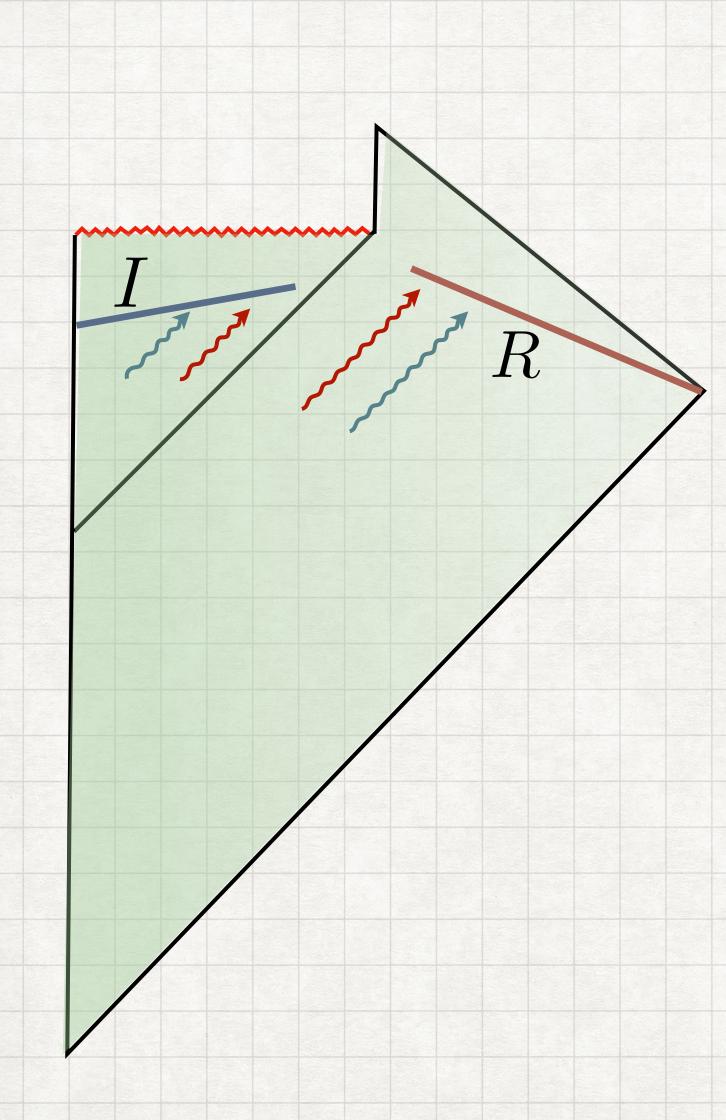
# Islands and Cosmology

Review talk

RIKEN

Kanato Goto



# Island Formula

[Penington, Almheiri-Engelhardt-Marolf-Maxfield '19]

The entropy of Hawking radiation in quantum gravity

$$S_{\mathrm{Rad}} = \min \operatorname{ext}_{I} \left[ \frac{\operatorname{Area}(\partial I)}{4G_{N}} + S_{\mathrm{matter}} \left( I \cup R \right) - S_{\mathrm{ct}}(\partial I) \right]$$

At late stage of the black hole evaporation,

$$S_{
m Rad} pprox rac{{
m Area(horizon)}}{4G_N} = S_{
m BH} 
ightarrow 0$$

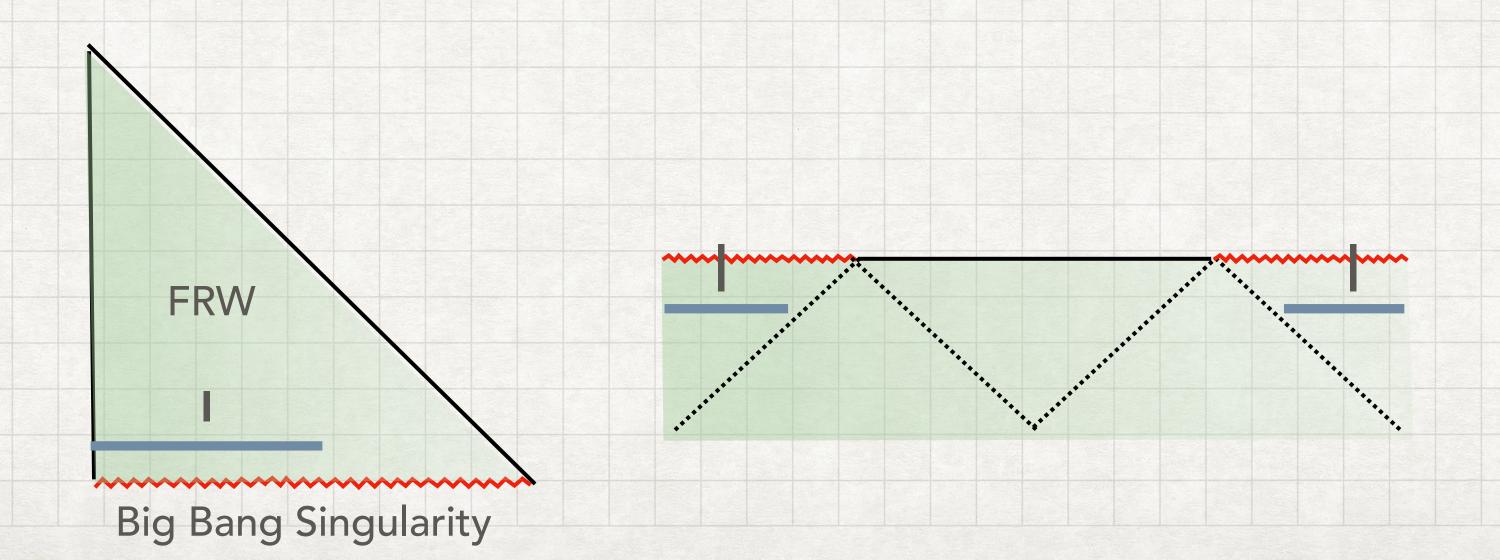
The island region I is encoded in Hawking radiation!

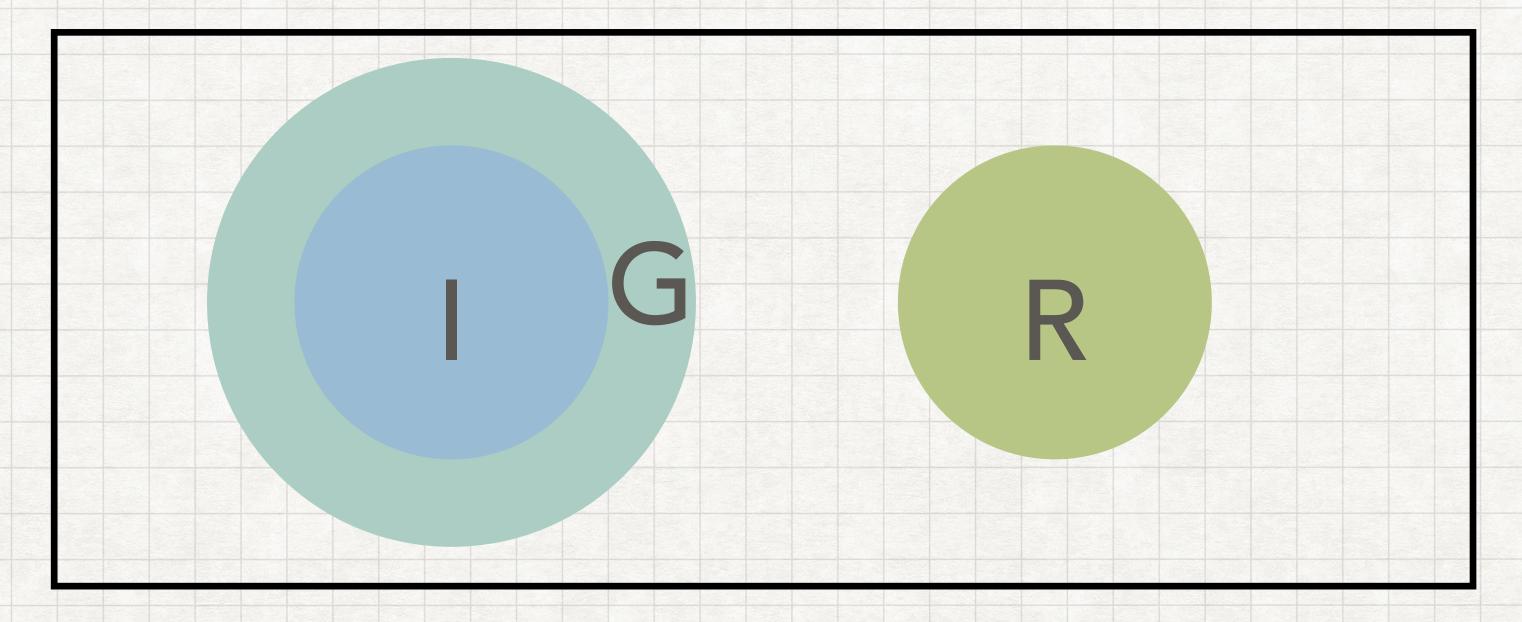
In quantum gravity, information is stored in a non-local form.

- Q. Do the effects of the existence of islands (wormholes) affect us living in a low energy world Are they observable?
  - Any implications for early universe, where the effects of quantum gravity are crucial?

In this talk I will review one (direct) application of the island formula to cosmology, arXiv:2008.01022 "Island in cosmology".

In this paper, the Island formula is applied to the FRW universe and de Sitter spacetime.





- 1. Bekenstein bound is violated for the region I
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- 3. A region G that surrounds island I and has a common boundary with it exists in the quantum normal region

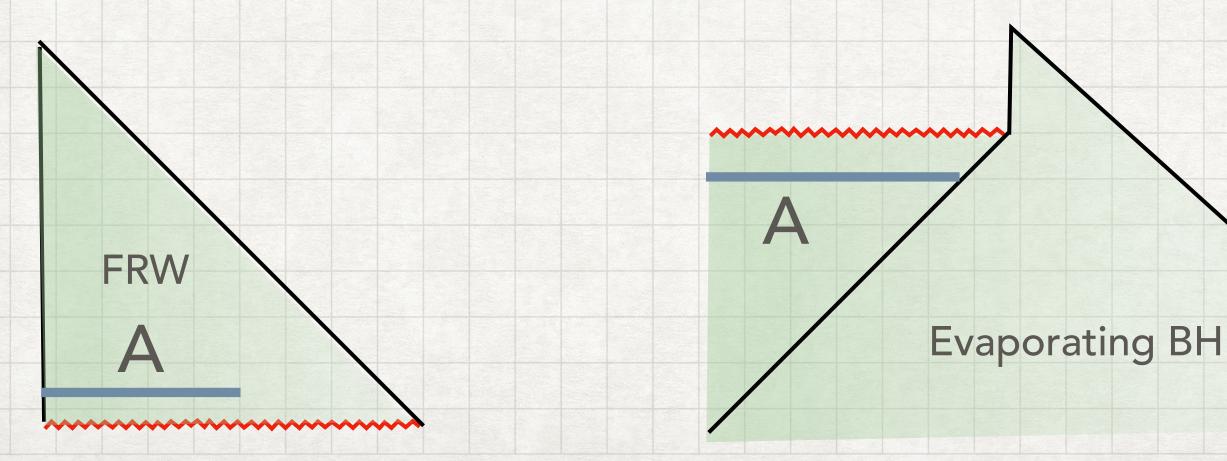
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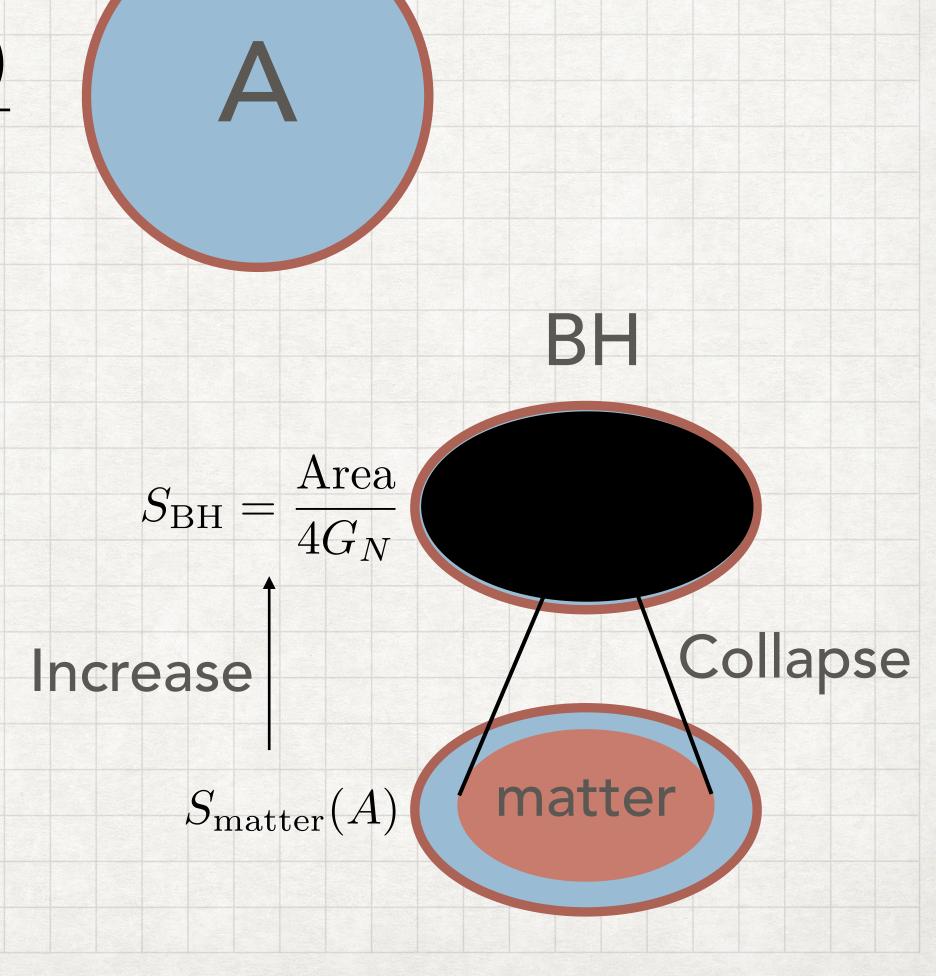
Bekenstein bound

For a region A defined in quantum gravity:

$$S_{\mathrm{matter}}(A) \le \frac{\mathrm{Area}(\partial A)}{4G_N}$$

- The foundation of the holographic principle (sometimes called "holographic bound")
- · It can be violated `near the singularity"

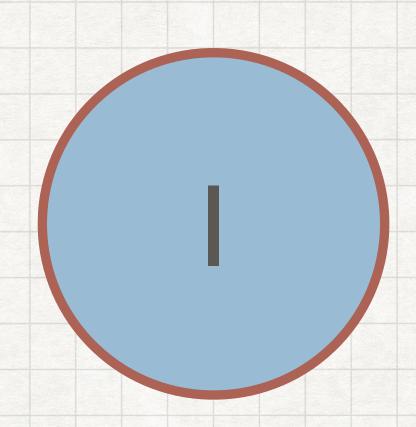




 $\partial A$ 

#### 1. Bekenstein bound is violated for the region I

$$S_{
m matter} \ (I) \gtrsim rac{{
m Area}(\partial I)}{4G_N}$$



Conditions for non-trivial islands to dominate

$$\frac{1}{4G_N}\operatorname{Area}(\partial I) + S_{\mathrm{mat}}(I \cup R) - S_{\mathrm{ct}}(\partial I) < S_{\mathrm{mat}}(R)$$

can be re-combined and written as the following UV finite quantities

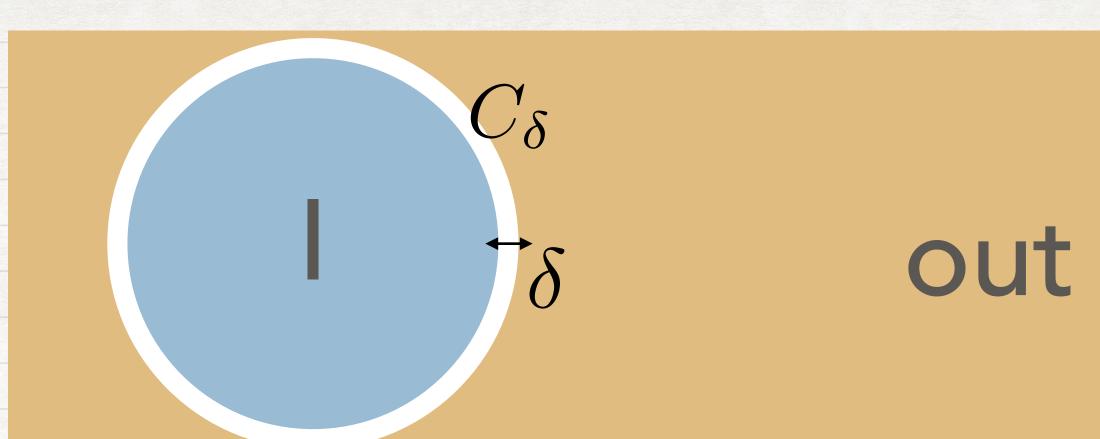
$$I_{\mathrm{mat}}(I,R) \geq S_{\mathrm{gen}}(I)$$

Mutual Information

$$I_{\text{mat}}(I, R) = S_{\text{mat}}(I) + S_{\text{mat}}(R) - S_{\text{mat}}(I \cup R)$$
  
=  $\hat{S}_{\text{mat}}(I) + \hat{S}_{\text{mat}}(R) - \hat{S}_{\text{mat}}(I \cup R)$ 

Generalized entropy

$$\frac{1}{4G_N}\operatorname{Area}(\partial I) + S_{\mathrm{mat}}(I)$$



Strong subadditivity:  $I(I, \text{out}) \ge I(I, R) \ge S_{\text{gen}}(I)$ 

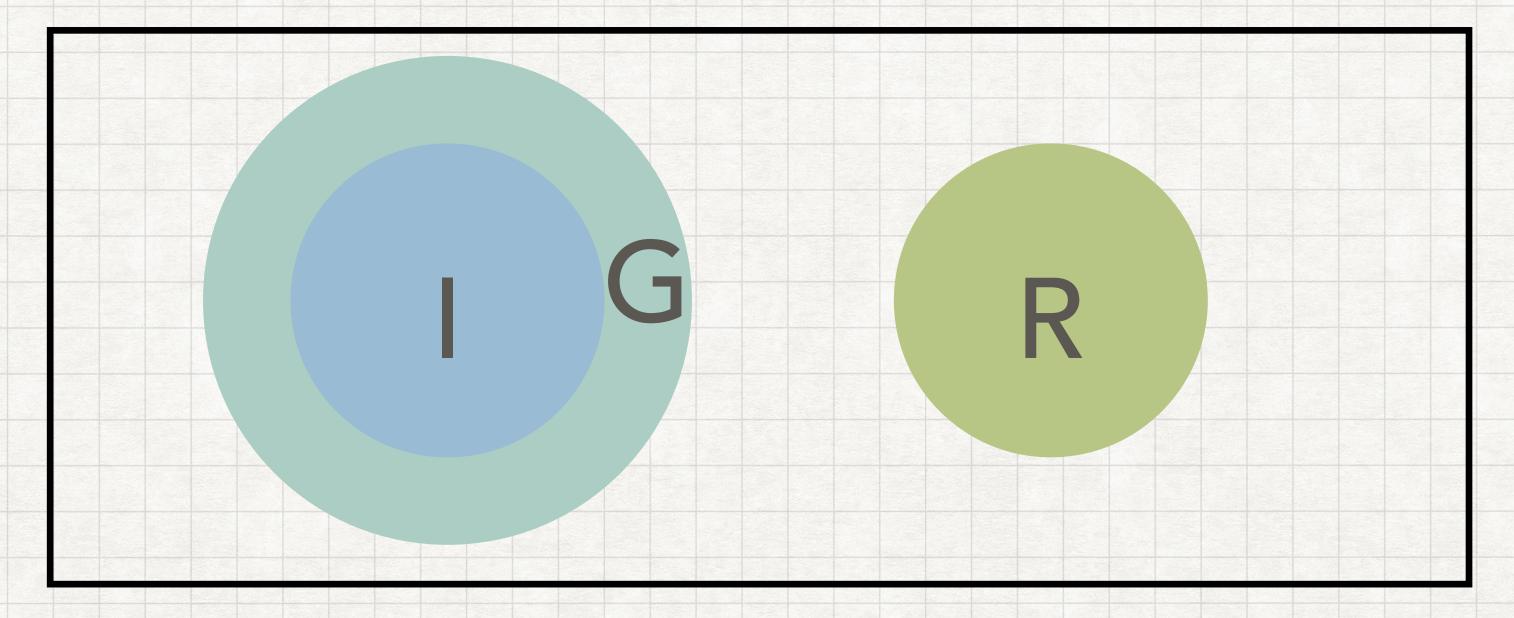
$$= I(I, I \cup C_{\delta}) = \widehat{S}_{\text{mat}}(I) + \widehat{S}_{\text{mat}}(I \cup C_{\delta}) - \widehat{S}_{\text{mat}}(C_{\delta})$$

$$\simeq 2\widehat{S}_{\text{mat}}(I) + \text{(subleading)}$$

$$\delta \to 0$$

これからBekenstein boundの破れが導かれる

$$\widehat{S}_{\mathrm{mat}}(I) \gtrsim rac{\mathrm{Area}(\partial I)}{4G_N}$$



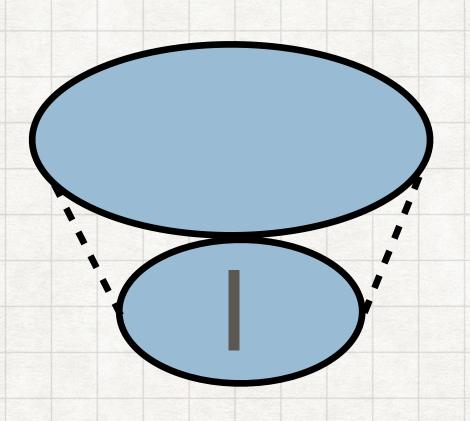
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#### 2. The island I exists in the quantum normal region

Classical normal region ("untrapped" surface)

 $\frac{d}{d\lambda^{+}} \text{Area} \ge 0$ Deformation in outward null direction

Deformation in inward null direction 
$$\dfrac{d}{d\lambda^-} {\rm Area} \leq 0$$

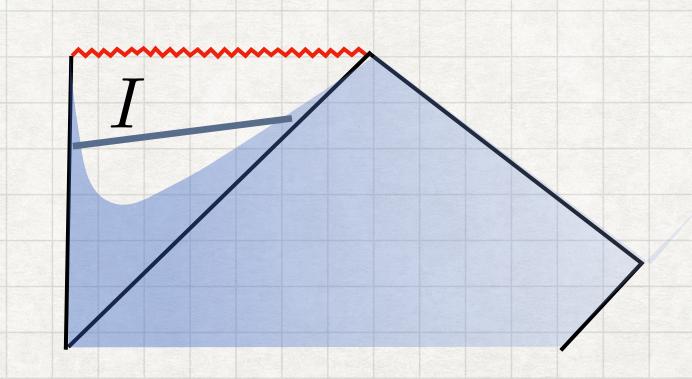


#### Quantum normal region

Deformation in outward null direction  $\frac{d}{d\lambda^+}S_{\mathrm{gen}} \geq 0$ 

Deformation in inward null direction

$$rac{d}{d\lambda} S_{
m gen} \leq 0$$



2. The island I exists in the quantum normal region

Extremal condition  $\frac{d}{d\lambda}S_{\mathrm{gen}}(R\cup I)=0$  for the island

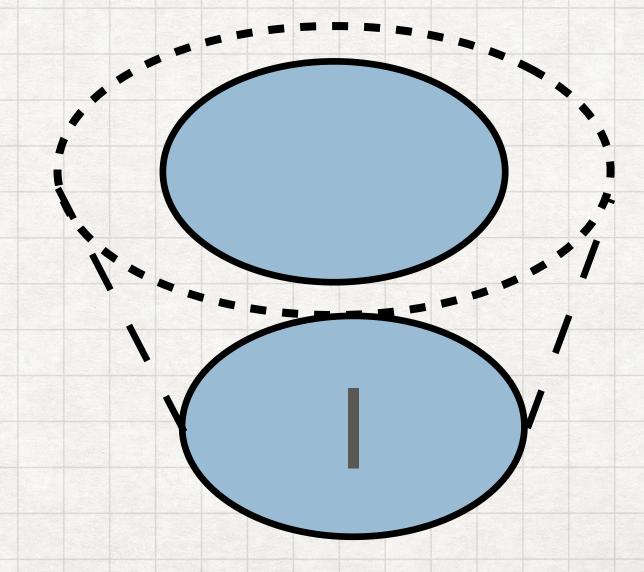
$$\frac{d}{d\lambda} \left[ S_{\text{mat}}(I) + \frac{\text{Area}(\partial I)}{4G_N} - S_{\text{ct}}(\partial I) \right] = \frac{d}{d\lambda} \left[ S_{\text{mat}}(I) + S_{\text{mat}}(R) - S_{\text{mat}}(I \cup R) \right] = 0$$

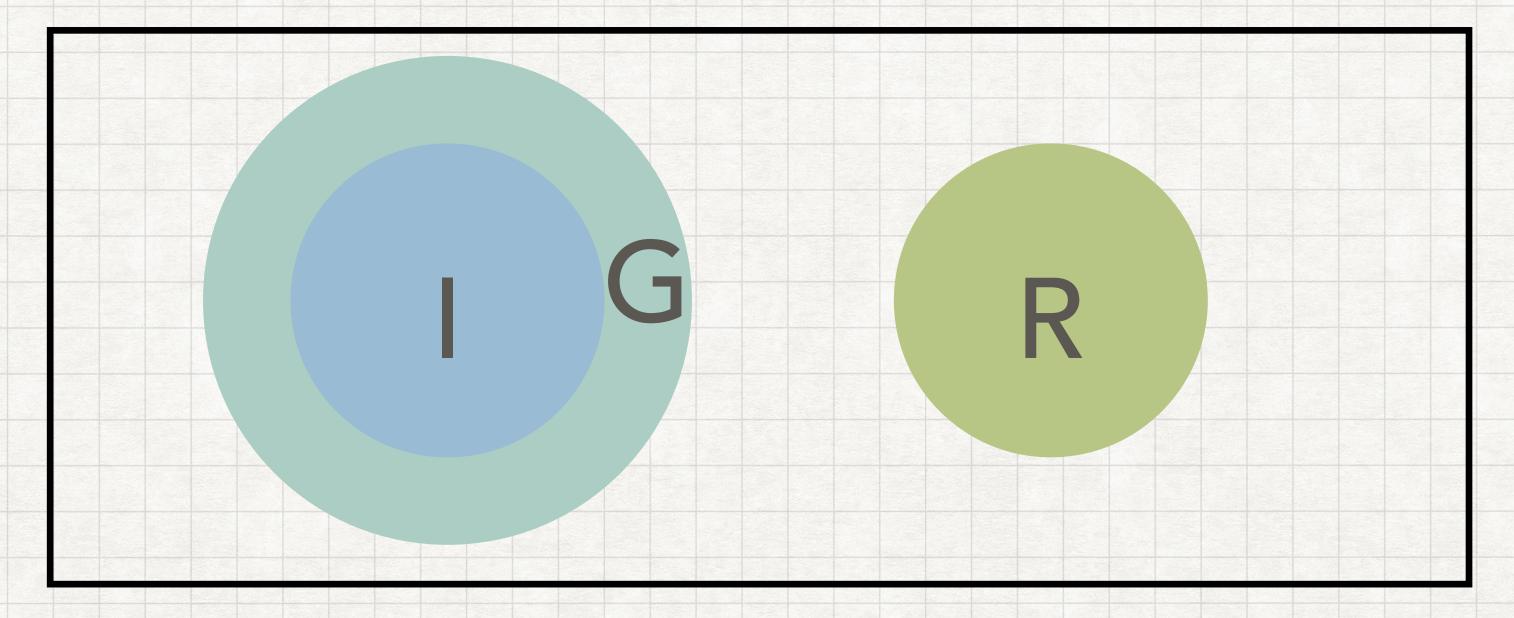
$$= S_{\text{gen}}(I)$$

$$= I(I, R)$$

For the outward null deformation, the r.h.s. is positive due to the monotonicity of the mutual information

$$\pm \frac{d}{d\lambda_{\pm}} S_{\text{gen}} (I) \ge 0$$

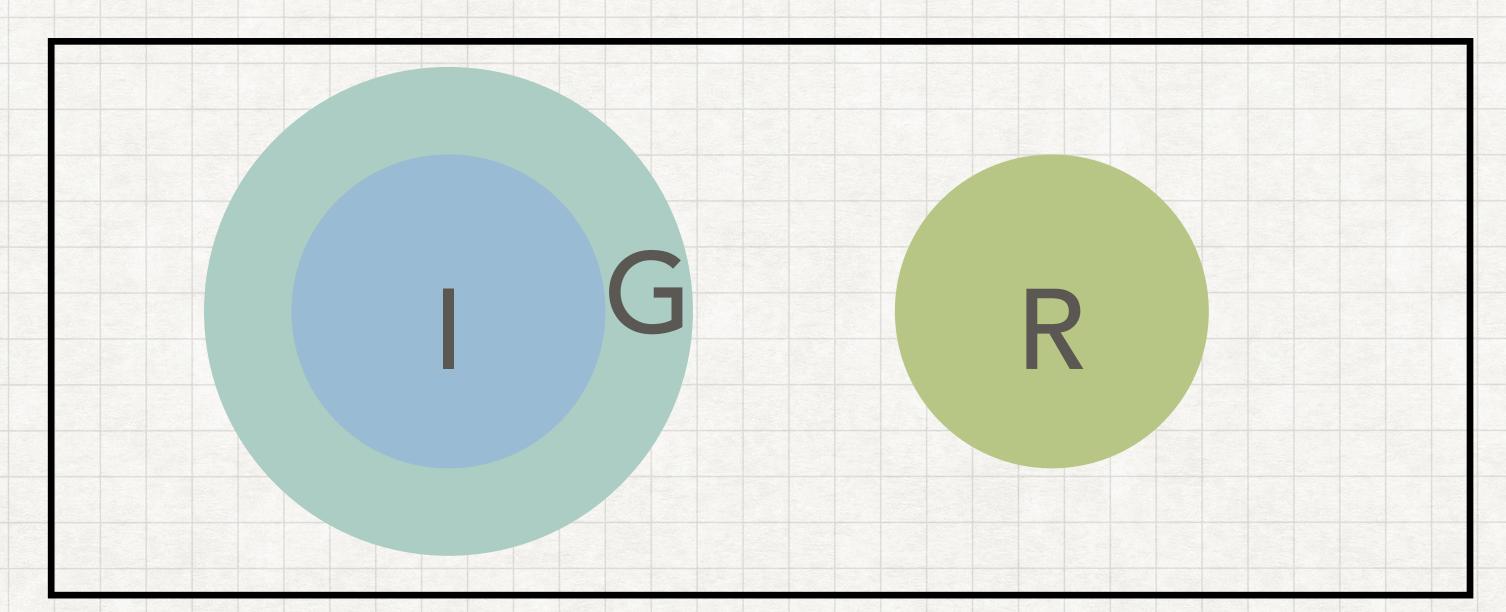




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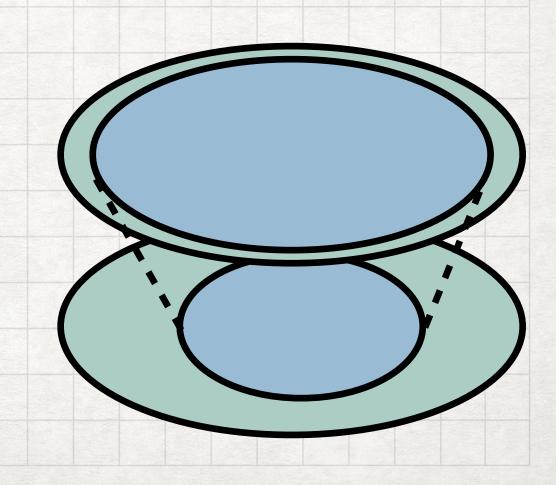


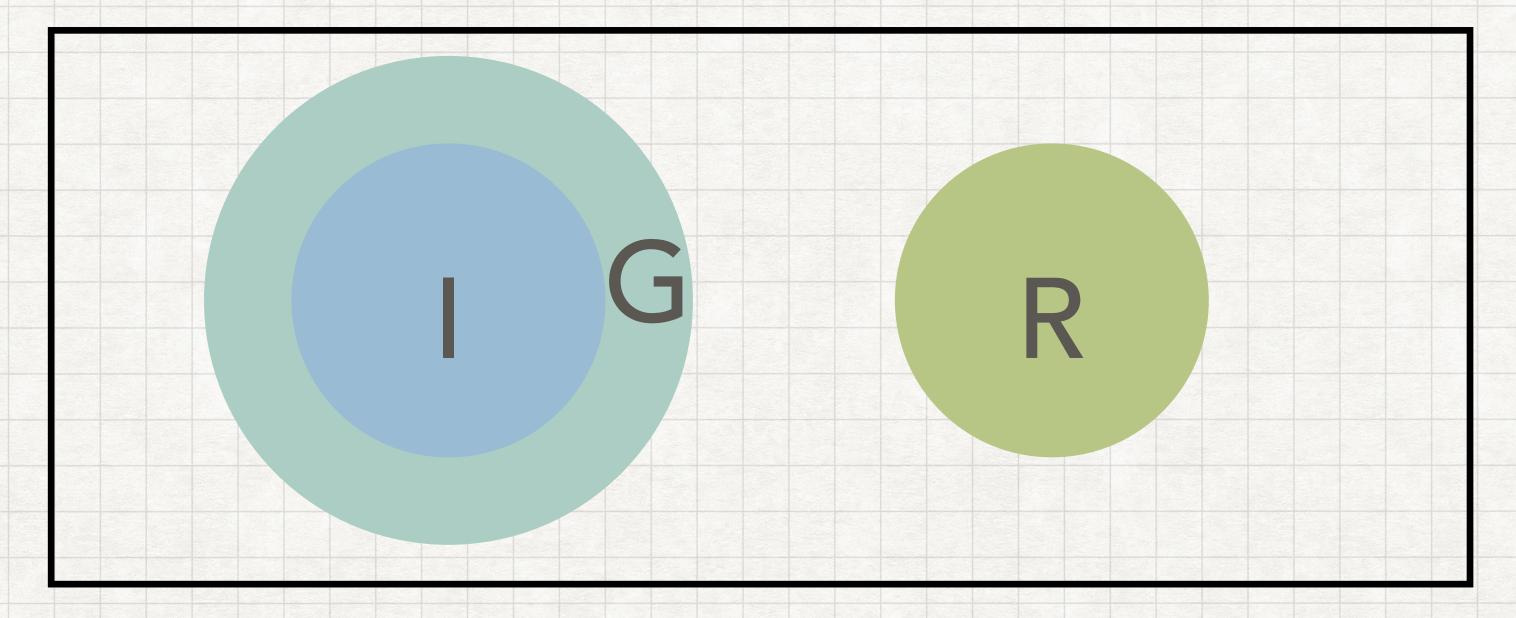
$$\pm \frac{d}{d\lambda_{\pm}} S_{\text{gen}}(G) \le 0$$

Derived from the extremal condition of the island & SSA of entropy

Outside of a sphere usually has a larger area when it is deformed "inward" (outward with respect to the sphere)

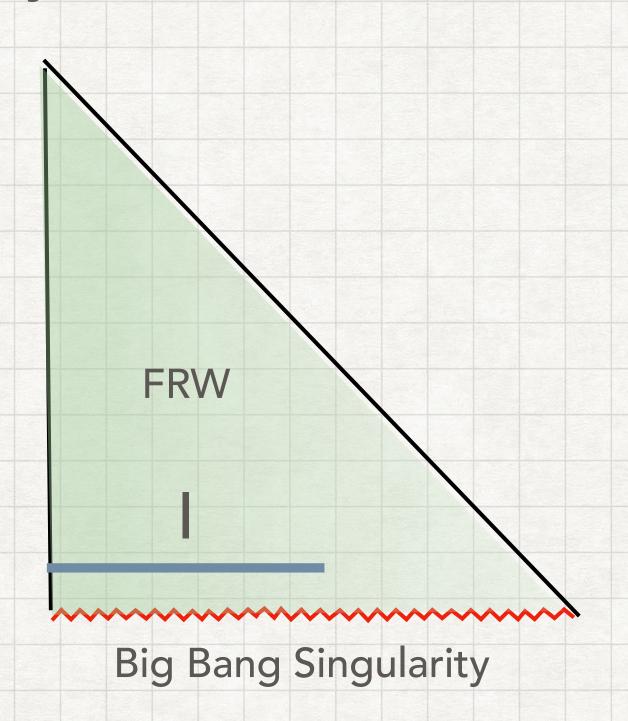
→Requires a large increase of the matter entropy in region G to compensate for the decrease in area





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## Spatially Flat FRW (Friedman-Robertson-Walker) Universe with $\Lambda=0$

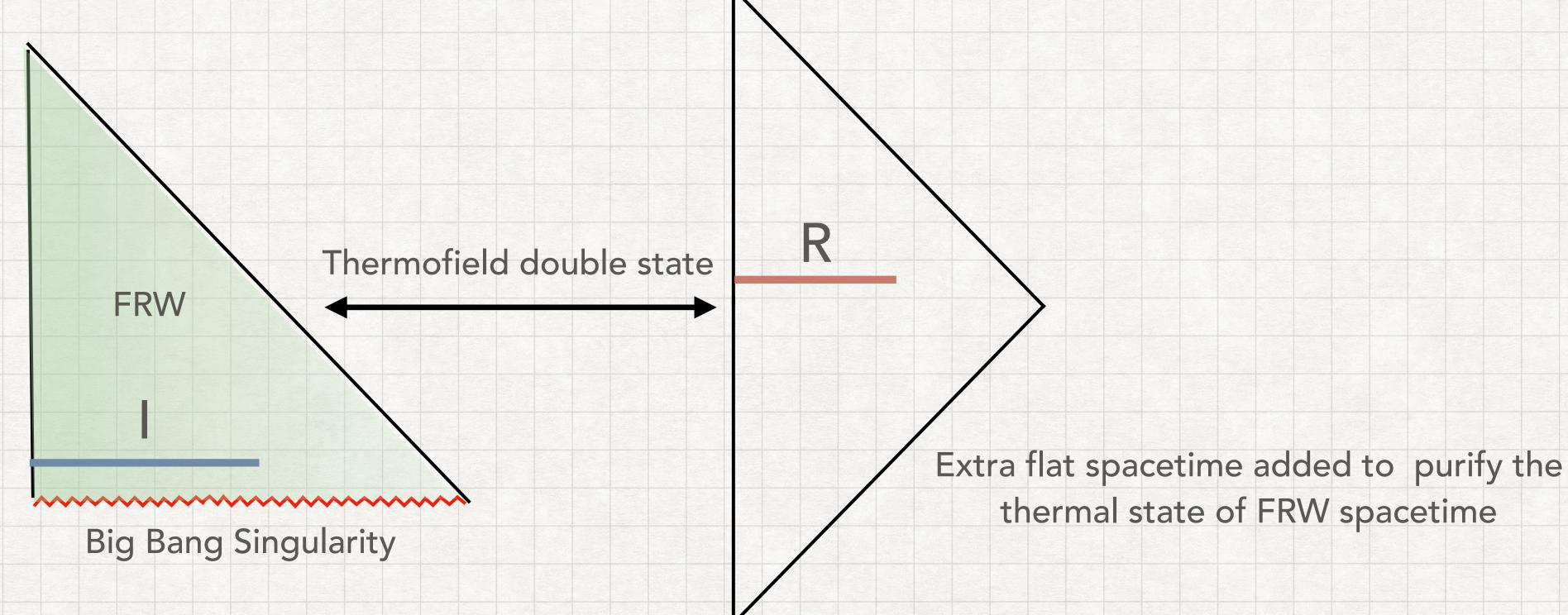


$$ds^{2} = -dt^{2} + a(t)^{2}dx^{2} = a(\eta)^{2}(-d\eta^{2} + dx^{2})$$

Matter is in a thermal state due to the thermal flux from the Big Bang singularity.

$$S_{\mathrm{matter}}(I) \approx s_{th} \mathrm{Vol}(I)$$

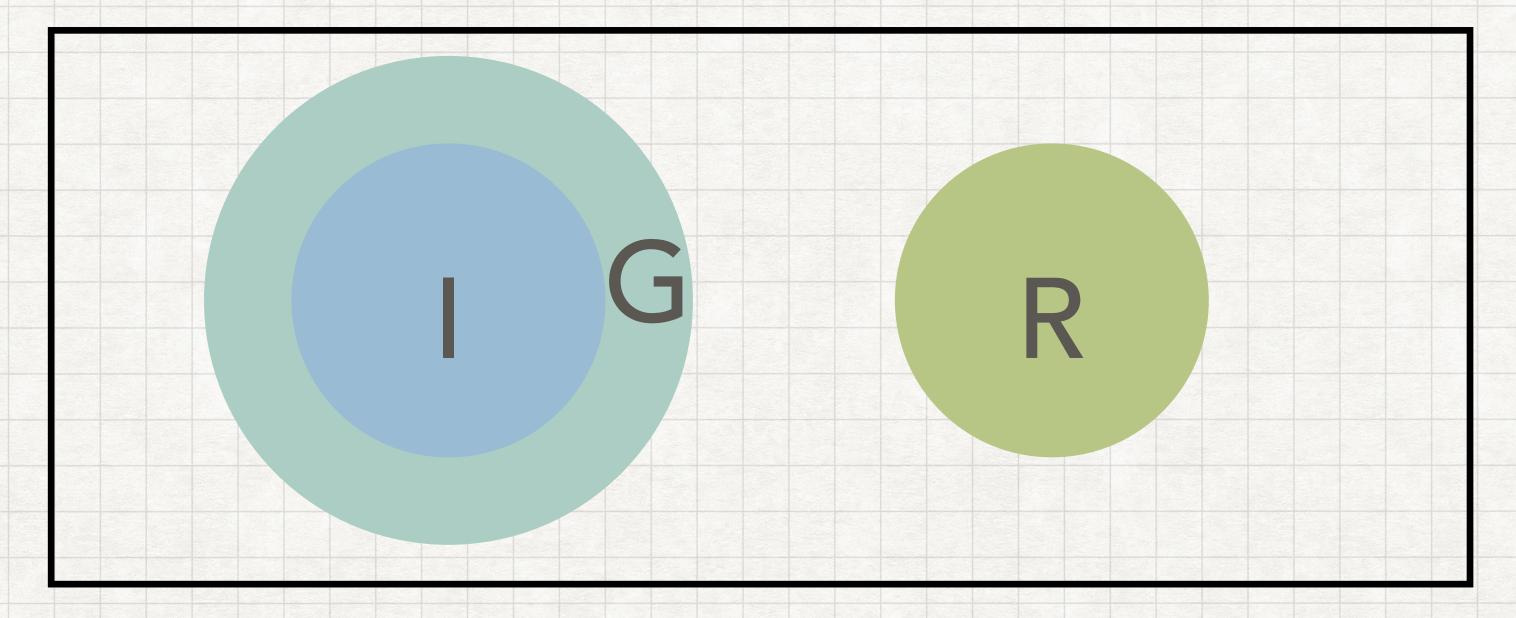




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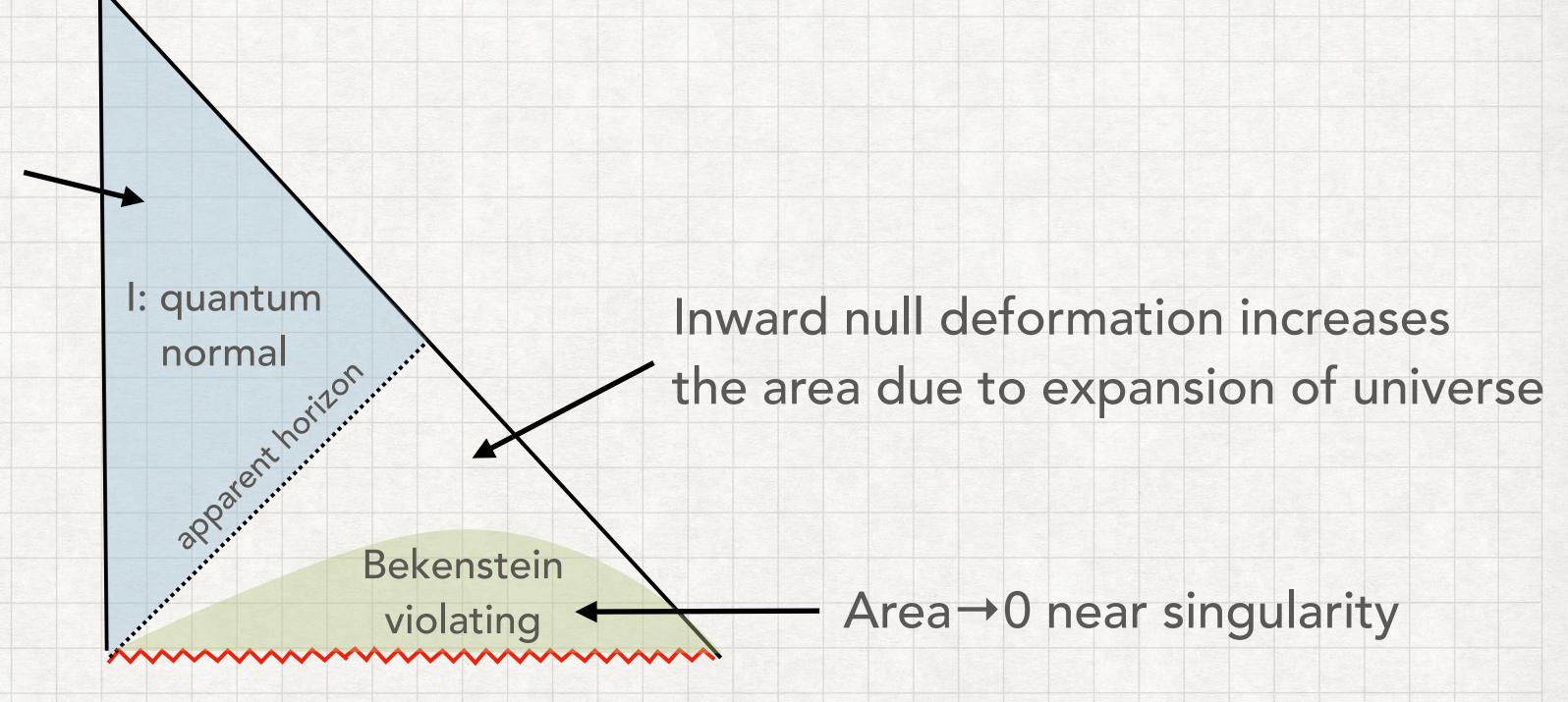
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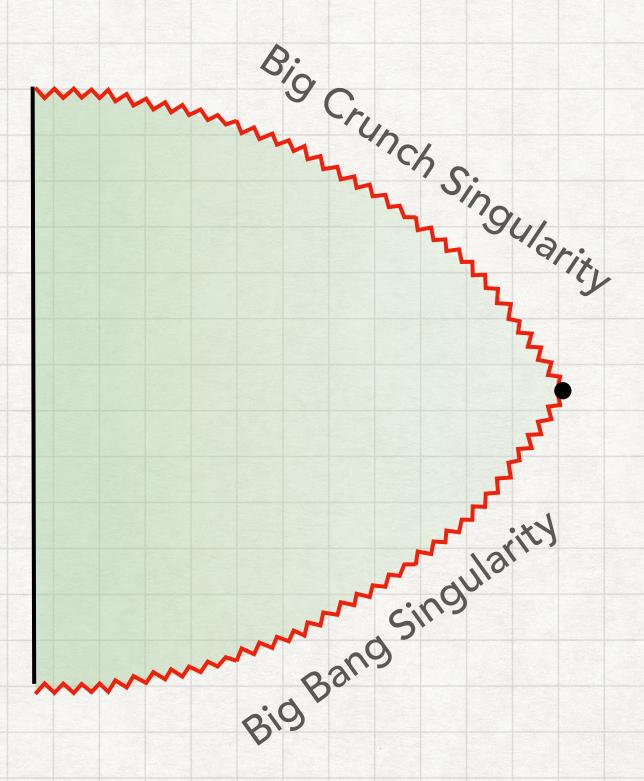
# Spatially Flat FRW (Friedman-Robertson-Walker) Universe with $\Lambda=0$

Quantum normal region: the effect of the expansion of the universe is not significant.



No island region that satisfies the conditions

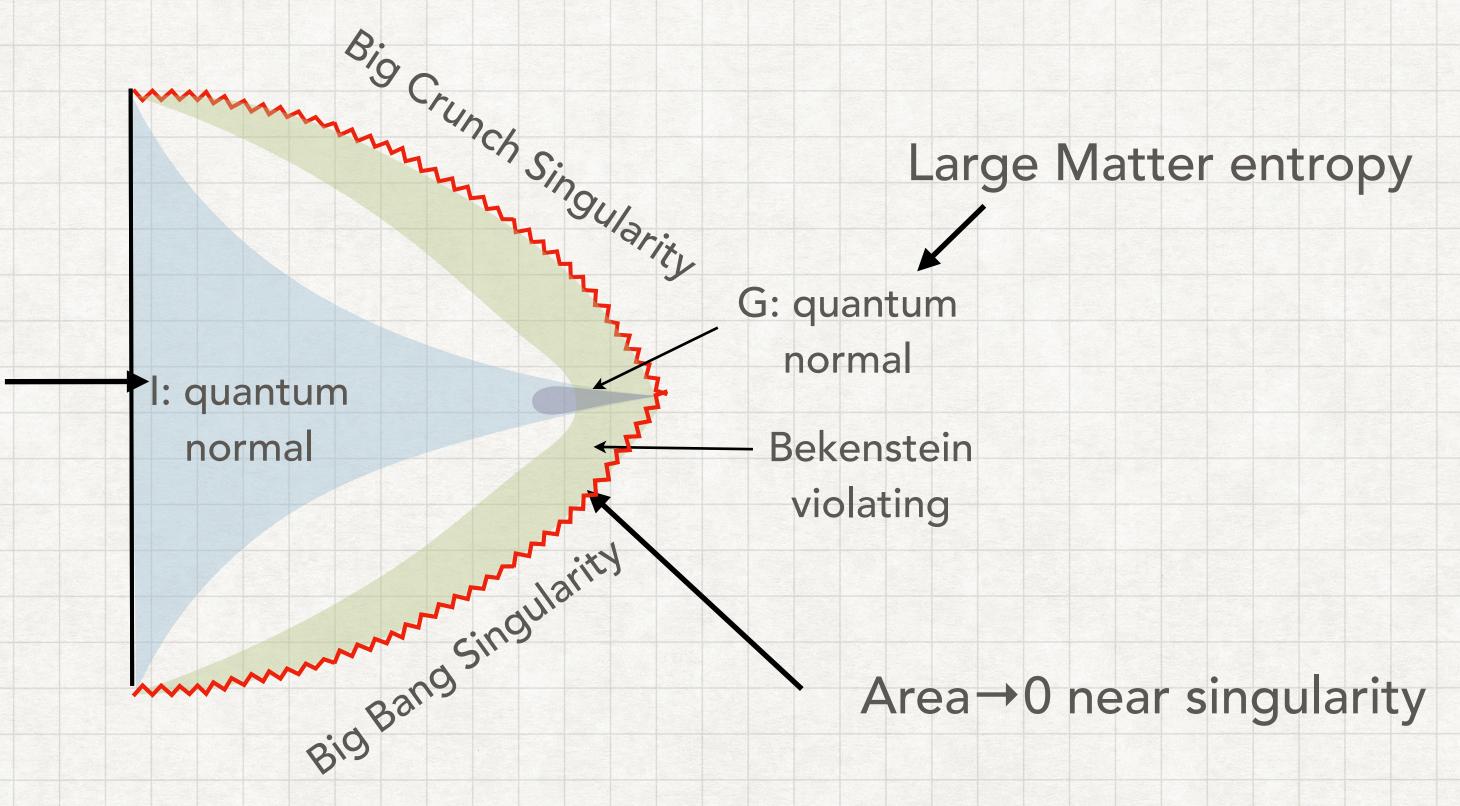




After the Big Bang, this universe expands and then re-collapses

## Spatially Flat FRW (Friedman-Robertson-Walker) Universe with $\Lambda < 0$

Quantum normal region: the effect of inflation & deflation of the universe is not significant.



There is an island region that satisfies the conditions!

Re-collapsing universe can be reconstructed from the extra flat spacetime degrees of freedom

2D de Sitter black hole There is also an island inside a black hole on de Sitter spacetime described by JT gravity. See also Ugajin et al. arXiv:2008.05275 [hep-th] We saw the most direct application of the island formula to early universe

Non-trivial island exists in recollapsing FRW universe, but their interpretation is unclear

In the future

• Can we apply the "idea" of the island formula/wormhole to early universe? & Can we consider observables that have imprints of quantum gravity effects?

cf. Eternal inflation & Baby universes

· Can we extend the idea of the holographic principle?