqrt{N^2-1}}$$

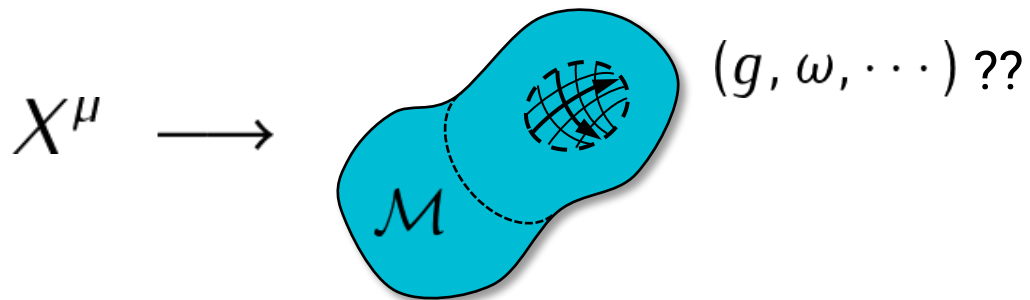
Eigenstates & energies

$$E_0(y) = \frac{1 + |y|^2}{2} - \frac{N|y|}{\sqrt{N^2-1}}$$
$$\xrightarrow{N \rightarrow \infty} \frac{1}{2} \underline{(1 - |y|)^2}$$

3. Information metric

Geometric structures on \mathcal{M}

- In general, smooth space has geometric structures on it
- Do X^μ have information of those geometric structures ??



- Information metric \rightarrow Riemannian structure on \mathcal{M}

Density matrix for X^μ

– In general, $E_0(y)$ is “degenerate” (degree of degeneracy $\equiv k$)

– There are k linearly independent the ground states

$$|0, y\rangle_\alpha \quad \alpha = 1, 2, \dots, k$$

– The ground states $\rightarrow N \times N$ density matrix

$$\rho(y) = \frac{1}{k} \sum_{\alpha=1}^k |0, y\rangle_\alpha \langle 0, y|_\alpha \quad (\text{weight} = \frac{1}{k})$$

– This invariant under changing basis

$$|0, y\rangle_\alpha \rightarrow c^\beta_\alpha(y) |0, y\rangle_\beta \quad (c^\dagger c = 1)$$

Berry phase

[Ishiki-TM-Muraki 2016]

- **Berry phase** (in the case for degenerate)

$$A_{\alpha\beta} = -i_{\alpha} \langle 0, y | \frac{\partial}{\partial y} | 0, y \rangle_{\beta}$$

- Under changing basis, it transform as a gauge field

$$\left\{ \begin{array}{l} |0, y\rangle_{\alpha} \rightarrow c^{\beta}_{\alpha}(y) |0, y\rangle_{\beta} \quad (c^{\dagger} c = 1) \\ A_{\alpha\beta}(y) \rightarrow (c^{\dagger}(y) A(y) c(y))_{\alpha\beta} - i(c^{\dagger}(y) \partial c(y))_{\alpha\beta} \end{array} \right.$$

Information metric

- Information metric for density matrices ρ

$$\begin{cases} ds^2 = \frac{1}{2} \text{Tr}[G d\rho] \\ d\rho = G\rho + \rho G \end{cases}$$

- This measure the distance between density matrices

- In the case for pure state: $\rho = |\psi\rangle\langle\psi|$

$$ds^2 = \langle d\psi|d\psi\rangle - \langle d\psi|\psi\rangle\langle\psi|d\psi\rangle$$

Riemannian metric via pullback

- The density matrix using the ground states is a mapping

$$\left\{ \begin{array}{l} \rho : y \in \mathcal{M} \mapsto \rho(y) \in \{ \text{Set of all } N \times N \text{ density matrices} \} \\ \rho(y) = \frac{1}{k} \sum_{\alpha=1}^k |0, y\rangle_{\alpha} \langle 0, y| \end{array} \right.$$

- In most cases, it is also an “embedding”
- $\rho(y)$ induces a Riemannian metric on \mathcal{M} via the pullback

$$ds^2 = \frac{1}{2} \text{Tr}[G(y) d\rho(y)]$$

Example 1: fuzzy two sphere

– Fuzzy two sphere: MR of two sphere $S^2 \subset \mathbb{R}^3$

$$\left\{ \begin{array}{l} X^i = \frac{2}{\sqrt{N^2-1}} L^i \quad (i = 1, 2, 3) \\ L^i: N \text{ dim irrep of } \text{SU}(2) \text{ generators} \end{array} \right.$$

– Hamiltonian & the ground states

$$\left\{ \begin{array}{l} H(y) = U(\theta, \phi) \left(\frac{1+|y|^2}{2} - \frac{2|y|L^3}{\sqrt{N^2-1}} \right) U^\dagger(\theta, \phi) \\ U(\theta, \phi) = e^{-i\phi L_3} e^{-i\theta L_2} e^{i\phi L_3} \\ |\Omega\rangle = U(\theta, \phi) |JJ\rangle \quad (2J + 1 = N) \end{array} \right.$$

Bloch coherent states

Example 1: fuzzy two sphere

– Density matrix & information metric

$$\left\{ \begin{array}{l} \rho = |\Omega\rangle\langle\Omega| \\ ds^2 = (d\theta)^2 + \sin^2 \theta (d\phi)^2 \end{array} \right.$$

Riemannian metric on two sphere !!

Example 2: fuzzy “four” sphere

– Fuzzy four sphere: fuzzy surface for $S^4 \subset \mathbb{R}^5$

[Castelino-Lee-Taylor]

$$\left\{ \begin{array}{l} \overbrace{X^\mu = \frac{1}{N} (\Gamma^\mu \otimes 1_4 \otimes \cdots \otimes 1_4 + \cdots + 1_4 \otimes \cdots \otimes 1_4 \otimes \Gamma^\mu)}^{N \text{ tensor product}}_{\text{sym}} \\ \Gamma^\mu (\mu = 1, 2, 3, 4, 5) : \text{five dimensional gamma matrices} \end{array} \right.$$

$$\Gamma^5 = \begin{pmatrix} 1_2 & 0 \\ 0 & -1_2 \end{pmatrix}$$

However, fuzzy four sphere satisfies at least $X^\mu X_\mu = 1 + \mathcal{O}(1/N)$ & $[X^\mu, X^\nu] \rightarrow 0 (N \rightarrow \infty)$

– Four sphere \neq symplectic manifold i.e. fuzzy four sphere \neq MR

Example 2: fuzzy “four” sphere

– Hamiltonian & the ground states for fuzzy four sphere

$$\left\{ \begin{array}{l} H(y) = U \left(\frac{1}{2}(1 + |y|^2) - |y|X^5 \right) U^\dagger + \mathcal{O}(1/N) \\ |0, y\rangle_\alpha = U | \text{eigenstates of } X^5 \text{ with } +1 \rangle_\alpha \quad (\alpha = 1, 2, \dots, N + 1) \\ U = e^{-\chi\Gamma^{21}/2} e^{-\psi\Gamma^{32}/2} e^{-\phi\Gamma^{43}/2} e^{-\theta\Gamma^{54}/2} \end{array} \right.$$

Coherent states are
Degenerate

– Density matrix & information metric

$$\left\{ \begin{array}{l} \rho = \frac{1}{N+1} \sum_\alpha |0, y\rangle_\alpha \langle 0, y| \\ ds^2 = (d\theta)^2 + \sin^2 \theta (d\phi)^2 + \sin^2 \theta \sin^2 \phi (d\psi)^2 + \sin^2 \theta \sin^2 \phi \sin^2 \psi (d\chi)^2 \end{array} \right.$$

Riemannian metric on four sphere !!

4. Summary

Summary

- Hamiltonian & its ground states \rightarrow commutative space \mathcal{M}
- The ground states \rightarrow mapping from \mathcal{M} to set of all density matrices
- The pullback of the information metric \rightarrow Riemannian metric on \mathcal{M}
 - It works for fuzzy four sphere (\neq MR)
- For other examples ?? e.g. fuzzy “three” sphere
- Apply to matrix models ?? e.g. emergent gravity

[Ramgoolam 2000]

[Steinacker 2011]