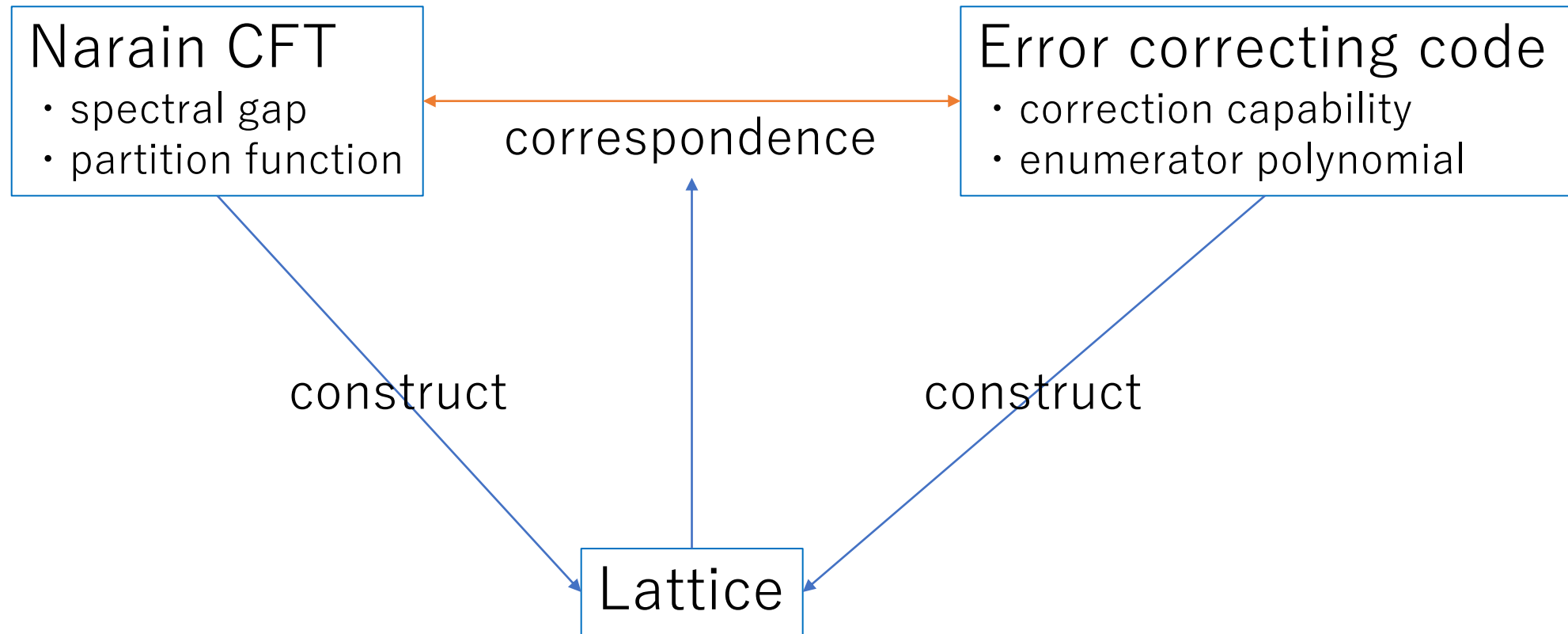


Narain CFTs and error correcting codes

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based on arXiv: 2203.10848

Overview



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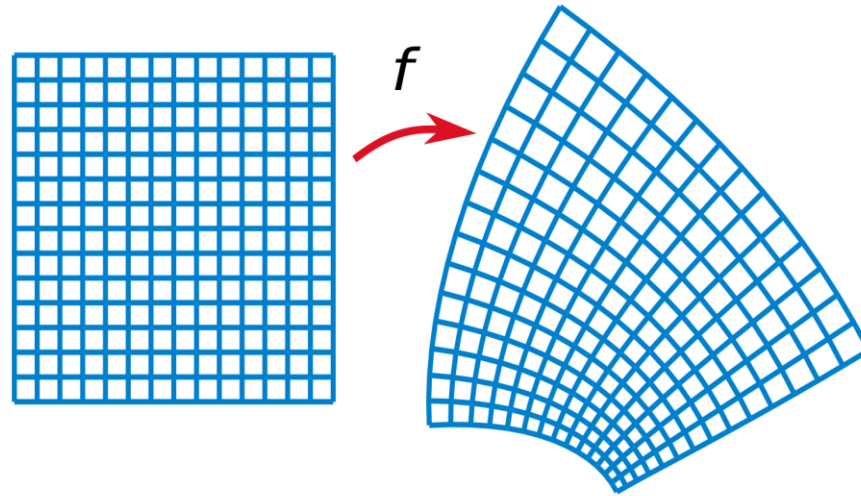
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1. Narain CFT

CFT

- A conformal field theory (CFT) is a quantum field theory that is invariant under conformal transformations.

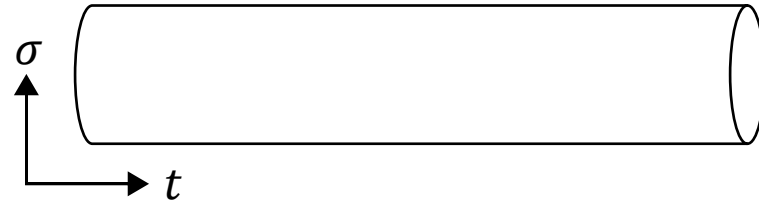
= angle preserving :



- A two-dimensional CFT has rich mathematical structure and is used to describe condensed matter, critical phenomena, and **string theory**.

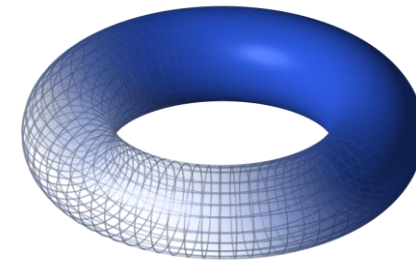
Narain CFT

- We consider “a closed string” : $X(t, \sigma)$, $\sigma \cong \sigma + 2\pi$



- A Narain CFT is a 2d CFT that describes a closed string on the compactified space :

$$X^i \cong X^i + 2\pi R, R : \text{radius}, i = 1, \dots, n$$



$T^{n=2}$

- The action :

$$S = \frac{1}{4\pi\alpha'} \int dt \int_0^{2\pi} d\sigma [G_{ij}(\partial_t X^i \partial_t X^j - \partial_\sigma X^i \partial_\sigma X^j) - 2B_{ij} \partial_t X^i \partial_\sigma X^j]$$

G_{ij} : metric, B_{ij} : antisymmetric background

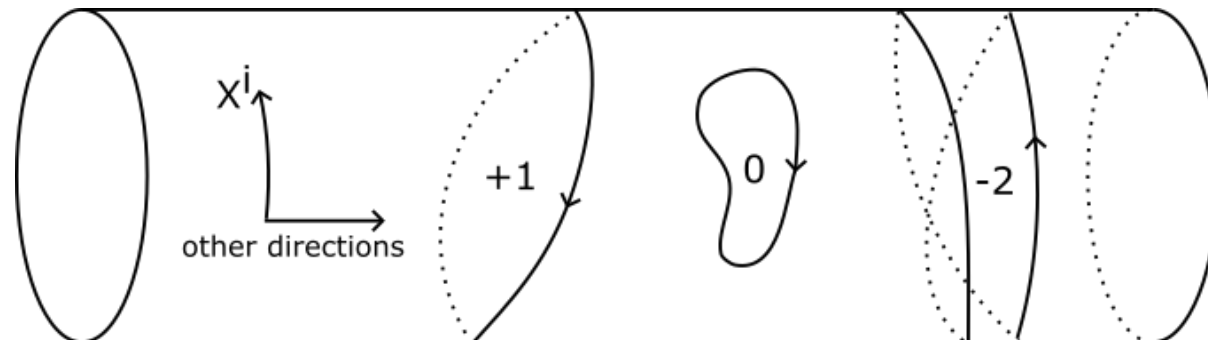
Effects of the compactification

- The center-of-mass momentum P
- The operator $\exp(2\pi i R \hat{P}_i)$, which translates strings once around the i -th direction, must be the identity for states.

$$P_i := \frac{\partial L}{\partial(\partial_t X^i)} = \frac{1}{R} m_i, \quad m_i \in \mathbb{Z} \quad \textcircled{1}$$

- winding number w
- A string can wind around the compact direction.

$$X^i(t, \sigma) - X^i(t, \sigma + 2\pi) = 2\pi R w^i, \quad w^i \in \mathbb{Z} \quad \textcircled{2}$$



Momentum

- From the equation of motion, the mode expansion of X^i is

$$X^i(t, \sigma) = X_L^i(t - \sigma) + X_R^i(t + \sigma),$$

$$X_L^i(t - \sigma) = \hat{x}_L^i + \frac{\alpha'}{2} \hat{p}_L^i(t - \sigma) + i \sqrt{\frac{\alpha'}{2}} \sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{\hat{\alpha}_n^i}{n} e^{-in(t-\sigma)},$$

$$X_R^i(t + \sigma) = \hat{x}_R^i + \frac{\alpha'}{2} \hat{p}_R^i(t + \sigma) + i \sqrt{\frac{\alpha'}{2}} \sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{\hat{\tilde{\alpha}}_n^i}{n} e^{-in(t+\sigma)}$$

- By substituting these for ①②, eigenvalues of \hat{p}_L, \hat{p}_R on orthogonal basis are

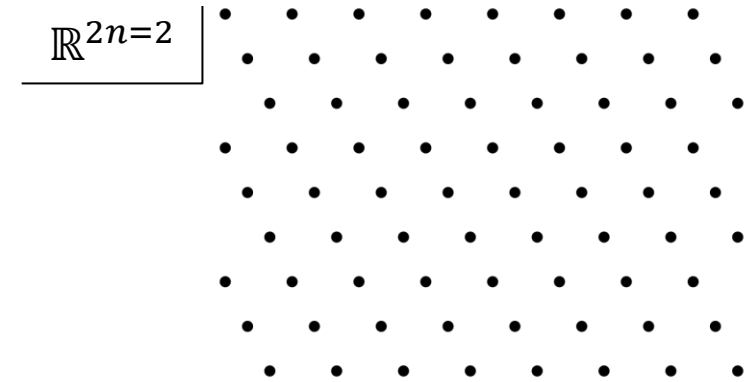
$$k_{L\mu} = e_\mu^i \left[\frac{1}{R} m_i + \frac{R}{2} (B + G)_{ij} w^j \right], \quad k_{R\mu} = e_\mu^i \left[\frac{1}{R} m_i + \frac{R}{2} (B - G)_{ij} w^j \right],$$

$$e_\mu^i : \text{tetrad} \quad (G_{ij} e_\mu^i e_\nu^j = \delta_{\mu\nu})$$

A lattice from a Narain CFT

- The momenta form a lattice :

$$\Lambda(R, G, B) = \left\{ \begin{pmatrix} \vec{k}_L \\ \vec{k}_R \end{pmatrix} \mid \vec{m}, \vec{w} \in \mathbb{Z}^n \right\} \subset \mathbb{R}^{2n}$$



- For later convenience, we define another lattice.

$$\Lambda_N(R, G, B) = \left\{ \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \mid \vec{m}, \vec{w} \in \mathbb{Z}^n \right\} \subset \mathbb{R}^{2n},$$

$$\alpha_\mu = \frac{k_{L\mu} + k_{R\mu}}{\sqrt{2}} = e_\mu^i \left[\frac{\sqrt{2}}{R} m_i + \frac{R}{\sqrt{2}} B_{ij} w^j \right],$$

$$\beta_\mu = \frac{k_{L\mu} - k_{R\mu}}{\sqrt{2}} = e_\mu^i \frac{R}{\sqrt{2}} G_{ij} w^j.$$

← We will associate a code with this lattice

Even self-duality

- **Prop.** The lattice $\Lambda_N(R, G, B)$ is **even** and **self-dual** with a metric

$$g = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}.$$

- Even
 - A lattice Λ is even $:\Leftrightarrow \forall x \in \Lambda, x \cdot x \in 2\mathbb{Z}$
- Self-dual
 - A dual lattice of $\Lambda \subset \mathbb{R}^n$: $\Lambda^* = \{x' \in \mathbb{R}^n \mid \forall x \in \Lambda, x \cdot x' \in \mathbb{Z}\}$
 - A lattice Λ is self-dual $:\Leftrightarrow \Lambda = \Lambda^*$
- We can verify these properties directly.

Proof (even)

- (A lattice Λ is even $:\Leftrightarrow \forall x \in \Lambda, x \cdot x \in 2\mathbb{Z}$)

$$\alpha_\mu = \frac{k_{L\mu} + k_{R\mu}}{\sqrt{2}} = e_\mu^i \left[\frac{\sqrt{2}}{R} m_i + \frac{R}{\sqrt{2}} B_{ij} w^j \right],$$
$$\beta_\mu = \frac{k_{L\mu} - k_{R\mu}}{\sqrt{2}} = e_\mu^i \frac{R}{\sqrt{2}} G_{ij} w^j.$$

- For $\forall x = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \in \Lambda_N(R, G, B)$,

$$x \cdot x = (\alpha^T \quad \beta^T) \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = 2\alpha^T \beta = 2m_i G_{ij} w^j + R^2 B_{ij} w^i w^j = 2m_i w^i$$

B is antisymmetric \rightarrow vanishes

Spectrum

$$\alpha_\mu = \frac{k_{L\mu} + k_{R\mu}}{\sqrt{2}},$$

$$\beta_\mu = \frac{k_{L\mu} - k_{R\mu}}{\sqrt{2}}$$

- We can describe important quantities of the CFT in the language of the lattice.
- The **spectral gap** (of primary states) = the energy difference between its ground state and first excited state :

$$\Delta = \min_{\substack{(\vec{k}_L, \vec{k}_R) \in \Lambda(R, G, B) \\ (\vec{k}_L, \vec{k}_R) \neq 0}} \frac{\vec{k}_L^2 + \vec{k}_R^2}{2} = \min_{\substack{(\alpha, \beta) \in \Lambda_N(R, G, B) \\ (\alpha, \beta) \neq \vec{0}}} \frac{\alpha^2 + \beta^2}{2}$$

- The **partition function** = $\text{Tr}_{\text{states}}[\exp(2\pi i \tau_1 P - 2\pi \tau_2 H)]$:

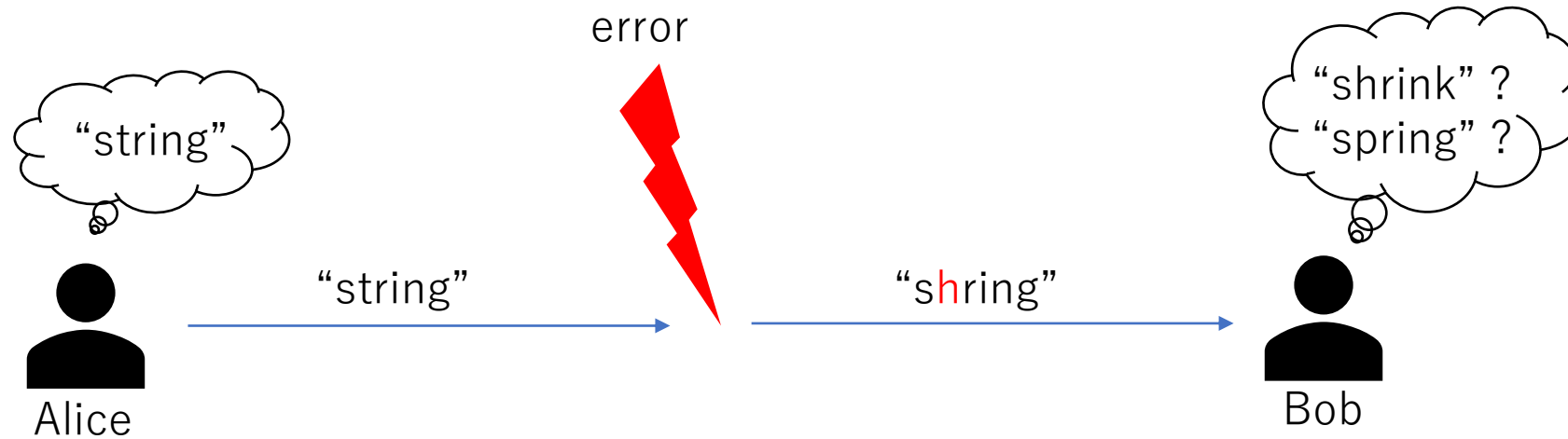
$$Z(\tau) = |\eta(\tau)|^{-2n} \sum_{(\vec{k}_L, \vec{k}_R) \in \Lambda(R, G, B)} q^{\vec{k}_L^2/2} \bar{q}^{\vec{k}_R^2/2}$$

$$= |\eta(\tau)|^{-2n} \sum_{(\alpha, \beta) \in \Lambda_N(R, G, B)} q^{(\alpha+\beta)^2/4} \bar{q}^{(\alpha-\beta)^2/4}$$

$\eta(\tau)$: Dedekind eta function,
 $q = e^{2\pi i \tau}$, $\bar{q} = e^{-2\pi i \bar{\tau}}$

2. Error correcting code

Error correcting code

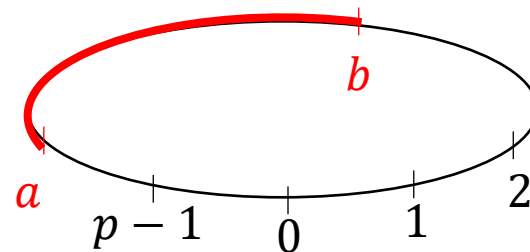


- An error correcting code is a concept in information theory for transmitting information correctly in spite of errors.

Finite field

- A finite field F is a field that contains a finite number of elements, which is a prime p or a prime power p^l .
- For a prime p , $F_p = \mathbb{Z}/p\mathbb{Z} = \{0, 1, \dots, p-1\}$
- We define a distance d between $a, b \in F_p^n$ by

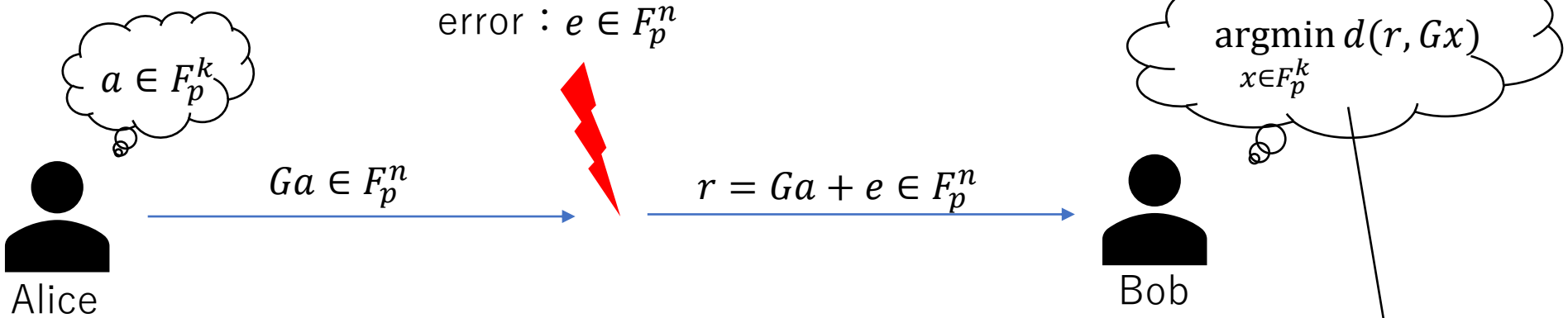
$$d(a, b) = \sqrt{\sum_{i=1}^n |a_i - b_i|^2}, \quad |a_i - b_i| = \min\{a_i - b_i, b_i - a_i\} (\in \mathbb{Z})$$



Error correction

$$n > k$$

- The error correction using an $n \times k$ matrix G on F_p :

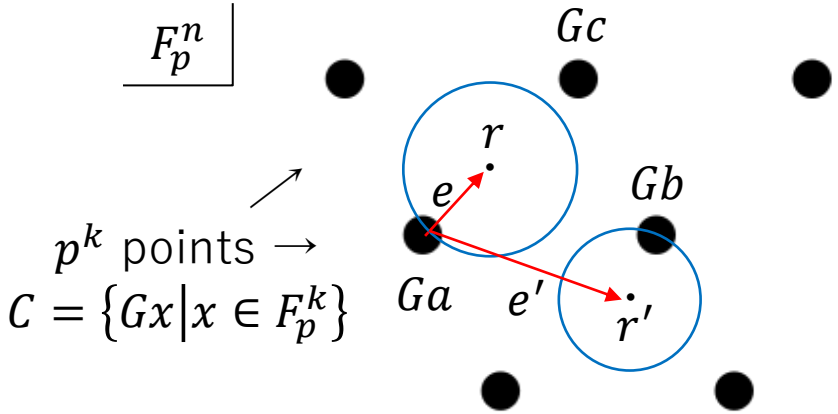


- We call $C = \{Gx \mid x \in F_p^k\} \subset F_p^n$ a code.

- Bob can get the correct message if

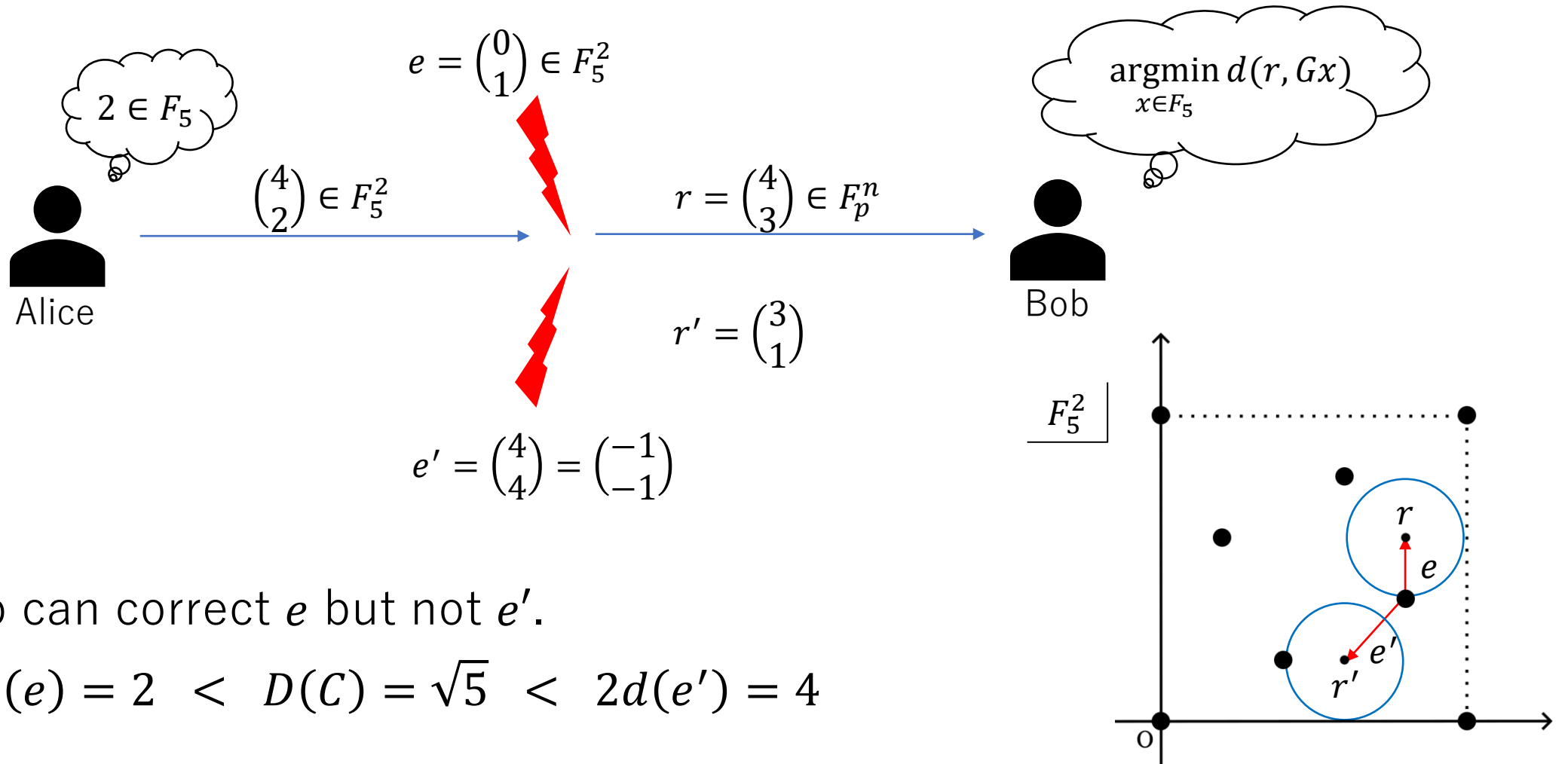
$$2d(e, 0) < D(C) := \min_{c, c' \in C, c \neq c'} d(c, c')$$

→ $D(C)$: error correction capability



Example : A code on F_5

- $G = \begin{pmatrix} 2 \\ 1 \end{pmatrix} \rightarrow C = \left\{ \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \begin{pmatrix} 2 \\ 1 \end{pmatrix}, \begin{pmatrix} 4 \\ 2 \end{pmatrix}, \begin{pmatrix} 1 \\ 3 \end{pmatrix}, \begin{pmatrix} 3 \\ 4 \end{pmatrix} \right\} \subset F_5^2, \quad D(C) = \sqrt{2^2 + 1^2} = \sqrt{5}$



- Bob can correct e but not e' .

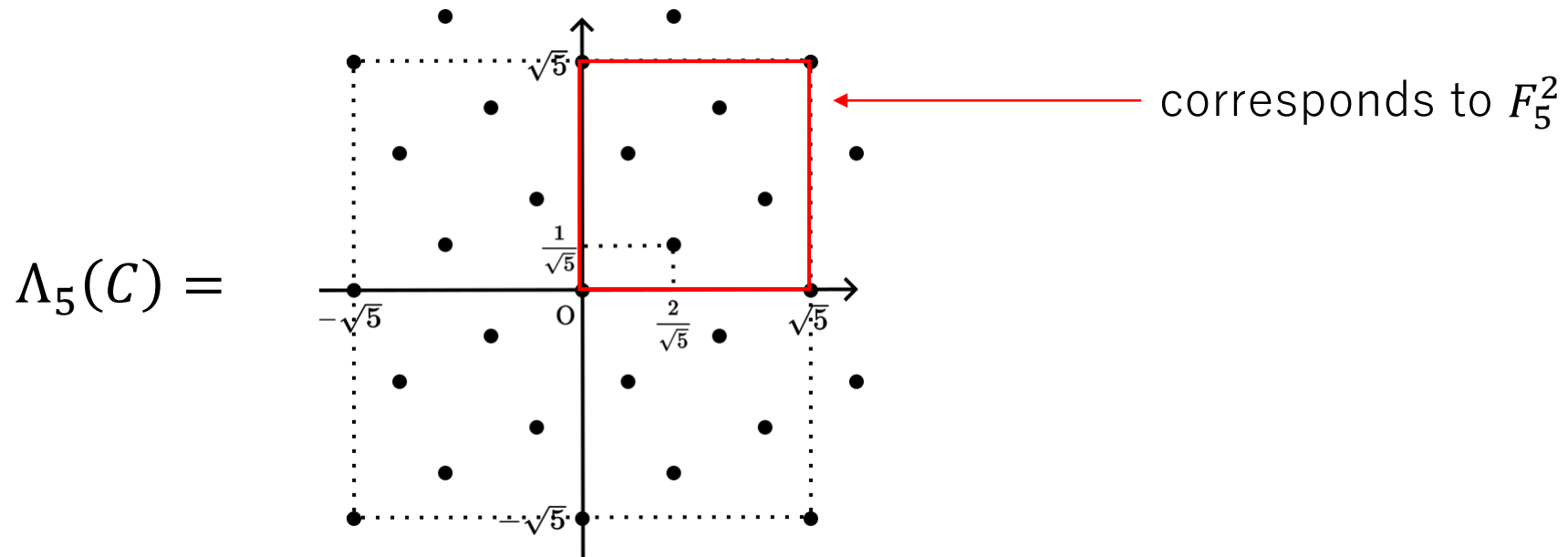
$$2d(e) = 2 < D(C) = \sqrt{5} < 2d(e') = 4$$

A lattice from a code

- We construct a lattice from a code $C \subset F_p^n$ by

$$\Lambda_p(C) = \left\{ \frac{c + pm}{\sqrt{p}} \mid c \in C, m \in \mathbb{Z}^n \right\} \subset \mathbb{R}^n$$

- e.g. For $C = \left\{ \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \begin{pmatrix} 2 \\ 1 \end{pmatrix}, \begin{pmatrix} 4 \\ 2 \end{pmatrix}, \begin{pmatrix} 1 \\ 3 \end{pmatrix}, \begin{pmatrix} 3 \\ 4 \end{pmatrix} \right\} \subset F_5^2$,



Even self-duality

- The case $p = 2$ was studied by Dymarsky and Shapere [1].
- **Prop.** For $p > 2$, the lattice $\Lambda_p(\mathcal{C})$ is **even** and **self-dual** with the metric $g = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$ if and only if \mathcal{C} is **self-dual**.

● Self-dual

- A dual code of $\mathcal{C} \subset F_p^n$: $\mathcal{C}^* = \{c' \in F_p^n \mid \forall c \in \mathcal{C}, c \cdot c' = 0\}$
- A code \mathcal{C} is self-dual : $\Leftrightarrow \mathcal{C} = \mathcal{C}^*$

- e.g. A code on F_5 generated by $G = \begin{pmatrix} 1 & 2 \\ 2 & 3 \\ 4 & 1 \\ 3 & 1 \end{pmatrix}$ is self-dual.

$$\because \begin{pmatrix} 1 \\ 2 \\ 4 \\ 3 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 2 \\ 4 \\ 3 \end{pmatrix} = 1 * 4 + 2 * 3 + 4 * 1 + 3 * 2 = 0, \quad \begin{pmatrix} 1 \\ 2 \\ 4 \\ 3 \end{pmatrix} \cdot \begin{pmatrix} 2 \\ 3 \\ 1 \\ 1 \end{pmatrix} = 1 * 1 + 2 * 1 + 4 * 2 + 3 * 3 = 0 \text{ etc.}$$

Proof (self-dual)

- The dual lattice of the code is the lattice of the dual code.

$$\begin{aligned} \because \quad & x' \in (\Lambda_p(\mathcal{C}))^* \\ & \Leftrightarrow \forall x \in \Lambda_p(\mathcal{C}), x \cdot x' \in \mathbb{Z} \\ & \Leftrightarrow \forall c \in \mathcal{C}, \forall m \in \mathbb{Z}^{2n}, \frac{1}{\sqrt{p}}(R(c) + pm) \cdot x' \in \mathbb{Z} && R \text{ is a map : } F_p \rightarrow \mathbb{Z} \\ & \Leftrightarrow \exists c' \in F_p^{2n}, \exists m' \in \mathbb{Z}^{2n}, x' = \frac{1}{\sqrt{p}}(R(c') + pm') \text{ and} \\ & \quad \forall c \in \mathcal{C}, \forall m \in \mathbb{Z}^{2n}, \frac{1}{p}(R(c) + pm) \cdot (R(c') + pm') \in \mathbb{Z} \\ & \Leftrightarrow \exists c' \in \mathcal{C}^*, \exists m' \in \mathbb{Z}^{2n}, x' = \frac{1}{\sqrt{p}}(R(c') + pm') \\ & \Leftrightarrow x' \in \Lambda_p(\mathcal{C}^*) \end{aligned}$$

- Thus, $(\Lambda_p(\mathcal{C}))^* = \Lambda_p(\mathcal{C}) \Leftrightarrow \mathcal{C}^* = \mathcal{C}$.

Self-dual code

- **Prop.** A code $C \subset F_p^n$ is self-dual if and only if n is even and C is generated by

$$G = \begin{pmatrix} I \\ X \end{pmatrix} \quad (\text{up to swapping rows})$$

where X is an $\frac{n}{2} \times \frac{n}{2}$ matrix s.t. $X + X^T = 0$ on F_p

- For the example on the previous page,

$$C = \left\{ \left(\begin{pmatrix} 1 & 2 \\ 2 & 3 \\ 4 & 1 \\ 3 & 1 \end{pmatrix} x \mid x \in F_5^2 \right) \right\} = \left\{ \left(\begin{pmatrix} 1 & 0 \\ 0 & 1 \\ 0 & 2 \\ 3 & 0 \end{pmatrix} y \mid y \in F_5^2 \right) \right\} \subset F_5^4$$

3. Relation

Relation through lattices

- Now, we constructed even self-dual lattices from a Narain CFT and a self-dual code.
- The simplest relation is the case where they form the same lattice.
- **Prop.** If a code $C \subset F_p^{2n}$ is generated by $G = \begin{pmatrix} I \\ X \end{pmatrix}$ where X is an $n \times n$ matrix s.t. $X + X^T = 0$,

$$\Lambda_N \left(R = \sqrt{\frac{2}{p}}, G = I, B = X \right) = \Lambda_p(C) \subset \mathbb{R}^{2n}$$

compactification radius \nearrow \uparrow metric \nwarrow antisymmetric background

\therefore Both lattices can be written as $\left\{ \begin{pmatrix} \sqrt{p}I & \frac{1}{\sqrt{p}}X \\ 0 & \frac{1}{\sqrt{p}}I \end{pmatrix} y \mid y \in \mathbb{Z}^{2n} \right\}$.

Correspondence in both theories

- Using this relation, we can consider the spectrum and the symmetries of the CFT in the language of the code.
- (rough summary in [3])

	CFT	Lattice Λ	Code \mathcal{C}
	modular invariance	even self-dual	self-dual
n	central charge	dimension	length
p	compactification radii $\sqrt{2/p}$	$\sqrt{p}\mathbb{Z}^{2n} \subset \Lambda$	on the finite field with p elements
	spectral gap	minimum length	correction capability
	partition function		enumerator polynomial

Partition function

- The partition function $Z(\tau)$ of the CFT can be written as the extended enumerator polynomial of the code.

$$Z(\tau) = |\eta(\tau)|^{-2n} \sum_{c \in \mathcal{C}} \prod_{x,y \in F_p} (t_{x,y})^{w_{x,y}(c)},$$

$$t_{x,y} = \sum_{m,l \in \mathbb{Z}} q^{(x+y+p(m+l))^2/4p} \bar{q}^{(x-y+p(m-l))^2/4p},$$

$$w_{x,y}(c) = |\{i \in \{1, \dots, n\} | (c_i, c_{i+n}) = (x, y)\}|$$

← code dependency

If $c = (1, 2, 2, 3, 1, 1)^T \in F_p^6$,
 $w_{1,3}(c) = 1, w_{2,1}(c) = 2$,
the others : 0

- We can relate

a symmetry of the CFT that keeps $Z(\tau)$ invariant to

a symmetry of the code that keeps polynomial invariant,

which have been studied separately.

Proof

$$\begin{aligned} & |\eta(\tau)|^{2n} Z(\tau) \\ &= \sum_{x \in \Lambda_N(r, I, B)} q^{\sum_{i=1}^n (x_i + x_{i+n})^2 / 4} \bar{q}^{\sum_{i=1}^n (x_i - x_{i+n})^2 / 4} \\ &= \sum_{y \in \Lambda_p(\mathcal{C})} q^{\sum_{i=1}^n (y_i + y_{i+n})^2 / 4} \bar{q}^{\sum_{i=1}^n (y_i - y_{i+n})^2 / 4} \\ &= \sum_{c \in \mathcal{C}} \sum_{m \in \mathbb{Z}^{2n}} \prod_{i=1}^n q^{(R(c_i) + pm_i + R(c_{i+n}) + pm_{i+n})^2 / 4p} \bar{q}^{(R(c_i) + pm_i - R(c_{i+n}) - pm_{i+n})^2 / 4p} \\ &= \sum_{c \in \mathcal{C}} \prod_{i=1}^n \sum_{m, l \in \mathbb{Z}} q^{(R(c_i) + R(c_{i+n}) + p(m+l))^2 / 4p} \bar{q}^{(R(c_i) - R(c_{i+n}) + p(m-l))^2 / 4p}. \end{aligned}$$

Spectral gap

- The spectral gap Δ of the CFT and the error correction capability $D(C)$ of the code satisfy

$$\Delta = \frac{1}{2p} \min\{D(C)^2, p^2\} \quad \leftarrow \text{In most cases, } D(C)^2 < p^2$$

- \rightarrow Searching for the code with high correction capability
= Searching for the Narain CFT with large spectral gap

include CFTs not related to codes

- The largest spectral gap among all Narain CFTs with n scalars is not well known for general n .
- Is this relation helpful?

Proof

$$\Delta = \min_{\substack{x \in \Lambda_N(r, I, B) \\ x \neq 0}} \frac{1}{2} x^T x = \min_{\substack{y \in \Lambda_p(\mathcal{C}) \\ y \neq 0}} \frac{1}{2} y^T y$$

$$= \min_{\substack{c \in \mathcal{C}, m \in \mathbb{Z}^{2n} \\ R(c) + pm \neq 0}} \frac{1}{2} \sum_{i=1}^{2n} \left(\frac{R(c_i) + pm_i}{\sqrt{p}} \right)^2$$

R is a map : $F_p \rightarrow \mathbb{Z}$

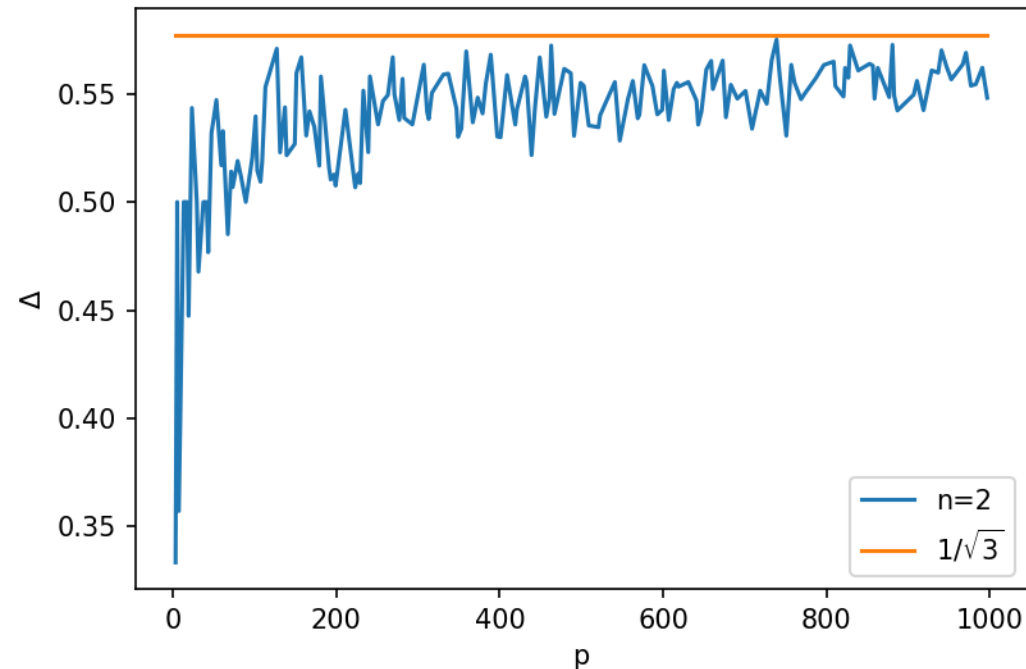
$$= \frac{1}{2p} \min \left\{ \min_{\substack{c \in \mathcal{C}, m \in \mathbb{Z}^{2n} \\ c \neq 0}} \sum_{i=1}^{2n} (R(c_i) + pm_i)^2, \min_{\substack{m \in \mathbb{Z}^{2n} \\ m \neq 0}} \sum_{i=1}^{2n} (pm_i)^2 \right\}$$

$$= \frac{1}{2p} \min \left\{ \min_{\substack{c \in \mathcal{C} \\ c \neq 0}} \sum_{i=1}^{2n} \min\{R(c_i)^2, (R(c_i) - p)^2\}, p^2 \right\}$$

$$= \frac{1}{2p} \min \{ D(\mathcal{C})^2, p^2 \}.$$

Spectral gap, $n = 2$

- From numerical calculations, the largest spectral gap of Narain CFTs corresponding to codes on F_p^2 is as follows :



- The values suggest that $1/\sqrt{3}$ is their upper bound, which can be checked analytically by reducing to the sphere packing in two dim.

Spectral gap, $n = 3$

- For $a \in \mathbb{Z}$ s.t. $p = (a^4 + 1)/2$ is a prime number, we consider a code C on F_p generated by

$$G = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & -a & -a^2 \\ a & 0 & -a^3 \\ a^2 & a^3 & 0 \end{pmatrix}.$$

- The error correction capability :

$$D(C) = \left| G \begin{pmatrix} (a-1)/2 \\ (a-1)/2 \\ 0 \end{pmatrix} \right| = \sqrt{(3a^4 - 4a^3 + 6a^2 - 4a + 3)/4}$$

- The spectral gap of the corresponding CFT :

$$\Delta = \frac{1}{2p} \min\{D(C)^2, p^2\} = \frac{(3a^4 - 4a^3 + 6a^2 - 4a + 3)}{4(a^4 + 1)} \xrightarrow{a \rightarrow \infty} \frac{3}{4} \quad \leftarrow \text{The largest known spectral gap for } n = 3 \text{ [2] !}$$

4. Future prospects

Code on finite field F_{p^l}

- We considered only finite fields with prime elements.
- For a prime power p^l , $F_{p^l} = F_p[x]/(f_{p,l}(x)) = \{ \sum_{t=0}^{l-1} a_t x^t \mid a_t \in F_p \}$

\nearrow
polynomial ring on F_p

\nwarrow
Conway polynomial
- E.g. $F_{3^2} = F_3[x]/(x^2 + 2x + 2) = \{ a_2 x^2 + a_1 x + a_0 \mid a_2, a_1, a_0 \in F_3 \}$
 $(x^2 + 1) \times (x + 2) = x^3 + 2x^2 + x + 2 = 2x + 2$
- It is difficult to relate a code on F_{p^l} to a self-dual lattice than on F_p .
→ Can it correspond to a more general CFT?

Spectral gap for large n

- Through the correspondence between quantum gravity and CFT, the spectral gap corresponds to the energy difference in gravity theory.
- We do not know
 - the largest spectral gap and
 - how to construct a CFT with large spectral gap for large n .
- Can we answer these using the relation between CFTs and codes?

Thank you for listening.

References

- [1] A. Dymarsky and A. Shapere, *Quantum stabilizer codes, lattices, and CFTs*, J. High Energ. Phys. 2021, 160 (2021) [arXiv:2009.01244].

- [2] N. Afkhami-Jeddi, H. Cohn, T. Hartman and A. Tajdini, *Free partition functions and an averaged holographic duality*, J. High Energ. Phys. 2021, 130 (2021) [arXiv:2006.04839].

- [3] Shinichiro Yahagi, *Narain CFTs and error-correcting codes on finite fields*, arXiv:2203.10848.