

Bubbling Geometries for Half BPS Wilson Lines

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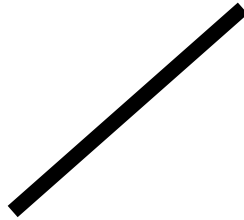
S. Yamaguchi, to appear



1. Overview

AdS₅

CFT₄



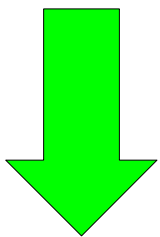
||

AdS₅ x S⁵

||

4dim N=4

Super YM theory



deform

⊃

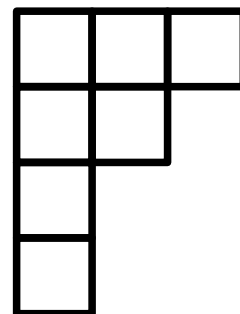
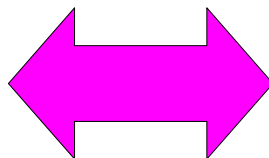
Goal

Supergravity
Solutions

?

Wilson loop
operators

$$\text{Tr}_R P \exp[i\oint (A + \phi)]$$



2. 4 dim N=4 U(N)SYM

Field Contents

Trivial reduction of 10 dim SYM

- Vector $A_\mu, \mu=0, 1, 2, 3$
- Spinors λ 16 real components
- Scalars $\phi_i, i=4, \dots, 9$

Each field is N x N Hermitian Matrix

Symmetries

R-symmetry $SO(6)=SU(4)$

A_μ : Singlet λ : Spinor ϕ_i : Vector

Supersymmetry

$$\delta A_\mu = i \bar{\lambda} \Gamma_\mu \epsilon$$

$$\delta \phi_i = i \bar{\lambda} \Gamma_i \epsilon$$

$$\delta \lambda = \dots$$

Conformal Symmetry $SO(2,4)=SU(2,2)$

Gauge invariant operators

1/2 BPS local operators

$$O_{\vec{k}}(x) = \prod_j \text{Tr} [Z^j]^{k_j}(x) \quad Z := \phi_4 + i\phi_5$$

$$\vec{k} := (k_1, k_2, k_3, \dots), k_j \in \mathbb{Z}_{\geq 0}$$

This operator is invariant under the SUSY transformation

$$\delta_\epsilon O_{\vec{k}} = 0, \quad \text{for } i\Gamma_{45}\epsilon = \epsilon.$$

1/2 BPS straight Wilson line

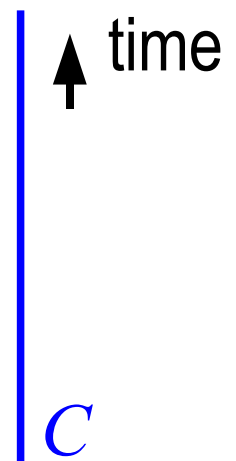
$$W_{\vec{k}} = \prod_j \text{Tr} [U^j]^{k_j},$$

$$U = P \exp \int_C ds [A_\mu \dot{x}^\mu + \phi_4 |\dot{x}|]$$

$$= P \exp \int dx^0 [A_0 + \phi_4]$$

This operator is invariant under the SUSY transformation

$$\delta_\epsilon W_{\vec{k}} = 0, \quad \text{for } \Gamma_{04}\epsilon = \epsilon.$$

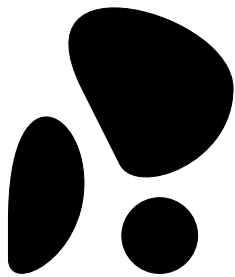


3. LLM Solution [Lin, Lunin, Maldacena]

$\frac{1}{2}$ BPS local operators that preserves
 $SO(4) \times SO(4)$ symmetry

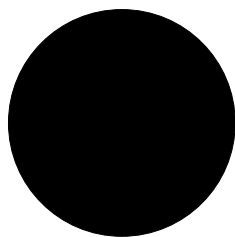
Supergravity solutions

“Droplet” $z : \mathbb{R}^2 \rightarrow \{-1, +1\}$



$\frac{1}{2}$ BPS smooth
solution

Example



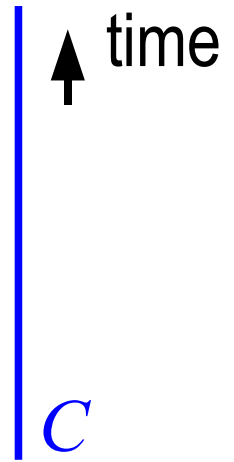
$AdS_5 \times S^5$

Interpretation

The phase space of free fermions derived
from Gaussian matrix quantum mechanics

4. Symmetry of a Straight Wilson line.

$$U := P \exp \int dx^0 [A_0 + \phi_4]$$

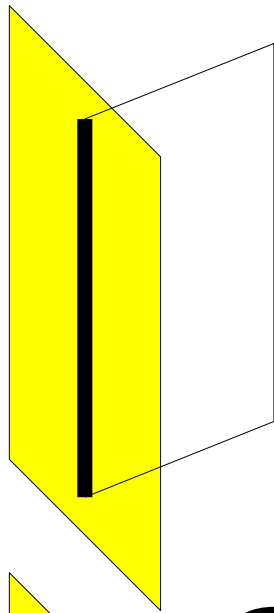


- Space-like rotation $SO(3)$
 - Dilatation
 - Time translation
 - Special Conformal Symmetry
 - time direction
 - Part of R-symmetry $SO(5)$
- } $SL(2, \mathbb{R})$

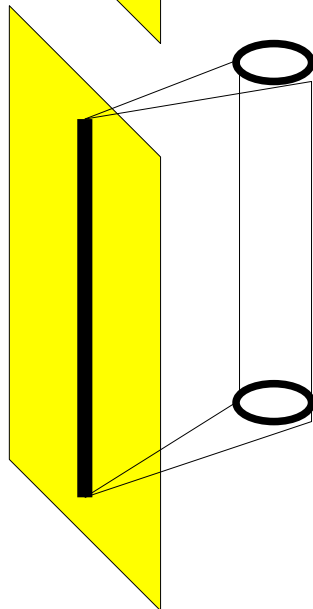
$$SL(2, \mathbb{R}) \times SO(3) \times SO(5)$$

5. Probe Picture

[Ray, Yee],[Maldacena]
[Drukker,Fiol]



AdS2 F-String



AdS2 x S2 D3-brane
with electric flux

Both preserves $SL(2,R) \times SO(3) \times SO(5)$ and
 $\frac{1}{2}$ Supersymmetry

Goal: The SUGRA solution (backreaction)



Stack of these branes

Note

AdS₂ x S² D3-brane

$$S_{CS} = \int F \wedge C_2, \quad C_2 : \text{RR 2-form}$$

The SUGRA solution has RR 2-form excitation

6. Ansatz

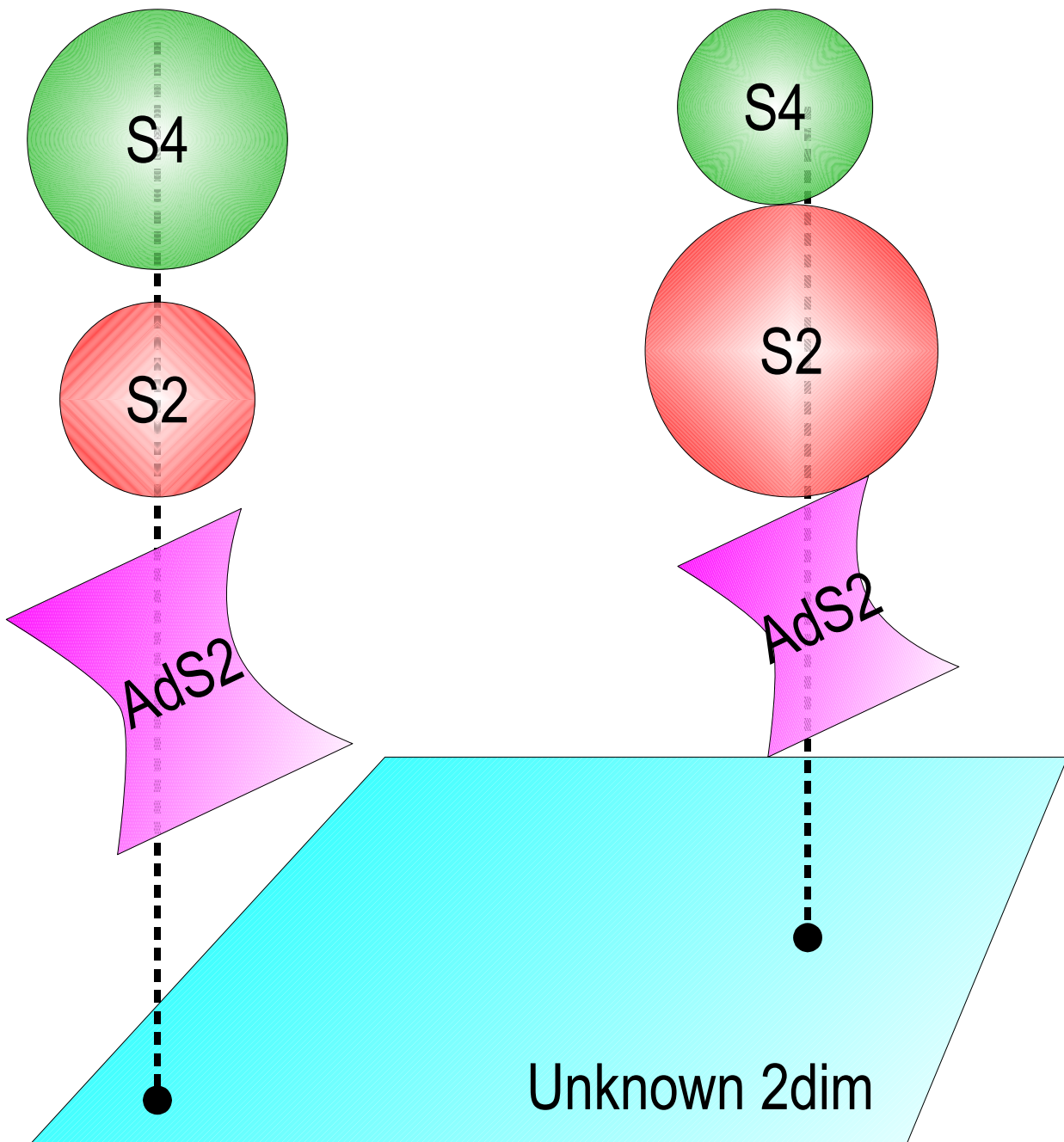
SL(2,R)

SO(3)

SO(5)

$$ds^2 = e^{2A} (AdS_2) + e^{2B} (S^2) + e^{2C} (S^4) + ds_2^2$$

Unknown



Field strength

AdS2

2dim

S2

S4

3-form

$$\boxed{\text{vol(AdS2)}} \wedge \boxed{\text{1-form}}$$

$$\boxed{\text{1-form}} \wedge \boxed{\text{vol(S2)}}$$

5-form

$$\boxed{\text{vol(AdS2)}} \wedge \boxed{\text{1-form}} \wedge \boxed{\text{vol(S2)}}$$

$$\boxed{\text{1-form}} \wedge \boxed{\text{vol(S4)}}$$

1-form

$$\boxed{\text{1-form}}$$

We can reduce every unknown quantities to 2 dimensional fields

Truncation

Background D3-brane

$$G_5 = J \wedge \text{vol}(AdS_2) \wedge \text{vol}(S^2) + \tilde{J} \wedge \text{vol}(S^4)$$

F-string

$$H_3 = F \wedge \text{vol}(AdS_2)$$

A equation of motion implies

$$G_3 = K \wedge \text{vol}(S^2)$$

Other fields can be put as 0 consistently

Unknown quantities

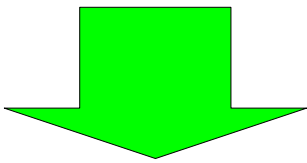
2dim metric	ds_2^2
4 scalars	A, B, C, ϕ
4 1-forms	J, \tilde{J}, F, K

7. The Structure of the Geometry

SUSY

$$\delta \psi_{\mu} = 0, \quad \delta \lambda = 0$$

Necessary conditions

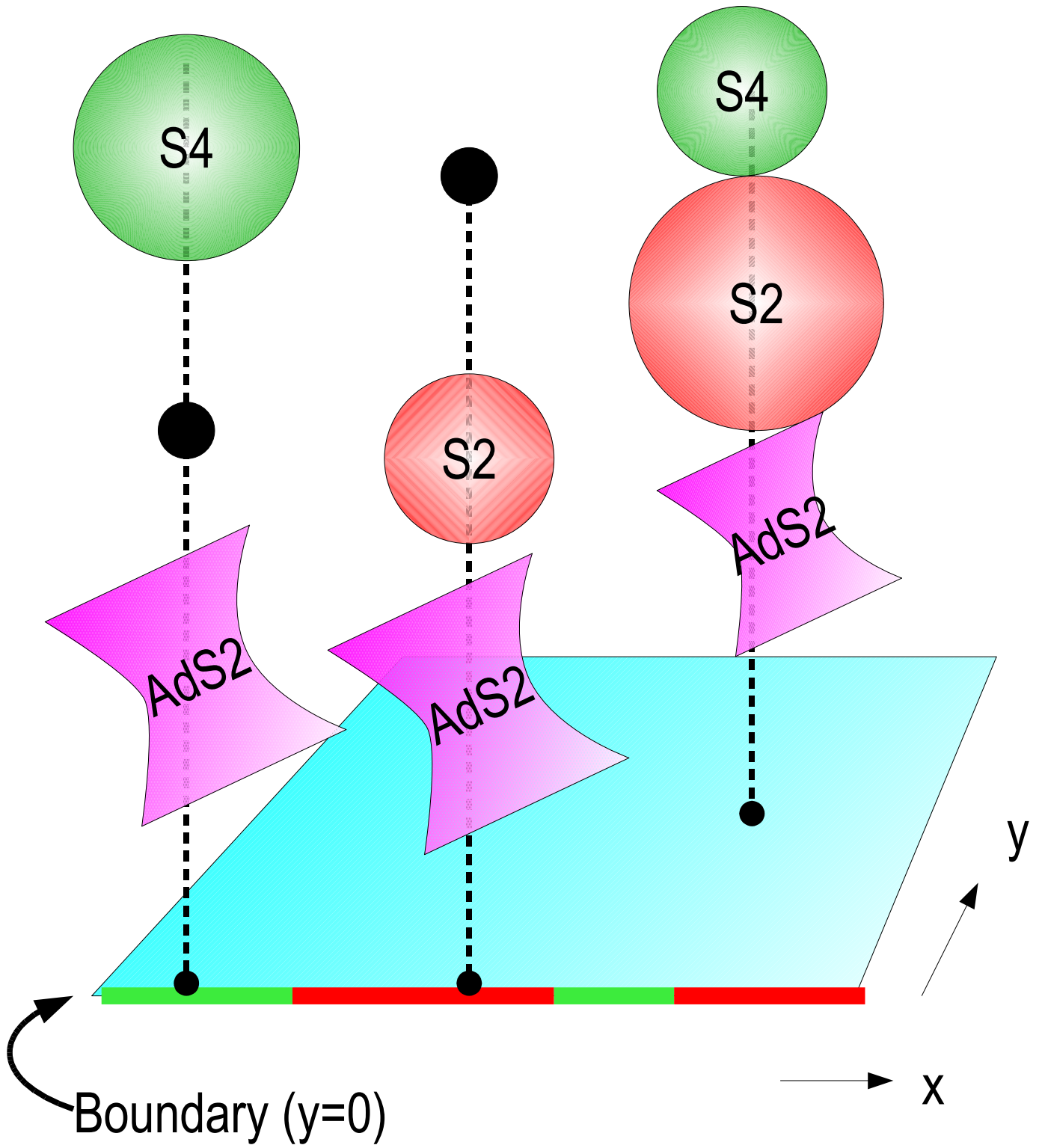


$$y := \exp(B + C)$$

Product of radii of
S² and S⁴

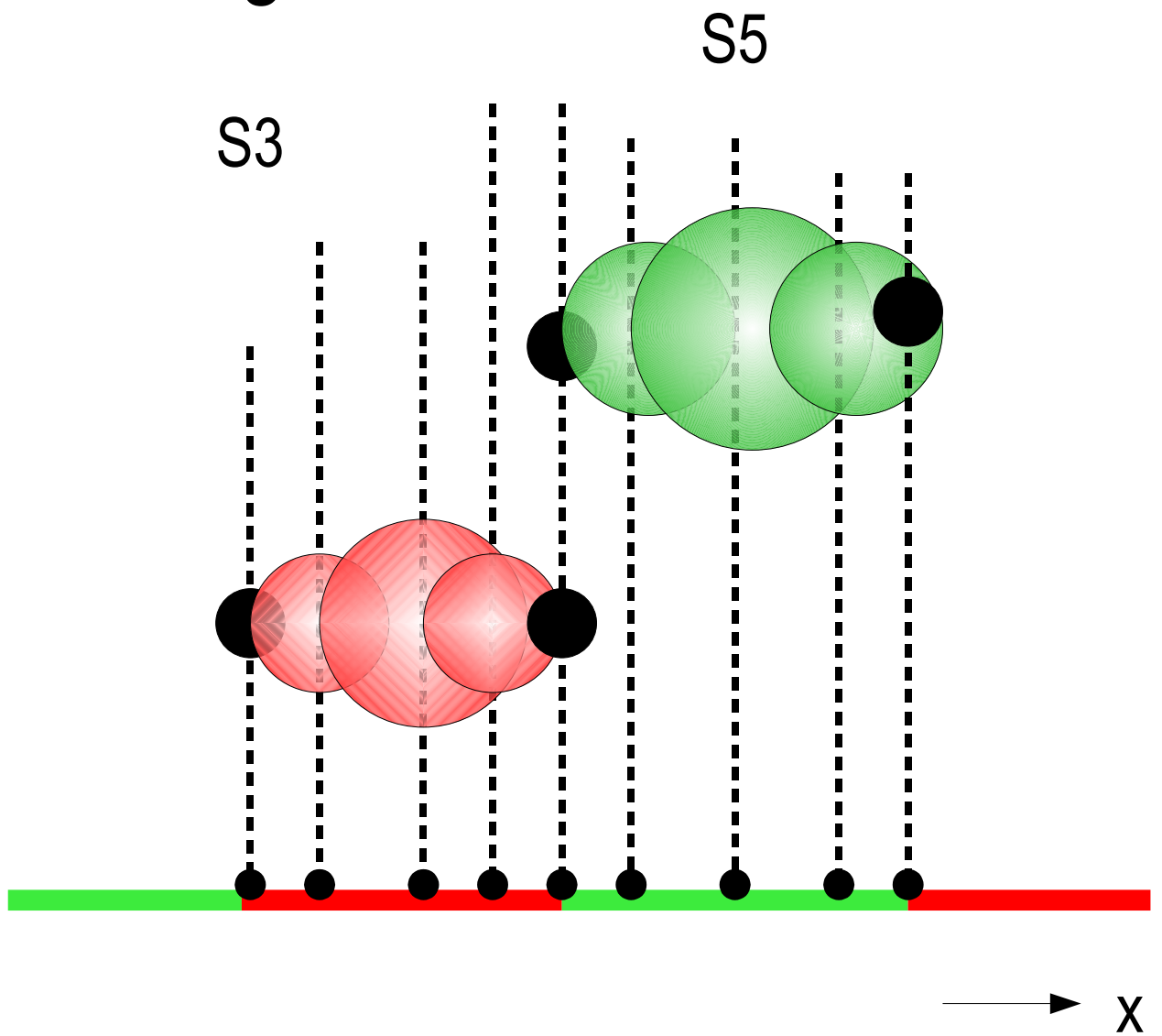
$$ds^2 = e^{2A} (AdS_2) + e^{2B} (S^2) + e^{2C} (S^4) + ds_2^2$$

$$ds_2^2 = \frac{1}{e^{2B} + e^{2C}} (dy^2 + dx^2)$$

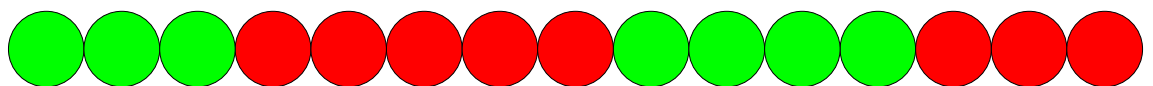


$$y = (\text{radius } S2) \times (\text{radius } S4)$$

Bubbling



discrete



Maya diagram

Relation to the label of Wilson line ?

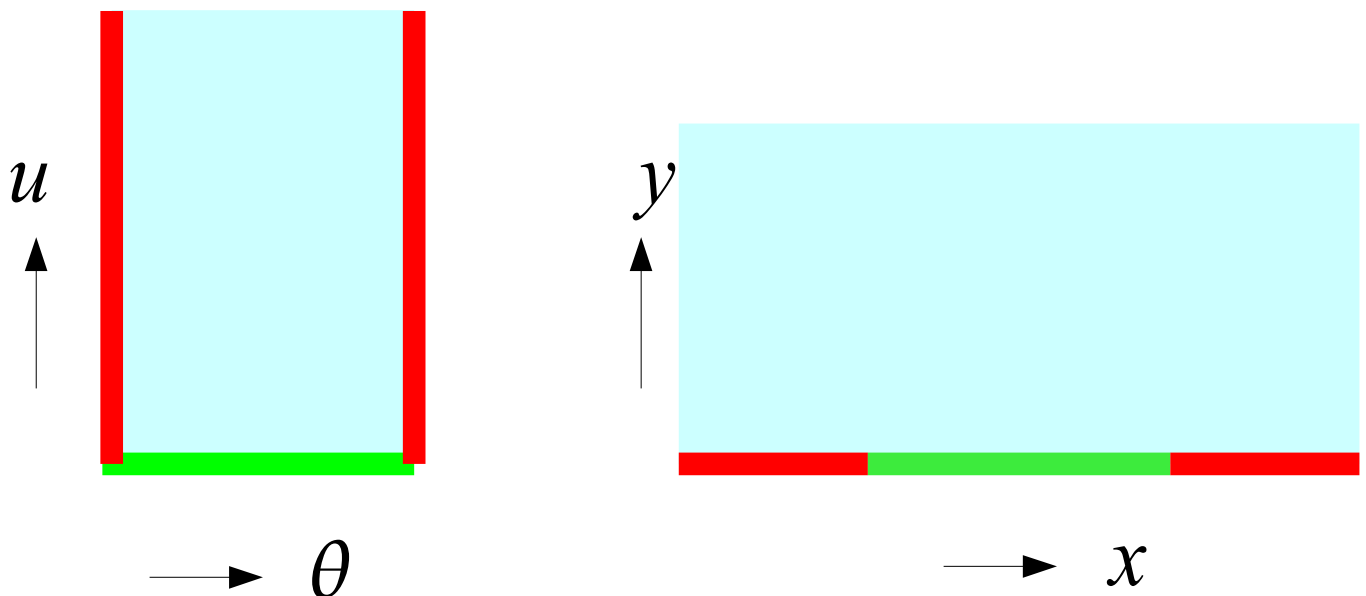
Example: AdS5 x S5 solution

$$ds^2 = R^2 \cosh^2 u (AdS_2) + R^2 \sinh^2 u (S^2) \\ + R^2 \sin^2 \theta (S^4) + R^2 (du^2 + d\theta^2)$$

$$R^2 = \sqrt{4\pi g_s N}$$

$$u \geq 0, \quad 0 \leq \theta \leq \pi.$$

$$y = R^2 \sinh u \sin \theta, \\ x = R^2 \cosh u \cos \theta.$$



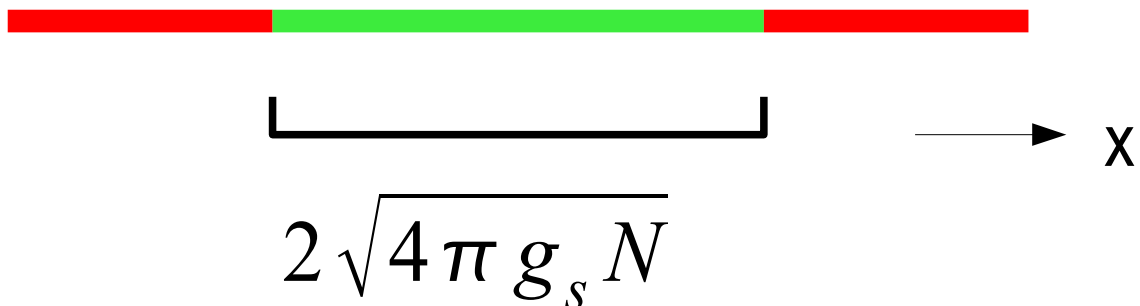
Comparing to the Gaussian matrix model

Vev of circular Wilson loops are related to Gaussian matrix model.

[Erickson, Semenoff, Zarembo], [Drukker, Gross]

Claim: The pattern of x-axis corresponds to the eigen value distribution of the matrix model

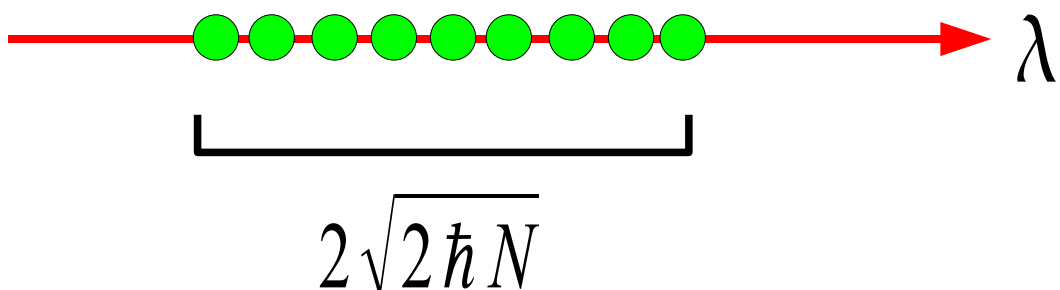
AdS5 x S5 solution



Gaussian matrix model

$$Z = \int dM \exp\left(-\frac{1}{\hbar} \text{Tr}[M^2]\right)$$

The classical distribution of the eigen values of Gaussian matrix model



Wilson line $W = \frac{1}{N} \text{Tr}[U^k],$

Geometry side [Drukker, Fiol]

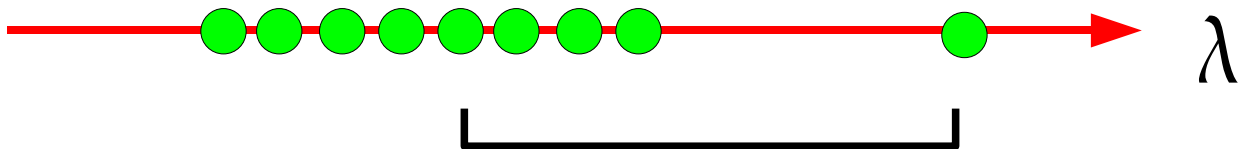
AdS2 x S2 D3-brane, k unit of electric flux



$$\sqrt{2 \cdot 2\pi g_s N + \frac{1}{4} k^2 (2\pi g_s)^2}$$

Matrix model side

$$\langle W \rangle = \int dM \frac{1}{N} \text{Tr}[e^{kM}] \exp\left(-\frac{1}{\hbar} \text{Tr}[M^2]\right),$$



$$\sqrt{2\hbar N + \frac{1}{4} k^2 \hbar^2}$$

These two picture completely match
if we identify $\hbar = 2\pi g_s$

Two same AdS2 x S2 D3-branes cannot coexist
Exclusion principle

Wilson line $W = \frac{1}{\dim} \text{Tr}_{antisym} [U],$

Geometry side

[Yamaguchi to appear]

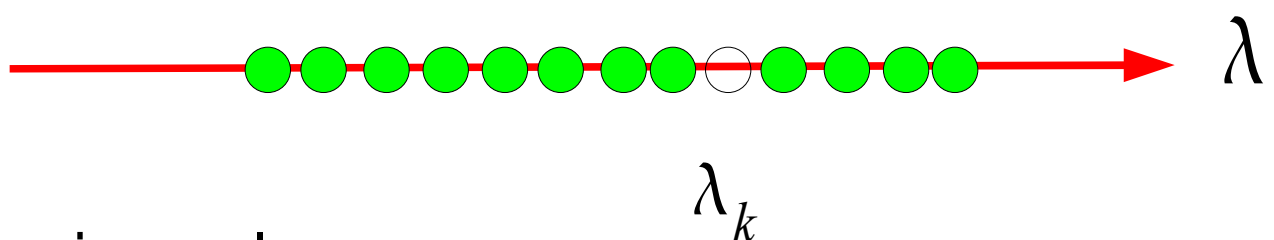
AdS2 x S4 D5-brane,



The position of D5-brane with F-string charge k

$$x_k = \sqrt{4\pi g_s N} \cos \theta, \quad k = \frac{2N}{\pi} \left(\frac{1}{2} \theta - \frac{1}{4} \sin 2\theta \right),$$

Matrix model side

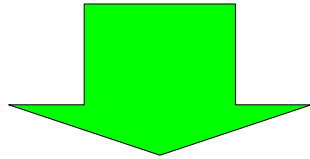


k -th eigenvalue

$$\lambda_k = \sqrt{2\hbar N} \cos \theta, \quad k = \frac{2N}{\pi} \left(\frac{1}{2} \theta - \frac{1}{4} \sin 2\theta \right),$$

8. Solutions in M-theory

Surface operator in 6dim(2,0)CFT
Wall operator in 3dim CFT



$SO(2,2) \times SO(4) \times SO(4)$ symmetry
16 SUSY

Anzats

$SO(2,2)$

$SO(4)$

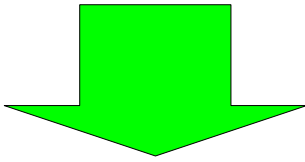
$SO(4)$

$$ds^2 = e^{2A} (AdS_3) + e^{2B} (S^3) + e^{2C} (S^3) + \underbrace{ds_2^2}$$

Unknown

Results

SUSY

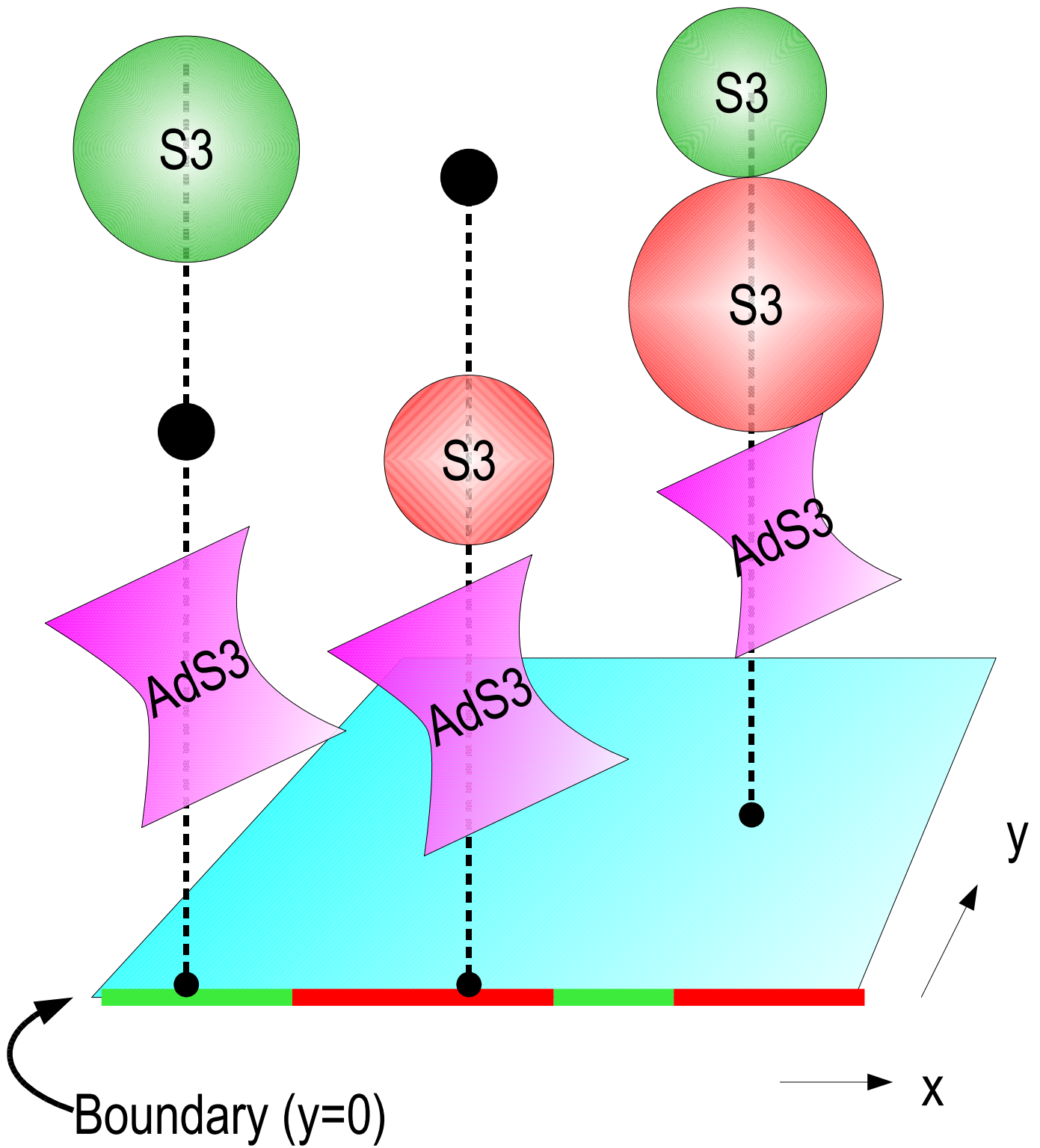


Necessary conditions

$y := \exp(A + B + C)$ Product of radii of
AdS₃, S³ and S³

$$ds^2 = e^{2A} (AdS_3) + e^{2B} (S^3) + e^{2C} (S^3) + ds_2^2$$

$$ds_2^2 = (-e^{2B+2C} + e^{2A+2B} + e^{2A+2C})^{-1} \\ \times (dy^2 + dx^2)$$



$y = (\text{radius AdS3}) \times (\text{radius S3}) \times (\text{radius S3})$

AdS3 never shrinks

A segment \longrightarrow S4



Examples

AdS7 x S4



→ X

AdS4 x S7



→ X

AdS3 x S3 x R4 x R

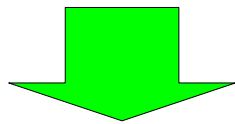


→ X

9. Conclusion

Necessary condition for supergravity solutions that preserves the same symmetry as $\frac{1}{2}$ BPS Wilson line operators.

Supergravity solutions



(Continuous) Maya diagram



The space of eigen values
of the Gaussian Matrix model

→ D-brane exclusion principle

The similar problem in M-theory.