

$\bar{B} \rightarrow D^* \tau \bar{\nu}$  における

$D^*$  偏極

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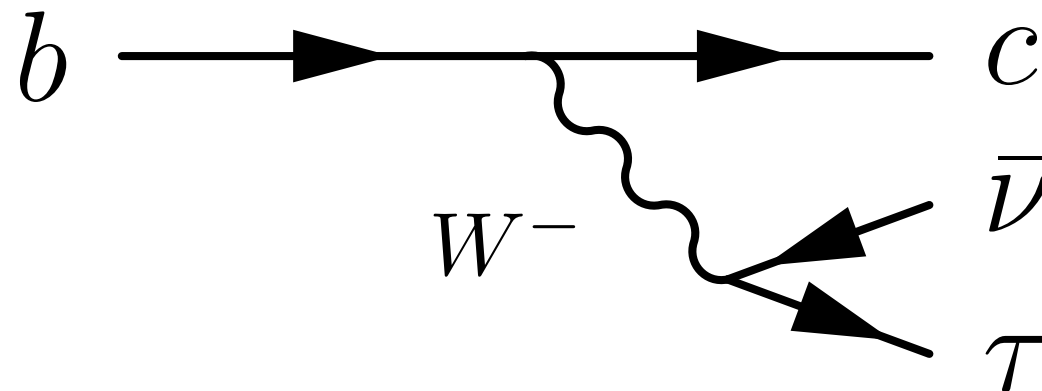
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# Introduction

Semi-tauonic B decays



$$\bar{B} \rightarrow D^{(*)} \tau \bar{\nu}_\tau$$

$$\frac{\mathcal{B}(\bar{B} \rightarrow D \tau \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D \ell \bar{\nu}_\ell)} = 0.40 \pm 0.08$$

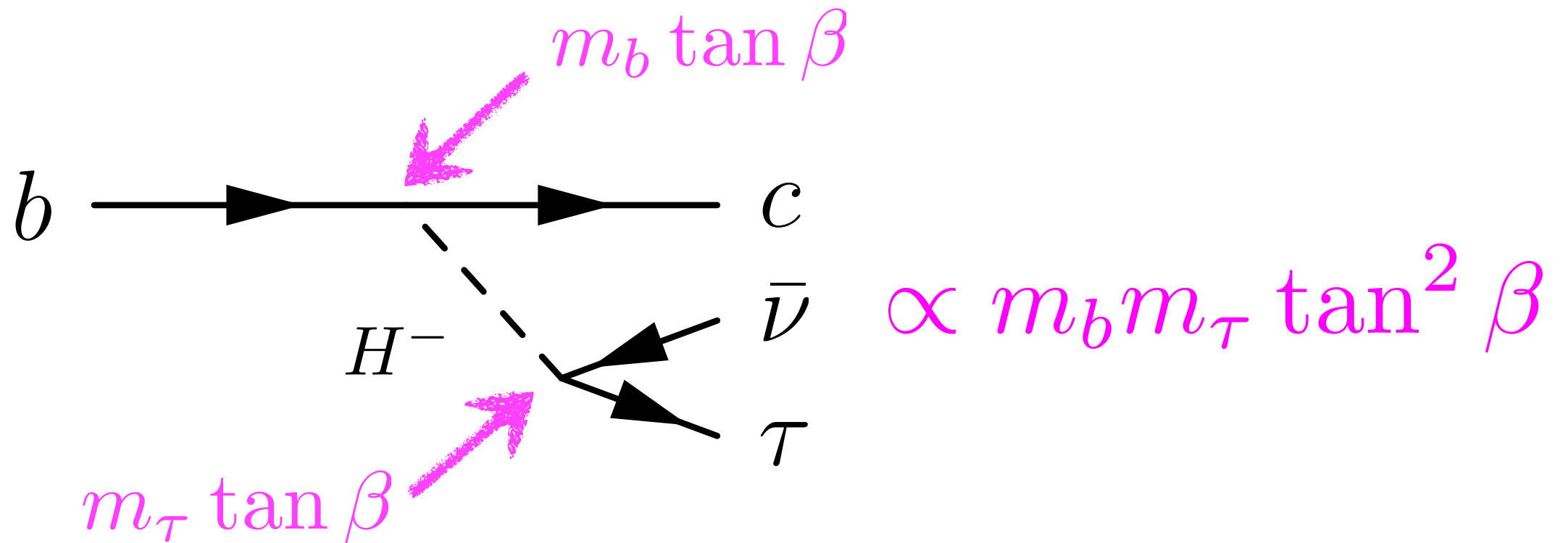
$$\frac{\mathcal{B}(\bar{B} \rightarrow D^* \tau \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D^* \ell \bar{\nu}_\ell)} = 0.35 \pm 0.04$$

(BABAR, Belle combined.)

# Charged Higgs contribution

W.S. Hou and B. Grzadkowski (1992),  
M.T. (1995), ....

## Type-II 2HDM (SUSY)



Sensitive to the charged Higgs  
if  $\tan \beta$  is large.

# D vs D\*

	pros	cons
D	simple (2 FFs) sensitive to scalar	less observables less statistics (Br~0.7%)
D*	D* polarization more statistics (Br~1.3%)	complicated (4 FFs) less sensitive to scalar

## Form factors in $\bar{B} \rightarrow D^* \tau \bar{\nu}$

$$w = v \cdot v'$$

$$\langle D^*(v') | \bar{c} \gamma_\mu b | B(v) \rangle = i h_V(w) \epsilon_{\mu\nu\alpha\beta} \epsilon^{*\nu} v'^\alpha v^\beta$$

$$\begin{aligned} \langle D^*(v') | \bar{c} \gamma_\mu \gamma_5 b | B(v) \rangle &= h_{A_1}(w) (1+w) \epsilon_\mu^* \\ &\quad - h_{A_2}(w) \epsilon^* \cdot v v_\mu \\ &\quad - h_{A_3}(w) \epsilon^* \cdot v v'_\mu \end{aligned}$$

## Heavy Quark Limit

$$h_V = h_{A_1} = h_{A_3} = \xi(w), \quad \xi(1) = 1$$

$$h_{A_2} = 0$$



Isgur-Wise func.

# Parametrization

Caprini et al. (1998)

$$z = \frac{\sqrt{w+1} - \sqrt{2}}{\sqrt{w+1} + \sqrt{2}}, \quad r = \frac{m_{D^*}}{m_B}$$

$$h_{A_1}(w) = h_{A_1}(1) \left\{ 1 - 8\rho^2 z + (53\rho^2 - 15)z^2 - (231\rho^2 - 91)z^3 \right\}$$

$$R_1(w) \equiv \frac{h_V(w)}{h_{A_1}(w)} \quad \boxed{\bar{B} \rightarrow D^* \ell \bar{\nu}} \quad \text{exp. data}$$

$$= R_1(1) - 0.12(w-1) + 0.05(w-1)^2$$

$$R_2(w) \equiv \frac{h_{A_3}(w) + r h_{A_2}(w)}{h_{A_1}(w)}$$

$$= R_2(1) + 0.11(w-1) - 0.06(w-1)^2$$

$$R_3(w) \equiv \frac{h_{A_3}(w) - r h_{A_2}(w)}{h_{A_1}(w)}$$

$$\boxed{\bar{B} \rightarrow D^* \tau \bar{\nu}} \quad \text{only}$$

# Numerical results

## Inputs

$$\bar{B} \rightarrow D^* \ell \bar{\nu} \quad (\ell = e, \mu)$$

$$\rightarrow \rho^2 = 1.24 \pm 0.04$$

**HFAG**

$$R_1(1) = 1.41 \pm 0.049$$

$$R_2(1) = 0.844 \pm 0.027$$

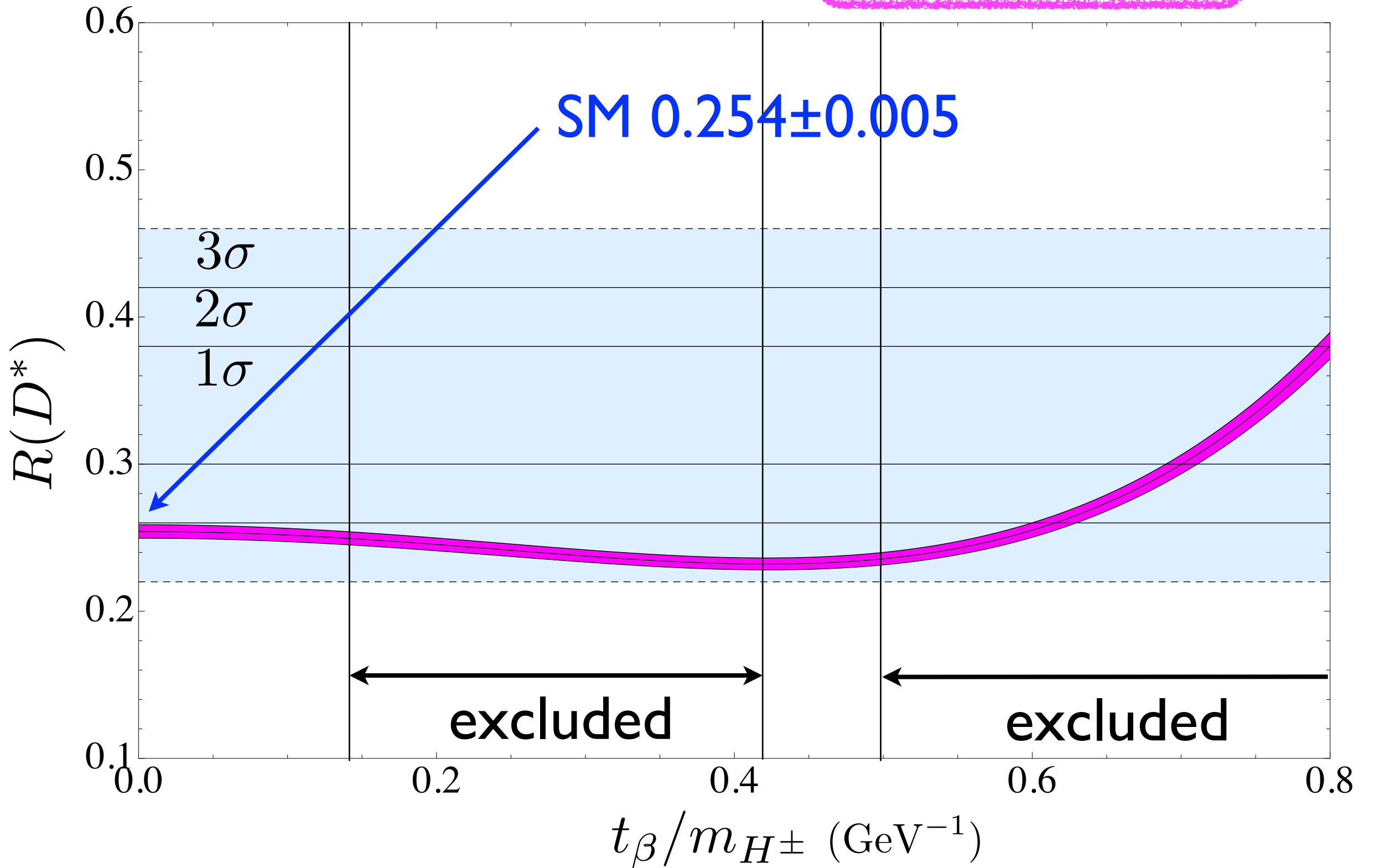
## HQET

$$R_3(w) = 1 + a \{ 0.22 - 0.052(w - 1) + 0.026(w - 1)^2 \}$$

$$a = 1 \pm 0.5$$

# Branching fraction ( $D^*$ )

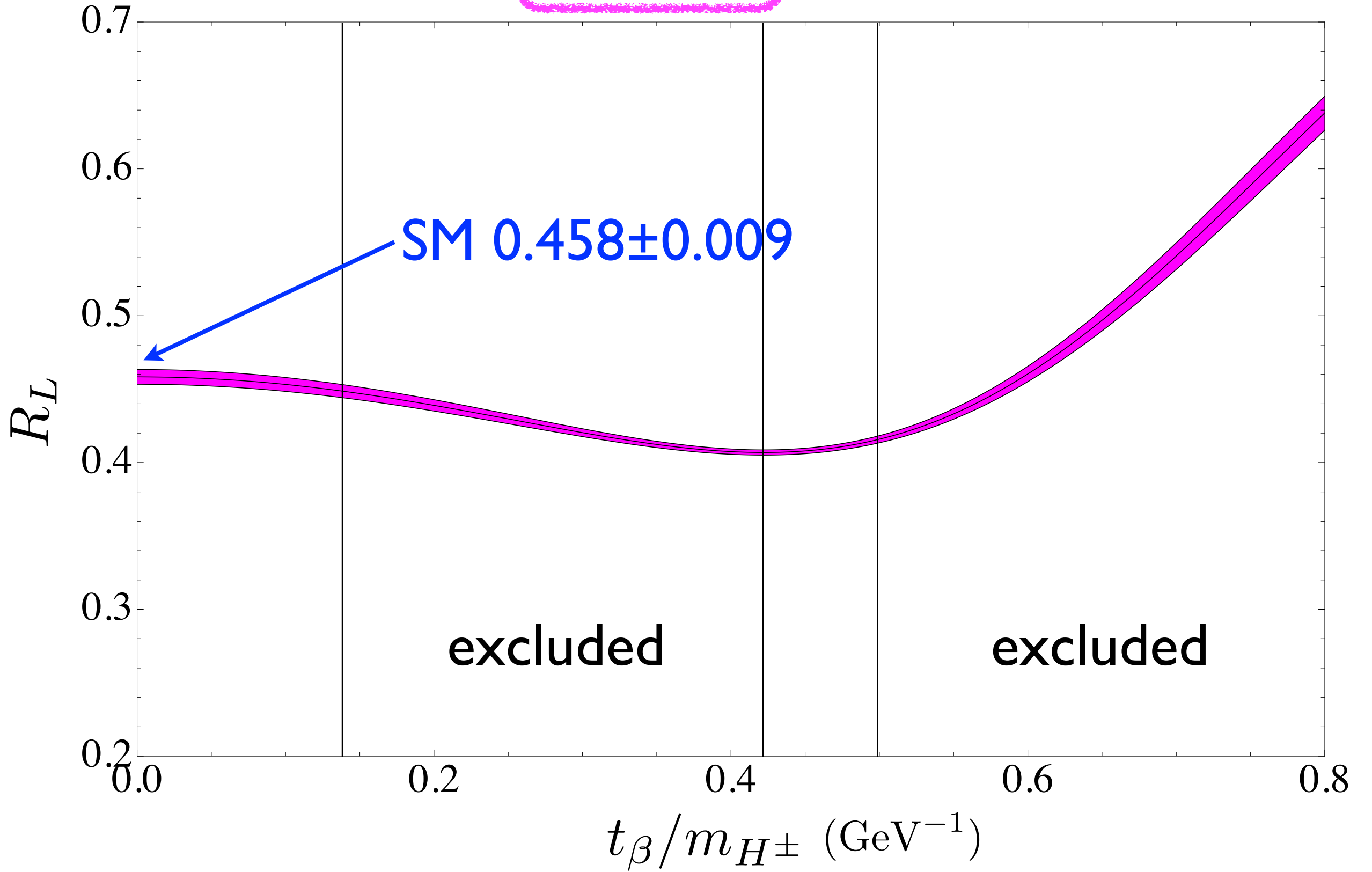
$$\frac{\mathcal{B}(\bar{B} \rightarrow D^* \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D^* \ell^- \bar{\nu}_\ell)}$$





# D\* polarization

$$\frac{D_L^*}{D_T^* + D_L^*}$$



# $D^*$ decay and polarization

## Helicity amplitudes

$$\mathcal{M}(D_T^* \rightarrow D\pi) \sim Y_{\pm}^1(\theta, \varphi)$$

$$\mathcal{M}(D_L^* \rightarrow D\pi) \sim Y_0^1(\theta, \varphi)$$

## Angular distribution

$$R_L = \frac{D_L^*}{D_T^* + D_L^*}$$

$$W(\cos \theta) = f(\cos \theta) + (2R_L - 1)g(\cos \theta)$$

$$f(\cos \theta) = \frac{3}{8}(1 + \cos^2 \theta), \quad g(\cos \theta) = \frac{3}{8}(3 \cos^2 \theta - 1)$$

## Sensitivity


$$S = \left[ \int \frac{g^2}{f + (2R_L - 1)g} d \cos \theta \right]^{1/2} \Rightarrow \sim 0.66$$

Statistical error:  $\delta R_L = \frac{1}{2S\sqrt{N}}$

**B factory**

$N \sim 500$    $\delta R_L \sim 0.03$

**Super B factory**

$N \sim 2 \times 10^4$    $\delta R_L \sim 5 \times 10^{-3}$

**cf. SM prediction:**  $R_L = 0.458 \pm 0.009$

# Summary

★ 4 form factors in  $\bar{B} \rightarrow D^* \tau \bar{\nu}$

$h_{A_1}, R_1, R_2$  determined by  $\bar{B} \rightarrow D^* \ell \bar{\nu}$ ,  
 $R_3$  determined by HQET.

Uncertainties are well controlled.  $\lesssim 5\%$

★  $D^*$  polarization

SM prediction  $R_L = 0.458 \pm 0.009$

B factory  $\delta R_L \sim 0.03$

Super B factory  $\delta R_L \sim 5 \times 10^{-3}$

- ★  $D^*$  less sensitive to charged Higgs than  $D$ .
- ★ But, different dependence on  $t_\beta/m_{H^\pm}$ .
- The 2-fold ambiguity can be solved within the semi-tauonic decays.
- ★  $D^*$  complementary to  $D$ .

# Backup Slides

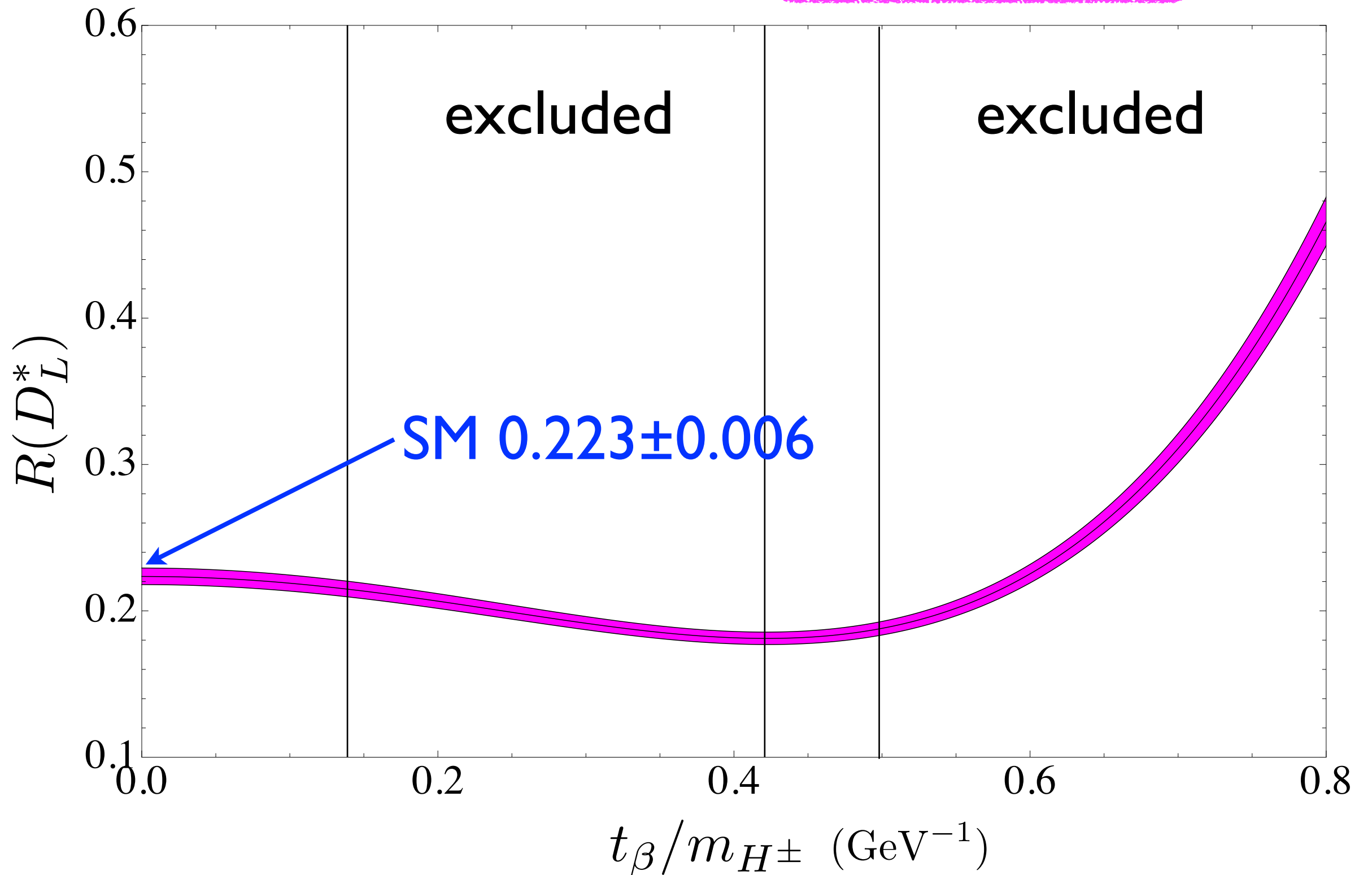
	$L = 0$	$L = 1$	$L = 2$
$B \rightarrow DW^*$	✓	✓	
$B \rightarrow D_T^* W^*$	✓	✓	✓
$B \rightarrow D_L^* W^*$	✓	✓	✓
$B \rightarrow DH^*$	✓		
$B \rightarrow D_T^* H^*$			
$B \rightarrow D_L^* H^*$		✓	



suppressed

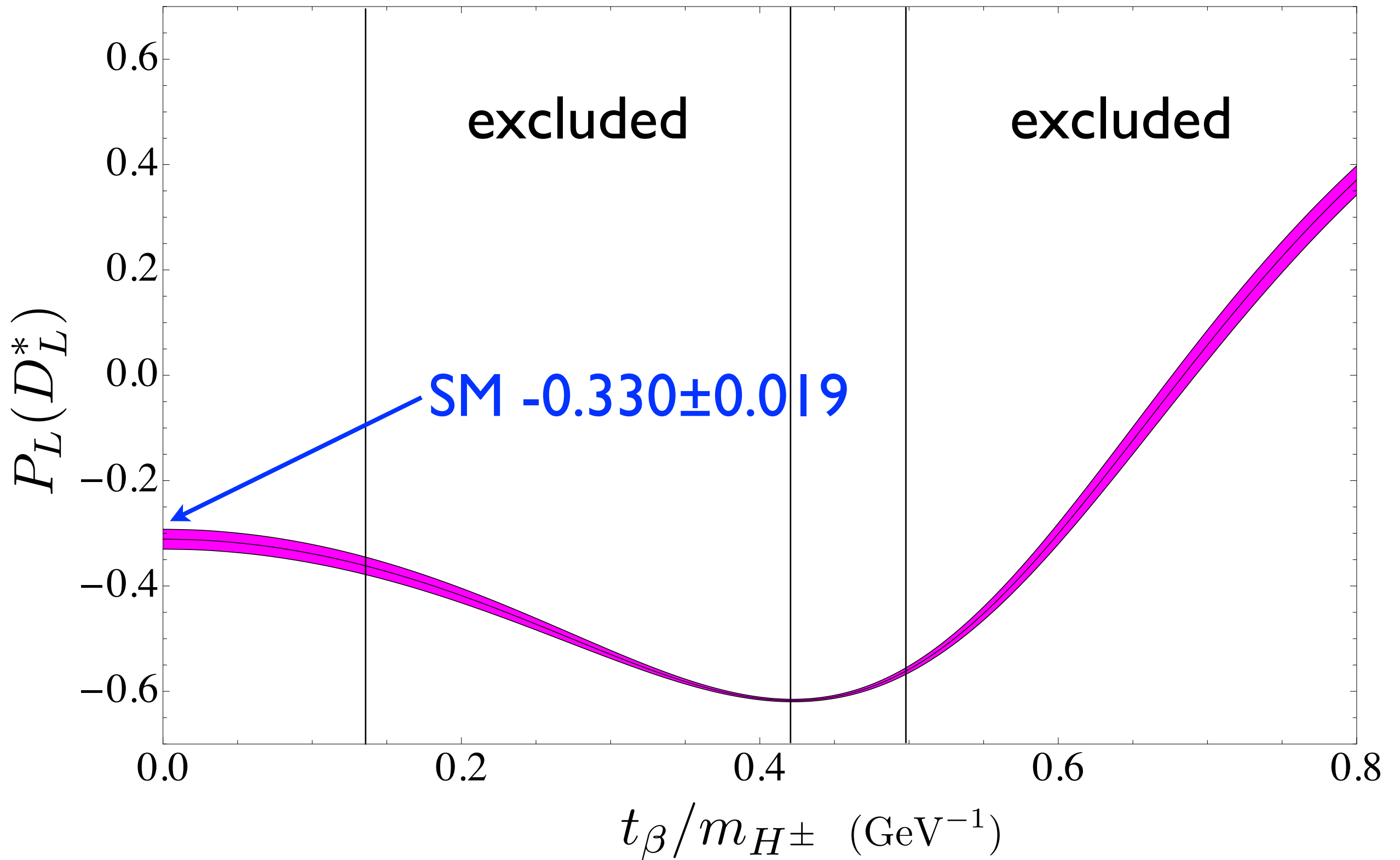
Branching fraction ( $D_L^*$ )

$$\frac{\mathcal{B}(\bar{B} \rightarrow D_L^* \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D_L^* \ell^- \bar{\nu}_\ell)}$$





# Tau longitudinal polarization



# Sensitivity to tau polarization

$$S(\tau \rightarrow \pi\nu(\ell\bar{\nu}\nu)) \simeq 0.61(0.20)$$

$$\delta P_\tau = \frac{1}{S\sqrt{N}}$$

**B factory**

$$N \sim 500 \quad \longrightarrow \quad \delta P_\tau \sim 0.07(0.22)$$

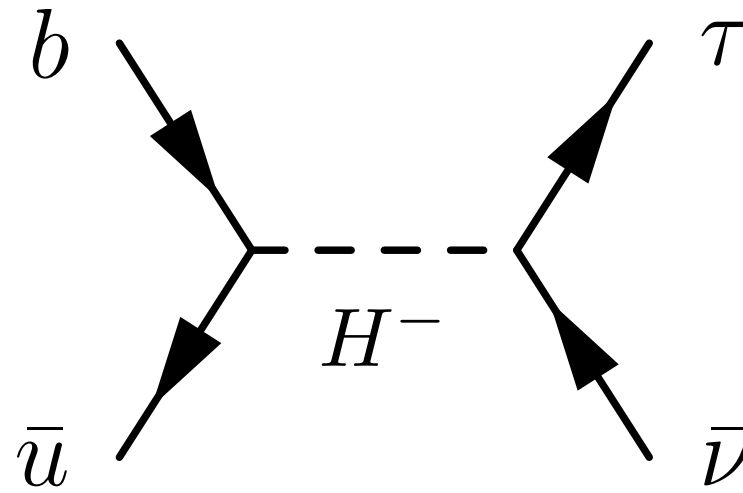
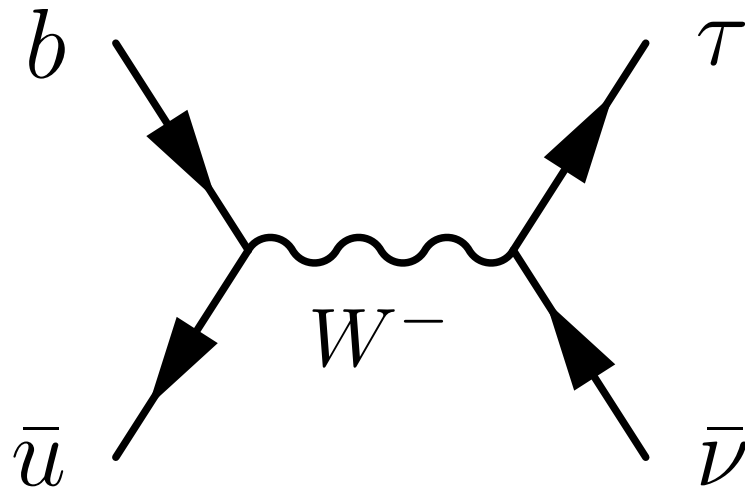
**Super B factory**

$$N \sim 2 \times 10^4 \quad \longrightarrow \quad \delta P_\tau \sim 0.01(0.04)$$

cf. SM prediction:  $P_\tau = -0.330 \pm 0.019$

# Pure tauonic B decay

$$B^- \rightarrow \tau \bar{\nu}$$



$$\mathcal{B} = (1.67 \pm 0.39) \times 10^{-4}$$

(HFAG)

Uncertainties

$$V_{ub}, f_B$$

Nierste et al.

