

# イッテルビウム同位体シフトデータの 一般化線形性同時解析による新物理探索

田中 実 (阪大理)

共同研究者 山本康裕 (NCTS→KEK)

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# Isotope shift (IS) and King linearity

Level-splitting difference between isotopes

$$\text{IS} = \nu_{A'A} := \nu_{A'} - \nu_A$$

IS of two transitions:  $t = 1, 2$

$$\nu_{A'A}^{(t)} = K_t \mu_{A'A} + F_t \langle r^2 \rangle_{A'A}$$

mass shift (MS) field shift (FS)

$$\text{Modified IS: } \tilde{\nu}_{A'A}^{(t)} := \nu_{A'A}^{(t)} / \mu_{A'A} = K_t + F_t \langle r^2 \rangle_{A'A} / \mu_{A'A}$$

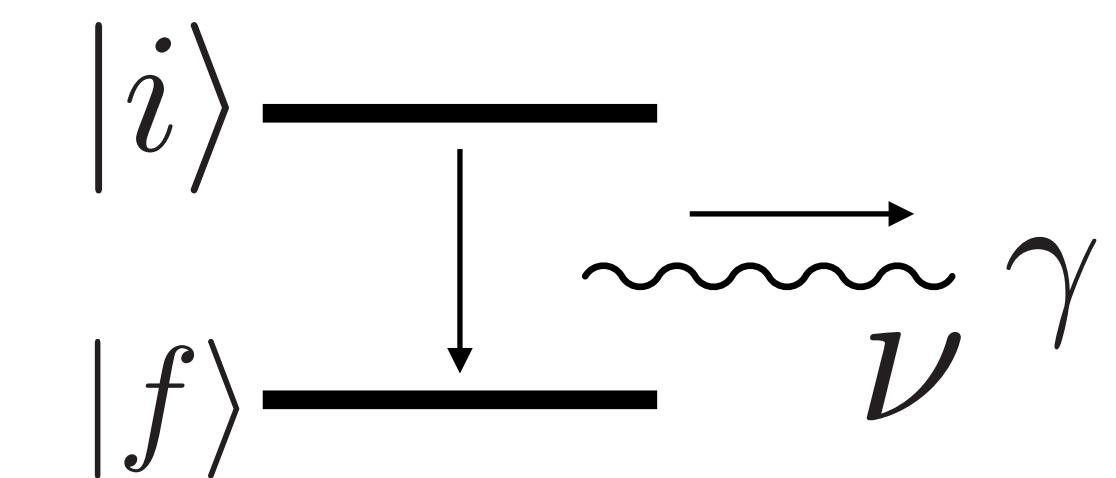
electronic factors nuclear factor

King linearity: eliminating the nuclear factor King, 1963

$$\tilde{\nu}_{A'A}^{(2)} = K_{21} + F_{21} \tilde{\nu}_{A'A}^{(1)}$$

$$K_{21} := K_2 - F_{21} K_1, \quad F_{21} := F_2 / F_1$$

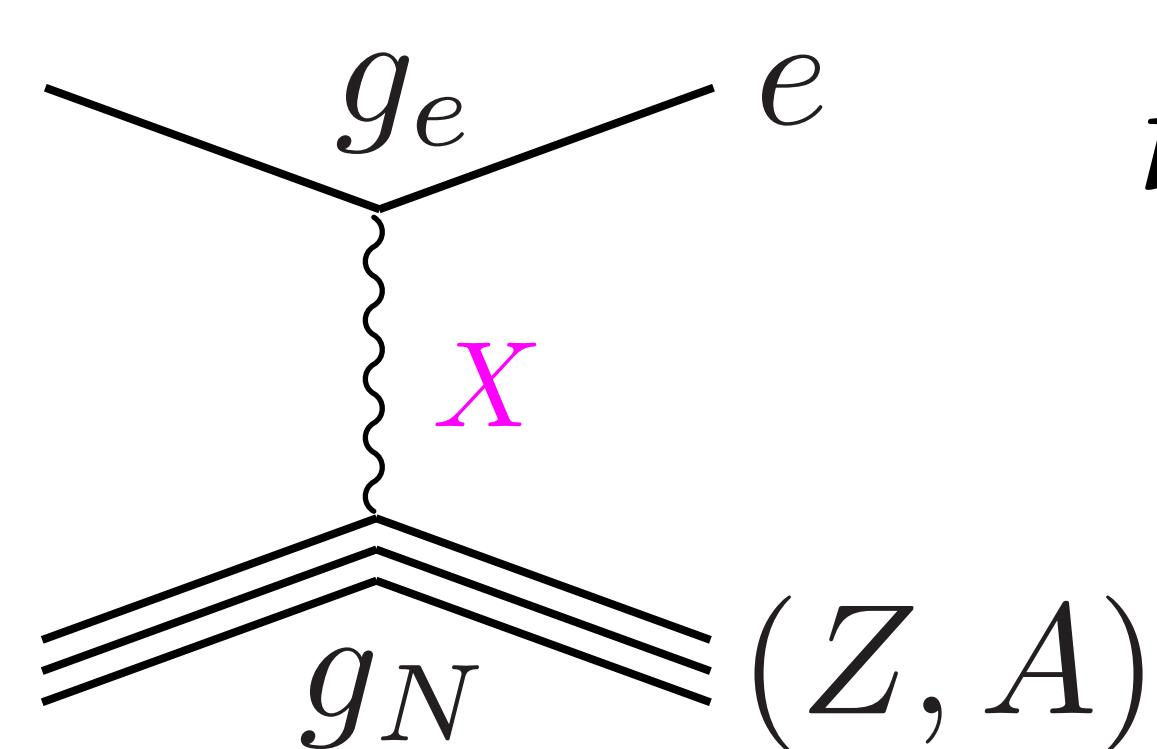
→  $(\tilde{\nu}_{A'A}^{(1)}, \tilde{\nu}_{A'A}^{(2)})$  on a straight line, King plot



# Nonlinearity

# IS by new neutron-electron interaction

Delaunay et al. arXiv:1601.05087v2



$$\nu_{A'A}^{(t)} = K_t \mu_{A'A} + F_t \langle r^2 \rangle_{A'A} + X_t (A' - A)$$

MS      
 FS      
 particle shift (PS)

# Nonlinearity due to subleading FS

$$\text{FS} = F_t \langle r^2 \rangle_{A'A} + F'_t [\langle r^2 \rangle_{A'A}]^2 + G_t \langle r^4 \rangle_{A'A} + \dots$$

# quadratic FS ( $\text{FS}_{22}$ )   higher moment ( $\text{FS}_4$ )

$$[\langle r^2 \rangle_{A'A}]^2 := (\langle r^2 \rangle_{A'A_0})^2 - (\langle r^2 \rangle_{AA_0})^2$$

# Generalized linearity

K. Mikami, MT, Y. Yamamoto EPJC77:896 (2017)

$$\nu_{A'A}^{(t)} = K_t \mu_{A'A} + F_t \langle r^2 \rangle_{A'A} + F'_t [\langle r^2 \rangle_{A'A}]^2 + X_t (A' - A)$$

3 transitions:  $t=1, 2, 3$

$\mathsf{FS}_{22}$

$\mathsf{PS}$

$$\begin{pmatrix} \nu_{A'A}^{(1)} - X_1(A' - A) \\ \nu_{A'A}^{(2)} - X_2(A' - A) \\ \nu_{A'A}^{(3)} - X_3(A' - A) \end{pmatrix} = \begin{pmatrix} K_1 & F_1 & F'_1 \\ K_2 & F_2 & F'_2 \\ K_3 & F_3 & F'_3 \end{pmatrix} \begin{pmatrix} \mu_{A'A} \\ \langle r^2 \rangle_{A'A} \\ [\langle r^2 \rangle_{A'A}]^2 \end{pmatrix} =: M \begin{pmatrix} \mu_{A'A} \\ \langle r^2 \rangle_{A'A} \\ [\langle r^2 \rangle_{A'A}]^2 \end{pmatrix}$$

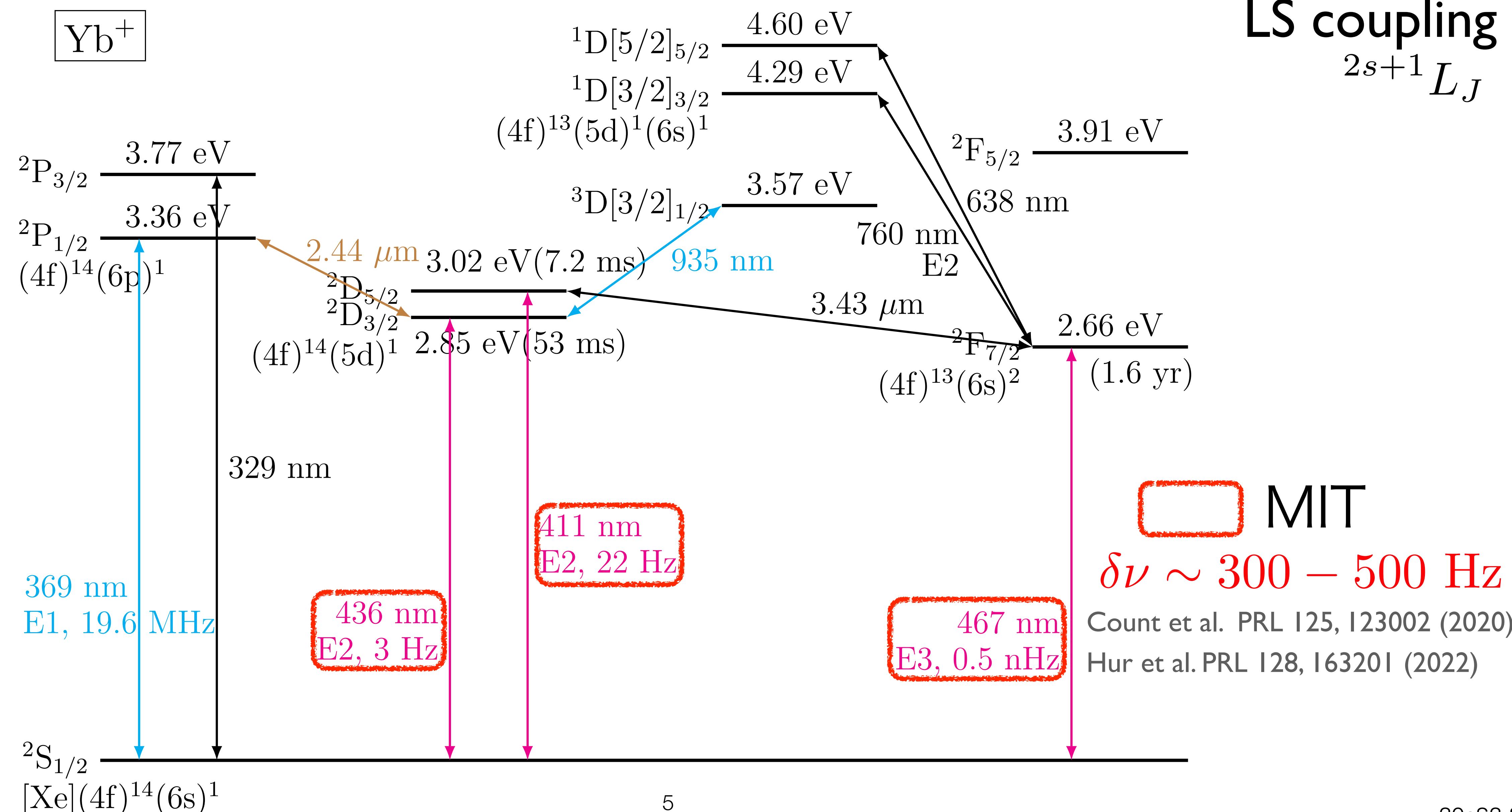
→  $(M^{-1})_{11}\nu_{A'A}^{(1)} + (M^{-1})_{12}\nu_{A'A}^{(2)} + (M^{-1})_{13}\nu_{A'A}^{(3)}$   
 $- \{(M^{-1})_{11}X_1 + (M^{-1})_{12}X_2 + (M^{-1})_{13}X_3\}(A' - A) = \mu_{A'A}$

$$(\nu_{A'A}^{(1)}, \nu_{A'A}^{(2)}, \nu_{A'A}^{(3)})/\mu_{A'A}$$

on a plane if  $X_t = 0$

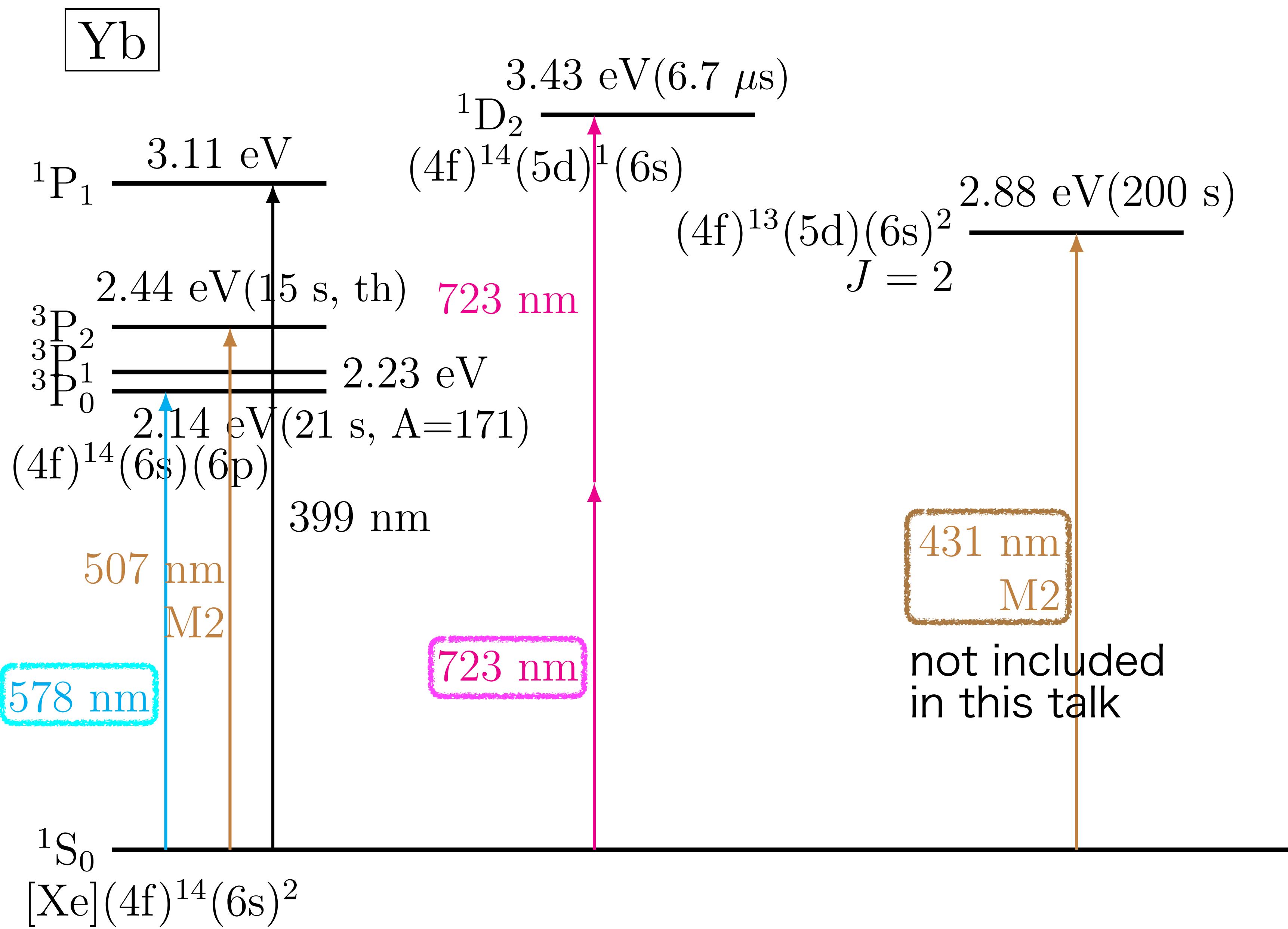
n transitions and n+1 IS pairs → NP search with n-2 NL's removed

# Data of Yb transitions



# LS coupling

$$2s+1 L_J$$



Yb

Kyoto

$\delta\nu \sim$  a few Hz

K. Ono, MT et al. PRX 12, 021033 (2022)

Mainz

Figueroa et al. PRL 128, 073001 (2022)

$\delta\nu \sim O(100)$  Hz

Kyoto

T. Ishiyama et al. PRL 130, 153402 (2023)

石山さんの講演 20pA2-4

NMIJ,AIST

A. Kawasaki et al. arXiv:2402.13541

# Simultaneous analysis of plural linearities

Y.Yamamoto, MT et al. in preparation

Yb: 4 independent IS pairs

3D linearity is nontrivial, while 5 transitions available.

→ Combined fit of 3 independent 3D relations

$$\nu_3 = k_3\mu + f_{31}\nu_1 + f_{32}\nu_2, \quad \nu_4 = k_4\mu + f_{41}\nu_1 + f_{42}\nu_2, \quad \nu_5 = k_5\mu + f_{51}\nu_1 + f_{52}\nu_2$$

$$\delta\chi^2/\text{dof} = 35.1/3 \quad (\text{A}'\text{A omitted})$$

Assigning this nonlinearity to PS,

→  $\delta\chi^2/\text{dof} = 9.98/1$

PS alone cannot explain the observed 3D nonlinearity.

# Consistency with SM sources

$$FS = F_t \langle r^2 \rangle_{A'A} + F'_t [\langle r^2 \rangle_{A'A}]^2 + G_t \langle r^4 \rangle_{A'A} + \dots$$

$$FS_{22} \qquad \qquad FS_4$$

Using theoretically calculated electronic factors,  
we subtract  $FS_{22}$  and/or  $FS_4$ .

	D5	D3	F7	P0	D2	
FS <sub>22</sub> (a)	81.908	83.247	-201.12	54.277	75.322	(AMBiT) Hur et al. PRL 128, 163201 (2022)
FS <sub>4</sub> (g)	14.934	15.159	-39.422	8.951		(GRASP2018)
FS <sub>4</sub> (r)	13.08				10.42	(RPA) Figueroa et al. PRL 128, 073001 (2022) (MHz/fm <sup>4</sup> )

Abbreviation: (Yb<sup>+</sup>) D5 := <sup>2</sup>S<sub>1/2</sub> - <sup>2</sup>D<sub>5/2</sub>, D3 := <sup>2</sup>S<sub>1/2</sub> - <sup>2</sup>D<sub>3/2</sub>, F7 := <sup>2</sup>S<sub>1/2</sub> - <sup>2</sup>F<sub>7/2</sub>  
(Yb) P0 := <sup>1</sup>S<sub>0</sub> - <sup>3</sup>P<sub>0</sub>, D2 := <sup>1</sup>S<sub>0</sub> - <sup>1</sup>D<sub>2</sub>

## Hypothesis: FS<sub>4</sub> and one unknown

No FS<sub>4</sub> in  $\nu'_i := \nu_i - p_{i/i_0} \nu_{i_0}$ ,  $p_{i/i_0} := G_i/G_{i_0}$  ratio of FS<sub>4</sub> electronic factor

→ Two independent 3D linear relations ( $i_0 = \text{D5}$ )

$$\nu'_i = k_i \mu + d_i \nu'_{\text{P0}} + e_i \nu'_{\text{D2}}, \quad i = \text{D3, F7}$$

$\delta\chi^2/\text{dof} = 13.3/2$  without theoretical uncertainty

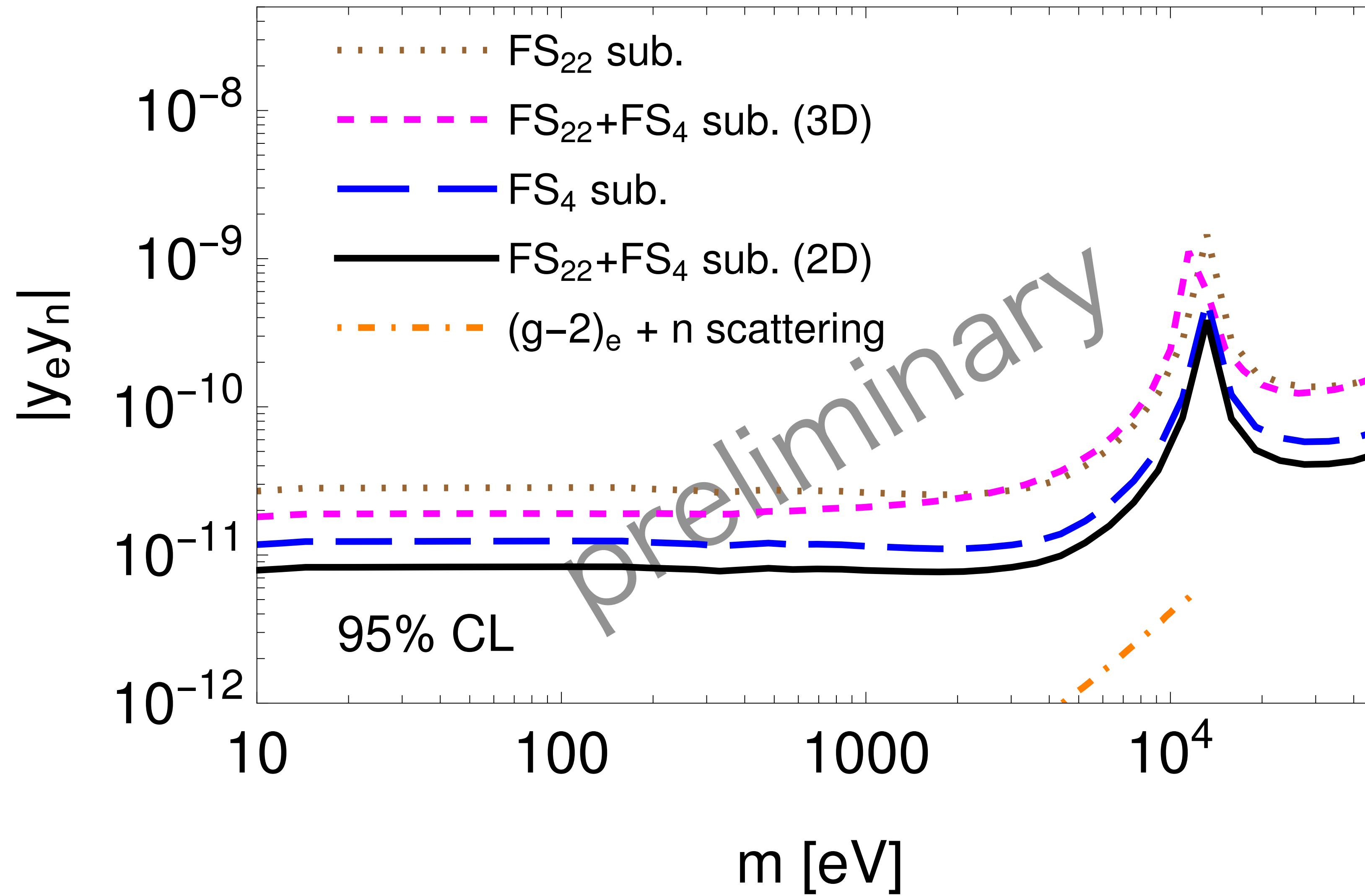
$\delta\chi^2/\text{dof} = 1.8/2$  with 6% theoretical uncertainty for  $p$ 's

→ Limit on new physics (not to make the fit worse)

$$\nu'_i := \nu_i - p_{i/i_0} \nu_{i_0} - \frac{y_e y_n}{4\pi} X_i(m)(A' - A)$$

$$|y_e y_n| < 1.2 \times 10^{-11}, \quad m = 10 \text{ eV} \quad (95\% \text{ CL})$$

# Summary of limit on new physics



# まとめと展望

- Isotope shift and King linearity  $\tilde{\nu}_{A'A}^{(2)} = K_{21} + F_{21}\tilde{\nu}_{A'A}^{(1)}$   
IS=MS+FS, linear relation of mIS of two transitions
- Nonlinearities: New physics and/or SM higher order
- Generalized linearity  
SM nonlinearity removed, improved sensitivity to new physics  
Two or more SM higher order contributions revealed  
One SM higher order + PS, excluded by our combined analysis
- New precise Yb IS data  
Yb<sup>+</sup> ion O(10) Hz, MPI-PTB Door et al. arXiv2403.07792  
improvement of MIT data and Yb masses
- Yb atom O(?) Hz, Kyoto 石山さんの講演 20pA2-4  
 $^1S_0(6s)^2 - (4f)^{13}(5d)(6p)^2 (J = 2)$ , 431 nm